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## TWO APPLICATIONS OF AN INTEGRAL FORMULA

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We are going to present two consequences of a general integral formula presented in [1].

## 1. Harmonic mappings of Riemannian manifolds

Be given a Riemannian manifold $\left(M, \mathrm{~d} s^{2}\right), \operatorname{dim} M=m$. In a suitable domain $U \subset M$, let us write $(i, j, \ldots=1, \ldots, m)$

$$
\begin{equation*}
\mathrm{d} s^{2}=\sum_{i}\left(\omega^{i}\right)^{2} \tag{1.1}
\end{equation*}
$$

$\omega^{i}$ being linearly independent 1-forms on $U$. Then there are on $U$ 1-forms $\omega_{i}^{j}$ such that

$$
\begin{equation*}
\mathrm{d} \omega^{i}=\sum_{j} \omega^{j} \wedge \omega_{j}^{i}, \quad \omega_{i}^{j}+\omega_{i}^{j}=0 \tag{1.2}
\end{equation*}
$$

the forms $\omega_{i}^{j}$ are uniquely determined by (1.2). The components of the curvature tensor be introduced by

$$
\begin{equation*}
\mathrm{d} \omega_{i}^{j}=\sum_{k} \omega_{i}^{k} \wedge \omega_{k}^{j}-\frac{1}{2} R_{i k l}^{j} \omega^{k} \wedge \omega^{l}, \quad R_{i k l}^{j}+R_{i l k}^{j}=0 \tag{1.3}
\end{equation*}
$$

they satisfy the symmetry relations

$$
\begin{equation*}
R_{i k l}^{j}+R_{j k l}^{i}=0, \quad R_{i k l}^{j}=R_{k i j}^{l}, \quad R_{i k l}^{j}+R_{i l j}^{k}+R_{i j k}^{l}=0 \tag{1.4}
\end{equation*}
$$

Let $v_{1}, \ldots, v_{m}$ be the field of orthonormal frames on $U$ dual to the field of coframes $\omega^{1}, \ldots, \omega^{m}$. Denote by $K\left(v_{i}, v_{j}\right), i \neq j$, the sectional curvature of the 2-plane $\left\{v_{i}, v_{j}\right\}$; of course, $K\left(v_{i}, v_{j}\right)=R_{i i j}^{j}$.

Further, be given another Riemannian manifold $\left(N, \mathrm{~d} \sigma^{2}\right), \operatorname{dim} N=n$, and a mapping $f: M \rightarrow N$. Consider a neighbourhood $V \subset N$ such that $f(U) \subset V$ and there are 1 -forms $\varphi^{\alpha}(\alpha, \beta, \ldots=1, \ldots, n)$ satisfying

$$
\begin{equation*}
\mathrm{d} \sigma^{2}=\sum_{\alpha}\left(\varphi^{\alpha}\right)^{2} \tag{1.5}
\end{equation*}
$$

Let

$$
\begin{equation*}
\tau^{\alpha}:=f^{*} \varphi^{\alpha}=\sum_{i} A_{i}^{\alpha} \omega^{i}, \quad \tau_{\alpha}^{\beta}:=f^{*} \varphi_{\alpha}^{\beta} . \tag{1.6}
\end{equation*}
$$

The exterior differentiation of $\left(1.6_{1}\right)$ yields

$$
\begin{equation*}
\sum_{i}\left(\mathrm{~d} A_{i}^{\alpha}-\sum_{j} A_{j}^{\alpha} \omega_{i}^{j}+\sum_{\beta} A_{i}^{\beta} \tau_{\beta}^{\alpha}\right) \wedge \omega^{i}=0, \tag{1.7}
\end{equation*}
$$

and, according to E. Cartan's lemma, we get the existence of functions $A_{i j}^{\alpha}$ on $U$ satisfying

$$
\begin{equation*}
\mathrm{d} A_{i}^{\alpha}-\sum_{j} A_{j}^{\alpha} \omega_{i}^{j}+\sum_{\beta} A_{i}^{\beta} \tau_{\beta}^{\alpha}=\sum_{j} A_{i j}^{\alpha} \omega^{j}, \quad A_{i j}^{\alpha}=A_{j i}^{\alpha} \tag{1.8}
\end{equation*}
$$

A furhter exterior differentiation implies

$$
\begin{gather*}
\sum_{j}\left(\mathrm{~d} A_{i j}^{\alpha}-\sum_{k} A_{i k}^{\alpha} \omega_{j}^{k}-\sum_{k} A_{k j}^{\alpha} \omega_{i}^{k}+\sum_{\beta} A_{i j}^{\beta} \tau_{\beta}^{\alpha}\right) \wedge \omega^{j}= \\
=\frac{1}{2} \sum_{j, k}\left(\sum_{l} A_{l}^{\alpha} R_{i j k}^{l}-\sum_{\beta, \gamma, \delta} A_{i}^{\beta} A_{j}^{\gamma} A_{k}^{\delta} S_{\beta \gamma \delta}^{\alpha}\right) \omega^{j} \wedge \omega^{k}, \tag{1.9}
\end{gather*}
$$

$S_{\beta \gamma \delta}^{\alpha}$ being the components of the curvature tensor of $\left(N, \mathrm{~d} \sigma^{2}\right)$. Thus there are functions $A_{i j k}^{\chi}$ such that

$$
\begin{gather*}
\mathrm{d} A_{i j}^{\alpha}-\sum_{k} A_{i k}^{\alpha} \omega_{j}^{k}-\sum_{k} A_{k j}^{\alpha} \omega_{i}^{k}+\sum_{\beta} A_{i j}^{\beta} \tau_{\beta}^{\alpha}=\sum_{k} A_{i j k}^{\alpha} \omega^{k}, \quad A_{i j k}^{\alpha}=A_{j i k}^{\alpha},  \tag{1.10}\\
A_{i j k}^{\alpha}-A_{i k j}^{\alpha}=\sum_{l} A_{l}^{\alpha} R_{i k j}^{l}-\sum_{\beta, \gamma, \delta} A_{i}^{\beta} A_{k}^{\gamma} A_{j}^{\delta} S_{\beta \gamma \delta}^{\alpha} . \tag{1.11}
\end{gather*}
$$

Let us consider on $U$ the 1 -forms

$$
\begin{equation*}
\varphi_{1}=\sum_{\alpha, i, j} A_{i}^{\alpha} A_{i j}^{\alpha} \omega^{j}, \quad \varphi_{2}=\sum_{\alpha, i, j} A_{j}^{\alpha} A_{i i}^{\alpha} \omega^{j} . \tag{1.12}
\end{equation*}
$$

It is easy to see that the forms (1.12) are globally defined over all of $M$. The usual *-operator be defined by

$$
\begin{gather*}
* \omega^{i}=(-1)^{i+1} \omega^{1} \wedge \ldots \wedge \omega^{i-1} \wedge \omega^{i+1} \wedge \ldots \wedge \omega^{n},  \tag{1.13}\\
\text { i.e., d } t:=\omega^{1} \wedge \ldots \wedge \omega^{n}=\omega^{i} \wedge * \omega^{i} .
\end{gather*}
$$

Now,

$$
\begin{align*}
& \mathrm{d} * \varphi_{1}=\sum_{\alpha, i, j}\left\{\left(A_{i j}^{\alpha}\right)^{2}+A_{i}^{\alpha} A_{j i j}^{\alpha}\right\} \mathrm{d} o, \\
& \mathrm{~d} * \varphi_{2}=\sum_{\alpha, i, j}\left(A_{i i}^{\alpha} A_{j j}^{\alpha}+A_{i}^{\alpha} A_{j j i}^{\alpha}\right) \mathrm{d} o, \tag{1.14}
\end{align*}
$$

and, according to (1.11),

$$
\begin{equation*}
\mathrm{d} *\left(\varphi_{1}-\varphi_{2}\right)=\sum_{\alpha, i, j}\left\{\left(A_{i j}^{\alpha}\right)^{2}-A_{i i}^{\alpha} A_{j j}^{\alpha}+\sum_{k} A_{i}^{\alpha} A_{k}^{\alpha} R_{j j i}^{k}-\sum_{\beta, \gamma, \delta} A_{j}^{\alpha} A_{i}^{\dot{\beta}} A_{i}^{\gamma} A_{j}^{\delta} S_{\beta \gamma \delta}^{\alpha}\right\} \mathrm{d} o . \tag{1.15}
\end{equation*}
$$

Let us turn our attention to the geometrical interpretation of the above introduced invariants. Let $p \in U \subset M$ be a given point. The Euclidean connection on $M$ or $N$ resp. is given by

$$
\begin{array}{rll}
\nabla m=\sum_{i} \omega^{i} v_{i}, & \nabla v_{i}=\sum_{j} \omega_{i}^{j} v_{j} & \text { or }  \tag{1.16}\\
\nabla^{*} n=\sum_{\alpha} \varphi^{\alpha} w_{\alpha}, & \nabla^{*} w_{\alpha}=\sum_{\beta} \varphi_{\alpha}^{\beta} w_{\beta} & \text { resp.; }
\end{array}
$$

here, $w_{1}, \ldots, w_{n}$ is the dual basis to $\varphi^{1}, \ldots, \varphi^{n}$. Evidently,

$$
\begin{equation*}
\mathrm{d} f_{p}\left(v_{i}\right)=A_{i}^{\alpha} w_{\alpha} \tag{1.17}
\end{equation*}
$$

Let $v \in T_{p}(M)$ be a non-zero vector. Choose a curve $\gamma:(-\varepsilon, \varepsilon) \rightarrow M$ such that $\gamma(0)=p$; let $s$ be its arc and $v$ its tangent vector at $p$. Denote by $\gamma^{*}=f \circ \gamma:(-\varepsilon, \varepsilon) \rightarrow$ $\rightarrow N$ the corresponding curve. Then it is easy to see that

$$
\begin{equation*}
\frac{\nabla^{*} n}{\mathrm{~d} s^{2}}-\mathrm{d} f_{p}\left(\frac{\nabla^{2} m}{\mathrm{~d} s^{2}}\right)=\frac{L(v)}{|v|^{2}} \tag{1.18}
\end{equation*}
$$

where $|v|^{2}=\sum_{i}\left(\omega^{i}(v)\right)^{2}$ and

$$
\begin{equation*}
L(v)=A_{i j}^{\alpha}(p) \omega^{i}(v) \omega^{j}(v) w_{\alpha}(f(p)) \tag{1.19}
\end{equation*}
$$

This gives the geometrical interpretation of the quadratic mapping

$$
\begin{equation*}
L: T_{p}(M) \rightarrow T_{f(p)}(N) \tag{1.20}
\end{equation*}
$$

Let $L(.$, .) be the corresponding bilinear mapping.
At $p$, let us choose an orthonormal frame $v_{i}$, let $w_{\alpha}$ be an orthonormal frame at $f(p)$. Then

$$
\begin{equation*}
L\left(v_{i}, v_{j}\right)=A_{i j}^{\alpha} w_{\alpha} \tag{1.21}
\end{equation*}
$$

and the expressions

$$
\begin{equation*}
\sum_{i, j}\left|L\left(v_{i}, v_{j}\right)\right|^{2}=\sum_{i, j, \alpha}\left(A_{i j}^{\alpha}\right)^{2}, \quad\left|\sum_{i} L\left(v_{i}\right)\right|^{2}=\sum_{\alpha, i, j i} A_{i i}^{\alpha} A_{j j}^{\alpha} \tag{1.22}
\end{equation*}
$$

do not depend on the choice of the frames $v_{i}$ and $w_{\alpha}$. In the same way, the vector

$$
\begin{equation*}
t=\sum_{i} L\left(v_{i}\right) \tag{1.23}
\end{equation*}
$$

is invariant; the mapping

$$
\begin{equation*}
t: M \rightarrow T(N), \quad t(p) \in T_{f(p)}(N) \tag{1.24}
\end{equation*}
$$

is the so-called tension field. The mapping $f: M \rightarrow N$ is said to be harmonic if $t=0$ for each $p \in M$.

The frames $\left(v_{1}, \ldots, v_{m}\right)$ and $\left(w_{1}, \ldots, w_{t}\right)$ at $p$ and $f(p)$ resp. are called adapted to $f$ if

$$
\begin{array}{ll}
\mathrm{d} f_{p}\left(v_{i}\right)=A_{i} w_{i} & \text { for } i=1, \ldots, m \text { in the case } m \leqq n \text { and } \\
\mathrm{d} f_{p}\left(v_{\alpha}\right)=A_{\alpha} \mathrm{w}_{\alpha} & \text { for } \alpha=1, \ldots, n,  \tag{1.25}\\
\mathrm{~d} f_{p}\left(v_{Q}\right)=0 & \text { for } \varrho=n+1, \ldots, m \text { in the case } m>n
\end{array}
$$

Thus, we may always write $\left(25_{1}\right)$ setting $w_{i}=0$ for $i>n$. The adapted bases exist for each couple $(p, f(p))$. In the adapted bases, we have

$$
\begin{gather*}
\sum_{\alpha, i, j} \sum_{k} A_{i}^{\alpha} A_{k}^{\alpha} R_{j j i}^{k}=\sum_{i}\left(A_{i}\right)^{2} \sum_{j \neq i} K\left(v_{j}, v_{i}\right),  \tag{1.26}\\
\sum_{\alpha, i, j, j, \gamma, \delta} \sum_{j, \delta}^{\alpha} A_{j}^{\alpha} A_{i}^{\beta} A_{i}^{\gamma} A_{j}^{\delta} S_{\beta \gamma \delta}^{\alpha}=2 \sum_{i \neq j}\left(A_{i} A_{j}\right)^{2} K^{*}\left(w_{i}, w_{j}\right) . \tag{1.27}
\end{gather*}
$$

Further,

$$
\begin{equation*}
\varphi_{1}\left(v_{i}\right)=\sum_{j}\left\langle\mathrm{~d} f\left(v_{j}\right), L\left(v_{i}, v_{j}\right)\right\rangle, \quad \varphi_{2}\left(v_{i}\right)=\left\langle\mathrm{d} f\left(v_{i}\right), t\right\rangle, \tag{1.28}
\end{equation*}
$$

$\langle$,$\rangle being the scalar product in T_{f(p)}(N)$.
Choosing for each couple ( $p, f(p)$ ) the adapted bases, we have the integral formula

$$
\begin{gather*}
\int_{\partial M} *\left(\varphi_{1}-\varphi_{2}\right)=\int_{M}\left\{\sum_{i, j}\left|L\left(v_{i}, v_{j}\right)\right|^{2}-|t|^{2}+\right. \\
\left.+\sum_{i}\left(A_{i}\right)^{2} \sum_{j \neq i} K\left(v_{j}, v_{i}\right)-2 \sum_{i \neq j}\left(A_{i} A_{j}\right)^{2} K^{*}\left(w_{i}, w_{j}\right)\right\} \mathrm{d} v . \tag{1.29}
\end{gather*}
$$

Thus we get the following
Theorem. Let $M, N$ be Riemannian manifolds and $f: M \rightarrow N$ a harmonic mapping. Let $N$ have non-positive sectional curvatures and let $M$ have, at each point $p \in M$ and for each unit vector $v \in T_{p}(M)$ the following property: $v_{1}, \ldots, v_{m-1}, v$ being an orthonormal basis of $T_{p}(M)$, we have $\sum_{r=1, \ldots, m-1} K\left(v, v_{r}\right)>0$. Let $\varphi_{1}=\varphi_{2}$ on the boundary $\partial M$ of $M$. Then $f$ is a constant mapping.
2. Holomorphic curves in the Hermitian plane

Be given a Hermitian plane $H^{2}$ and let $m: D \rightarrow H^{2}$ be a holomorphic curve, $D \subset \mathscr{C}$ being a bounded domain. To each its point $m(d), d \in D$, let us associate an orthonormal frame $\left\{m, w_{1}, w_{2}\right\}$. Then we have the equations

$$
\begin{gather*}
\mathrm{d} m=\tau^{1} w_{1}+\tau^{2} w_{2} \\
\mathrm{~d} w_{1}=\tau_{1}^{1} w_{1}+\tau_{1}^{2} w_{2}, \quad \mathrm{~d} w_{2}=\tau_{2}^{1} w_{1}+\tau_{2}^{2} w_{2} \tag{2.1}
\end{gather*}
$$

clearly $(i, j, \ldots=1,2)$

$$
\begin{gather*}
\tau_{i}^{j}+\bar{\tau}_{j}^{i}=0,  \tag{2.2}\\
\mathrm{~d} \tau^{i}=\tau^{j} \wedge \tau_{j}^{i}, \quad \mathrm{~d} \tau_{i}^{j}=\tau_{i}^{k} \wedge \tau_{k}^{j} . \tag{2.3}
\end{gather*}
$$

Let us restrict ourselves to the tangent frames satisfying

$$
\begin{equation*}
\tau^{2}=0 \tag{2.4}
\end{equation*}
$$

By successive exterior differentiations we get the existence of functions $R, S, T, U: D \rightarrow \mathscr{C}$ such that

$$
\begin{gather*}
\tau_{1}^{2}=R \tau^{1}  \tag{2.5}\\
\mathrm{~d} R+R\left(\tau_{2}^{2}-2 \tau_{1}^{1}\right)=S \tau^{1},  \tag{2.6}\\
\mathrm{~d} S+S\left(\tau_{2}^{2}-3 \tau_{1}^{1}\right)+3 R^{2} \bar{R} \bar{\tau}^{1}=T \tau^{1}  \tag{2.7}\\
\mathrm{~d} T+T\left(\tau_{2}^{2}-4 \tau_{1}^{1}\right)+10 R \bar{R} S \bar{\tau}^{1}=U \tau^{1} . \tag{2.8}
\end{gather*}
$$

Let us consider another field of orthonormal frames

$$
\begin{equation*}
u_{1}=e^{i \alpha} w_{1}, \quad u_{2}=e^{i \beta} w_{2} ; \quad \alpha, \beta: D \rightarrow \mathscr{R} ; \tag{2.9}
\end{equation*}
$$

let us write

$$
\begin{equation*}
\mathrm{d} m=\varphi^{1} u_{1}, \quad \mathrm{~d} u_{1}=\varphi_{1}^{1} u_{1}+\varphi_{1}^{2} u_{2}, \quad \mathrm{~d} u_{2}=\varphi_{2}^{1} u_{1}+\varphi_{2}^{2} u_{2} . \tag{2.10}
\end{equation*}
$$

Then it is easy to see that

$$
\begin{gather*}
\varphi^{1}=e^{-i \alpha} \tau^{1}  \tag{2.11}\\
\varphi_{1}^{1}=\tau_{1}^{1}+i \mathrm{~d} \alpha, \quad \varphi_{2}^{2}=\tau_{2}^{2}+i \mathrm{~d} \beta, \quad \varphi_{1}^{2}=e^{i(\alpha-\beta)} \tau_{1}^{2} \tag{2.12}
\end{gather*}
$$

Write

$$
\begin{gather*}
\varphi_{1}^{2}=R^{\prime} \varphi_{1}  \tag{2.13}\\
\mathrm{~d} R^{\prime}+R^{\prime}\left(\varphi_{2}^{2}-2 \varphi_{1}^{1}\right)=S^{\prime} \varphi^{1}  \tag{2.14}\\
\mathrm{~d} S^{\prime}+S^{\prime}\left(\varphi_{2}^{2}-3 \varphi_{1}^{1}\right)+3 R^{\prime 2} \bar{R}^{\prime} \bar{\varphi}^{1}=T^{\prime} \varphi^{1}  \tag{2.15}\\
\mathrm{~d} T^{\prime}+T^{\prime}\left(\varphi_{2}^{2}-4 \varphi_{1}^{1}\right)+10 R^{\prime} \bar{R}^{\prime} S^{\prime} \bar{\varphi}^{1}=U^{\prime} \varphi^{1} \tag{2.16}
\end{gather*}
$$

Then

$$
\begin{array}{ll}
R^{\prime}=e^{i(2 \alpha-\beta)} R, & S^{\prime}=e^{i(3 \alpha-\beta)} S, \\
T^{\prime}=e^{i(4 \alpha-\beta)} T, & U^{\prime}=e^{i(5 \alpha-\beta)} U . \tag{2.17}
\end{array}
$$

The mappings $B^{(k)}: T_{m} \rightarrow N_{m}$ be introduced by

$$
\begin{gather*}
B\left(z w_{1}\right)=z^{2} R w_{2}, \quad B^{(1)}\left(z w_{1}\right)=z^{3} S w_{2},  \tag{2.18}\\
B^{(2)}\left(z w_{1}\right)=z^{4} T w_{2} ; \quad z \in \mathscr{C} .
\end{gather*}
$$

These mappings are invariant. Indeed: Let $w=z w_{1}=z^{\prime} u_{1}$, then $z^{\prime}=e^{-i \alpha_{z}}$ and $z^{\prime 2} R^{\prime} u_{2}=z^{2} R w_{2} ;$ similarly for $B^{(k)}$. Let $S^{1}=\left\{w \in T_{m} ;\langle w, w\rangle=1\right\}$, i.e., $S^{1}=$ $=\left\{z w_{1} ;|z|^{2}=1\right\}$. Then $B^{(k)}\left(S^{1}\right)$ is a circle; the radius of $B\left(S^{1}\right)$ is equal to $|R|^{1 / 2}$, the radius of $B^{(1)}\left(S^{1}\right)$ is equal to $|S|^{1 / 2}$, etc. The geometrical interpretation of the mappings $B^{(k)}$ will be presented later on.

The area element of $m$ is given by

$$
\begin{equation*}
\mathrm{d} o=\frac{1}{2} i \tau^{1} \wedge \bar{\tau}^{1} . \tag{2.19}
\end{equation*}
$$

The Hodge operator be introduced by

$$
\begin{equation*}
* \tau^{1}=-i \tau^{1}, \quad * \bar{\tau}^{1}=i \bar{\tau}^{1} . \tag{2.20}
\end{equation*}
$$

Let $f: D \rightarrow \mathscr{R}$ be a function. Then its Laplacian $\Delta f$ is given, as usually, by

$$
\begin{equation*}
\Delta f \mathrm{~d} o=\mathrm{d} * \mathrm{~d} f \tag{2.21}
\end{equation*}
$$

The straightforward calculations lead to ( $n \geqq 1$ )

$$
\begin{gather*}
\mathrm{d}|R|^{2 n}=2 n|R|^{2 n-2} \operatorname{Re}\left(\bar{R} S \tau^{1}\right),  \tag{2.22}\\
\Delta|R|^{2 n}=4 n|R|^{2 n-2}\left(n|S|^{2}-3|R|^{4}\right),  \tag{2.23}\\
\mathrm{d}|S|^{2 n}=2 n|S|^{2 n-2} \operatorname{Re}\left\{\left(\bar{S} T-3 S R \bar{R}^{2}\right) \tau^{1}\right\}, \tag{2.24}
\end{gather*}
$$

$$
\begin{align*}
\Delta|S|^{2 n}= & 4 n\left\{|S|^{2 n-2}\left(n|T|^{2}-16|S|^{2}|R|^{2}+9 n|R|^{6}\right)-\right. \\
& \left.-6(n-1)|S|^{2 n-4}|R|^{2} \operatorname{Re}\left(\bar{S}^{2} R T\right)\right\} . \tag{2.25}
\end{align*}
$$

Especially,

$$
\begin{align*}
& \Delta|R|^{2}=4\left(|\mathrm{~S}|^{2}-3|R|^{4}\right), \\
& \Delta|R|^{4}=8|R|^{2}\left(2|S|^{2}-3|R|^{4}\right),  \tag{2.26}\\
& \Delta|S|^{2}=4\left(|T|^{2}-16|S|^{2}|R|^{2}+9|R|^{6}\right)
\end{align*}
$$

and

$$
\begin{equation*}
\Delta\left(|S|^{2}+4|R|^{4}\right)=4\left(|T|^{2}-15|R|^{6}\right) \tag{2.27}
\end{equation*}
$$

Lemma. Let $S=0$ on $D$. Then $m(D)$ is a part of a straight line of $H^{2}$.
Proof. The equation (2.7) implies $R=0$. QED.
Theorem. Let $S=0$ on $\partial D$ and

$$
\begin{equation*}
3|R|^{4} \geqq|S|^{2} \quad \text { on } D . \tag{2.28}
\end{equation*}
$$

Then $m(D)$ is a part of a straight line of $H^{2}$.
Proof. Obviously, $* \mathrm{~d}|R|^{2 n}=0$ on $\partial D$. From the integral formula

$$
\begin{equation*}
0=\int_{M} \Delta|R|^{2} \mathrm{~d} v \tag{2.29}
\end{equation*}
$$

we get, because of (2.26),

$$
\begin{equation*}
3|R|^{4}=|S|^{2} \quad \text { on } D . \tag{2.30}
\end{equation*}
$$

The integral formula

$$
\begin{equation*}
\int_{\partial M} * \mathrm{~d}|R|^{4}=8 \int_{M}|R|^{2}\left(2|S|^{2}-3|R|^{4}\right) \mathrm{d} v \tag{2.31}
\end{equation*}
$$

reduces to

$$
\begin{equation*}
0=\int_{M}|R|^{6} \mathrm{~d} v, \tag{2.32}
\end{equation*}
$$

and we get $R=0$. QED.
The formulas (2.23), (2.25), (2.27) imply new characterizations of straight lines of $H^{2}$. It is sufficient to suppose $S=0$ on $\partial D$ and, for ex.,

$$
\begin{equation*}
14|R|^{6} \geqq|T|^{2} \tag{2.33}
\end{equation*}
$$

or

$$
\begin{equation*}
|T|^{2} \geqq 8|R|^{2}\left(2|S|^{2}-|R|^{4}\right) \tag{2.34}
\end{equation*}
$$

on $D$; see (2.27) and (2.26 $)$.
Now, the geometrical description of the mappings $B^{(k)}$ is given in [1]. To do this, let us consider $H^{2}$ as a space over m, i.e., $H^{2}$ becomes $E^{4}$. Write

$$
\begin{gather*}
v_{1}=w_{1}, \quad v_{2}=i w_{1}, \quad v_{3}=w_{2}, \quad v_{4}=i w_{2}, \\
\tau^{1}=\omega^{1}+i \omega^{2}, \quad \tau^{2}=\omega^{3}+i \omega^{4}, \quad \tau_{1}^{2}=\omega_{1}^{3}+i \omega_{1}^{4},  \tag{2.35}\\
\tau_{1}^{1}=i \omega_{1}^{2}, \quad \tau_{2}^{2}=i \omega_{3}^{4},
\end{gather*}
$$

i.e.,

$$
\begin{align*}
\mathrm{d} m & =\omega^{1} v_{1}+\omega^{2}{ }_{2}, \\
\mathrm{~d} v_{1} & =\omega_{1}^{2} v_{2}+\omega_{1}^{3} v_{3}+\omega_{1-4}^{4}, \\
\mathrm{~d} v_{2} & =-\omega_{1}^{2} v_{1} \quad-\omega_{1}^{4} v_{3}+\omega_{1}^{3} v_{4},  \tag{2.36}\\
\mathrm{~d} v_{3} & =-\omega_{1}^{3} v_{1}+\omega_{1}^{4} v_{2}+\omega_{3}^{4} v_{4}, \\
\mathrm{~d} v_{4} & =-\omega_{1}^{4} v_{1}-\omega_{1}^{3} v_{2}-\omega_{3}^{4} v_{3}
\end{align*}
$$

and

$$
\begin{equation*}
\omega_{1}^{3}=R_{1} \omega^{1}-R_{2} \omega^{2}, \quad \omega_{1}^{4}=R^{2} \omega^{1}+R_{1} \omega^{2} \tag{2.37}
\end{equation*}
$$

with $R_{1}=\operatorname{Re} R, R_{2}=\operatorname{Im} R$. In $E^{4}$, consider a general surface

$$
\begin{align*}
\mathrm{d} n & =\varrho^{1} v_{1}+\varrho^{2} v_{2}, \\
\mathrm{~d} v_{1} & =\varrho_{1}^{2} v_{2}+\varrho_{1}^{3} v_{3}+\varrho_{1}^{4} v_{4}, \\
\mathrm{~d} v_{2} & =-\varrho_{1}^{2} v_{1} \quad+\varrho_{2}^{3} v_{3}+\varrho_{2}^{4} v_{4},  \tag{2.38}\\
\mathrm{~d} v_{3} & =-\varrho_{1}^{3} v_{1}-\varrho_{2}^{3} v_{2}+\varrho_{3}^{4} v_{4}, \\
\mathrm{~d} v_{4} & =-\varrho_{1}^{4} v_{1}-\varrho_{2}^{4} v_{2}-\varrho_{3}^{4} v_{3}
\end{align*}
$$

with

$$
\begin{array}{r}
\varrho_{1}^{3}=a_{1} \varrho^{1}+a_{2} \varrho^{2}, \quad \varrho_{2}^{3}=a_{2} \varrho^{1}+a_{3} \varrho^{2}, \\
\varrho_{1}^{4}=b_{1} \varrho^{1}+b_{2} \varrho^{2}, \quad \varrho_{2}^{4}=b_{2} \varrho^{1}+b_{3} \varrho^{2}, \\
\mathrm{~d} a_{1}-2 a_{2} \varrho_{1}^{2}-b_{1} \varrho_{3}^{4}=\alpha_{1} \varrho^{1}+\alpha_{2} \varrho^{2}, \\
\mathrm{~d} a_{2}+\left(a_{1}-a_{3}\right) \varrho_{1}^{2}-b_{2} \varrho_{3}^{4}=\alpha_{2} \varrho^{1}+\alpha_{3} \varrho^{2}, \\
\mathrm{~d} a_{3}+2 a_{2} \varrho_{1}^{2}-b_{3} \varrho_{3}^{4}=\alpha_{3} \varrho^{1}+\alpha_{4} \varrho_{2},  \tag{2.40}\\
\mathrm{~d} b_{1}-2 b_{2} \varrho_{1}^{2}+a_{1} \varrho_{3}^{4}=\beta_{1} \varrho^{1}+\beta_{2} \varrho^{2}, \\
\mathrm{~d} b_{2}+\left(b_{1}-b_{3}\right) \varrho_{1}^{2}+a_{2} \varrho_{3}^{4}=\beta_{2} \varrho^{1}+\beta_{3} \varrho^{2}, \\
\mathrm{~d} b_{3}+2 b_{2} \varrho_{1}^{2}+a_{3} \varrho_{3}^{4}=\beta_{3} \varrho^{1}+\beta_{4} \varrho^{2} .
\end{array}
$$

Then it is known [1] that, for

$$
\begin{aligned}
\Phi & =\left(a_{1} \alpha_{3}+a_{2} \alpha_{4}-a_{2} \alpha_{2}-a_{3} \alpha_{3}+b_{1} \beta_{3}+b_{2} \beta_{4}-b_{2} \beta_{2}-b_{3} \beta_{3}\right) \varrho^{1}+ \\
& +\left(\mathrm{a}_{2} \alpha_{1}+a_{3} \alpha_{2}-a_{1} \alpha_{2}-a_{2} \alpha_{3}+b_{2} \beta_{1}+b_{3} \beta_{2}-b_{1} \beta_{2}-b_{2} \beta_{3}\right) \varrho^{2}
\end{aligned}
$$

we have

$$
\begin{gather*}
\int_{\partial N} * \Phi=\int_{N}\left[2\left(\alpha_{1} \alpha_{3}+\alpha_{2} \alpha_{4}-\alpha_{2}^{2}-\alpha_{3}^{2}+\beta_{1} \beta_{3}+\beta_{2} \beta_{4}-\beta_{2}^{2}-\beta_{3}^{2}\right)-\right. \\
-\left\{\left(a_{1}-a_{3}\right)^{2}+4 a_{2}^{2}+\left(b_{1}-b_{3}\right)^{2}+4 b_{2}^{2}\right\}\left(a_{1} a_{3}-a_{2}^{2}+b_{1} b_{3}-b_{2}^{2}\right)+ \\
\left.+2\left\{b_{2}\left(a_{1}-a_{3}\right)+a_{2}\left(b_{3}-b_{1}\right)\right\}^{2}\right] \mathrm{d} v . \tag{2.42}
\end{gather*}
$$

In our case, (2.42) is identical with

$$
\begin{equation*}
\int_{\partial N} * \mathrm{~d}|R|^{2}=\int_{N} \Delta|R|^{2} \mathrm{~d} v . \tag{2.43}
\end{equation*}
$$

## BIBLIOGRAPHY

[1] Švec A.: On a general integral formula. To appear.

## SOUHRN

## DVĚ APLIKACE JEDNÉ INTEGRÁLNÍ FORMULE

## ALOIS ŠVEC

V práci jsou vyloženy aplikace integrální formule [1] na teorii harmonických zobrazení a na teorii křivky v hermiteovské rovině.

## PE3ЮME

## ДВА ПРИМЕНЕНИЯ ОДНОЙ ИНТЕГРАЛЬНОЙ ФОРМУЛЫ

## АЛОИС ШВЕЦ

В работе излагаются приложения интегральной формулы [1] на теорию гармонических отображений и на теорию кривых в пространствах Эрмита.

