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# CONNECTIONS ON SOME FUNCTIONAL BUNDLES 

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## Introduction

Our starting point was the idea of the Schrödinger connection on a double fibered manifold by Jadczyk and Modugno, [4], [5]. We discuss the "pure case" of two classical fiber bundles $E_{1}$ and $E_{2}$ over the same base and define a connection $\Gamma$ on the bundle $\mathscr{F}\left(E_{1}, E_{2}\right)$ of all smooth maps from a fiber of $E_{1}$ into the fiber of $E_{2}$ over the same base point. We study systematically the geometry of the iterated tangent bundle of the infinite dimensional space $\mathscr{F}\left(E_{1}, E_{2}\right)$ as well as the jet prolongations of $\mathscr{F}\left(E_{1}, E_{2}\right)$ by means of the ideas introduced by the second author in [9]. Since we deal with functional bundles, our vector fields and connections represent a kind of differential operators. That is why we pay special attention to the case of finite order operators, in which we are able to deduce a very concrete description of the objects and operations in question.

In such a situation we found the simplest way for introducing the curvature of $\Gamma$ in a construction by Ehresmann, [2], which is based on the notion of semiholonomic 2 -jets. In the new context we were obliged to rearrange some results, deduced in the finite dimension by direct evaluation, into a more geometrical setting, which could be generalized to our infinite dimensional case. Only then we study the bracket of two vector fields on $\mathscr{F}\left(E_{1}, E_{2}\right)$. This is a modification of the bracket of two vertical prolongation operators on a classical fibered manifold by Kosmann-Schwarzbach, [11], and the second author, [8]. In Proposition 14 we deduce a satisfactory bracket formula for the curvature of $\Gamma$. We also discuss the absolute differentiation with respect to $\Gamma$ and the special case $E_{2}$ is a vector bundle.

If we deal with two finite dimensional manifolds and a map between them, we always assume they are of class $C^{\infty}$, i.e. smooth in the classical sense. On the other

[^0]hand, the idea of smoothness in the infinite dimension is taken from the theory of smooth structures by Frölicher, [3].

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## 1. The tangent bundle of $\mathscr{F}\left(E_{1}, E_{2}\right)$

Let $p_{1}: E_{1} \rightarrow M$ and $p_{2}: E_{2} \rightarrow M$ be two classical fiber bundles (i.e. locally trivial fibered manifolds) over the same base. Consider the set of all fiber maps

$$
\mathscr{F}\left(E_{1}, E_{2}\right)=\bigcup_{x \in M} C^{\infty}\left(E_{1 x}, E_{2 x}\right)
$$

and denote by $p: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow M$ the canonical projection. We define no topology on $\mathscr{F}\left(E_{1}, E_{2}\right)$, but we introduce the concept of a smooth map from a classical manifold $Q$ into $\mathscr{F}\left(E_{1}, E_{2}\right)$.

Definition 1. A map $f: Q \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ is called smooth, if
(i) $p \circ f: Q \rightarrow M$ is smooth and
(ii) the induced map $\tilde{f}:(p \circ f)^{*} E_{1} \rightarrow E_{2}$,

$$
\tilde{f}(q, y)=f(q)(y), \quad(q, y) \in(p \circ f)^{*} E_{1}
$$

is also smooth.
As usual, $(p \circ f)^{*} E_{1} \rightarrow Q$ denotes the bundle induced from $E_{1}$ by means of $p \circ f$, i.e.

$$
(p \circ f)^{*} E_{1}=\left\{(q, y) \in Q \times E_{1} \mid(p \circ f)(q)=p_{1}(y)\right\}
$$

Thus, $\mathscr{F}\left(E_{1}, E_{2}\right)$ is endowed with a smooth structure in the sense of Frölicher, [3].
For every smooth curve $f: \mathbb{R} \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ we first construct the tangent vector $X=\left.\frac{\partial}{\partial t}\right|_{0}(p \circ f) \in T M$ of its base map at $t=0$. Write

$$
T_{X} E_{1}=\left(T p_{1}\right)^{-1}(X) \subset T E_{1} \quad \text { or } \quad T_{X} E_{2}=\left(T p_{2}\right)^{-1}(X) \subset T E_{2}
$$

so that $T_{X} E_{1}$ or $T_{X} E_{2}$ is an affine bundle over $E_{1 x}$ or $E_{2 x}, x=p(f(0))$, with the derived vector bundle $T\left(E_{1 x}\right):=V_{x} E_{1}$ or $T\left(E_{2 x}\right):=V_{x} E_{2}$, respectively. Then $f$ defines a map $T_{0} f: T_{X} E_{1} \rightarrow T_{X} E_{2}$ by

$$
\begin{equation*}
T_{0} f\left(\left.\frac{\partial}{\partial t}\right|_{0} h(t)\right)=\left.\frac{\partial}{\partial t}\right|_{0} f(t)(h(t)) \tag{1}
\end{equation*}
$$

where we may assume that $h: \mathbb{R} \rightarrow E_{1}$ satisfies $p \circ f=p_{1} \circ h$.

Definition 2. We say that two smooth curves $f, g: \mathbb{R} \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ satisfying $\left.\frac{\partial}{\partial t}\right|_{0} p \circ f=\left.\frac{\partial}{\partial t}\right|_{0} p \circ g=X$ determine the same tangent vector at $f(0)=g(0)=\varphi$, if

$$
T_{0} f=T_{0} g: T_{X} E_{1} \rightarrow T_{X} E_{2}
$$

The set $T \mathscr{F}\left(E_{1}, E_{2}\right)$ of all equivalence classes will be called the tangent bundle of $\mathscr{F}\left(E_{1}, E_{2}\right)$.

We write $\left.\frac{\partial}{\partial t}\right|_{0} f(t) \in T \mathscr{F}\left(E_{1}, E_{2}\right)$ for the tangent vector determined by $f$ and $\pi: T \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ and $T p: T \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T M$ for the canonical projections. If $A \in T \mathscr{F}\left(E_{1}, E_{2}\right)$, then we denote by $\tilde{A}: T_{T p(A)} E_{1} \rightarrow T_{T p(A)} E_{2}$ the associated map (1).

Remark 1. Let $\varepsilon \subset \mathscr{F}\left(E_{1}, E_{2}\right)$ be any subset. Then we define $T \varepsilon \subset T \mathscr{F}\left(E_{1}, E_{2}\right)$ by restricting ourselves to the smooth curves with values in $\varepsilon$.

One sees easily that $T_{0} f=T_{0} g: T_{X} E_{1} \rightarrow T_{X} E_{2}$ is an affine bundle morphism over the base map $\varphi: E_{1 x} \rightarrow E_{2 x}$ with the derived linear morphism $T \varphi: T\left(E_{1 x}\right) \rightarrow$ $T\left(E_{2 x}\right)$. Indeed, let $x^{i}$ be some local coordinates on $M, y^{p}$ or $z^{a}$ be some additional coordinates on $E_{1}$ or $E_{2}$ and

$$
\begin{equation*}
x^{i}=f^{i}(t), \quad z^{a}=f^{a}\left(y^{p}, t\right) \tag{2}
\end{equation*}
$$

be the coordinate expression of $f(t)$. Write

$$
Y^{p}=\mathrm{d} y^{p}, \quad Z^{a}=\mathrm{d} z^{a}, \quad \varphi^{a}(y)=f^{a}(y, 0), \quad \Phi^{a}(y)=\frac{\partial f^{a}\left(y^{p}, 0\right)}{\partial t}
$$

Then the coordinate form of (1) is

$$
\begin{equation*}
Z^{a}=\frac{\partial \varphi^{a}(y)}{\partial y^{p}} Y^{p}+\Phi^{a}(y) \tag{3}
\end{equation*}
$$

Hence the tangent vector to (2) is locally characterized by two systems of numbers and two systems of functions

$$
\begin{equation*}
x^{i}=f^{i}(0), \quad X^{i}=\frac{\partial f^{i}(0)}{\partial t}, \quad \varphi^{a}\left(y^{p}\right), \quad \Phi^{a}\left(y^{p}\right) \tag{4}
\end{equation*}
$$

The following lemma gives a global assertion of such a type.

Lemma 1. Let $F: T_{X} E_{1} \rightarrow T_{X} E_{2}$ be an affine bundle morphism over $\varphi: E_{1 x} \rightarrow$ $E_{2 x}$ with the derived linear morphism $T \varphi: T\left(E_{1 x}\right) \rightarrow T\left(E_{2 x}\right)$. Then there exists a smooth curve $f: \mathbb{R} \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ such that $F=\tilde{A}$ for the tangent vector $A=$ $\left.\frac{\partial}{\partial t}\right|_{0} f(t)$.

Proof. Consider some local trivializations $U \times S_{1}$ and $U \times S_{2}$ of $E_{1}$ and $E_{2}$ over a neighborhood $U \subset M$ of $x$. Then $\mathscr{F}\left(U \times S_{1}, U \times S_{2}\right)=U \times C^{\infty}\left(S_{1}, S_{2}\right)$. The restriction of $F$ to $Y^{p}=0$ represents a map $\bar{F}: S_{1} \rightarrow T S_{2}$ along $\varphi$. By Proposition 5 from [16] there exists a smooth curve $\gamma: \mathbb{R} \rightarrow C^{\infty}\left(S_{1}, S_{2}\right)$ such that $\bar{F}(y)=\frac{\partial \bar{\gamma}(y, 0)}{\partial t}$, where $\tilde{\gamma}: \mathbb{R} \times S_{1} \rightarrow S_{2}$ is defined by $\tilde{\gamma}(y, t)=\gamma(t)(y)$. If $\delta: \mathbb{R} \rightarrow U$ is any curve with $\left.\frac{\partial}{\partial t}\right|_{0} \delta=X$, then the curve $(\delta, \gamma): \mathbb{R} \rightarrow U \times C^{\infty}\left(S_{1}, S_{2}\right)$ has the required property.

Now we show that each fiber of $T \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ is a vector space. Consider $\tilde{A}_{1}: T_{X_{1}} E_{1} \rightarrow T_{X_{1}} E_{2}$ and $\tilde{A}_{2}: T_{X_{2}} E_{1} \rightarrow T_{X_{2}} E_{2}$ over the same $\varphi$. Given $Y \in\left(T_{X_{1}+X_{2}} E_{1}\right)_{y}, y \in E_{1 x}$, we take any $W \in\left(T_{X_{1}} E_{1}\right)_{y}$, so that $Y-W \in\left(T_{X_{2}} E_{1}\right)_{y}$, and we define

$$
\widetilde{A_{1}+A_{2}}(Y)=\tilde{A}_{1}(W)+\tilde{A}_{2}(Y-W)
$$

If we select another $\bar{W} \in\left(T_{X_{1}} E_{1}\right)_{y}$, then $W-\bar{W}$ is a vertical vector. Hence

$$
\tilde{A}_{1}(\bar{W})=\tilde{A}_{1}(W)+T \varphi(\bar{W}-W), \quad \tilde{A}_{2}(Y-\bar{W})=\tilde{A}_{2}(Y-W)+T \varphi(W-\bar{W})
$$

so that our definition is correct. Further, for $0 \neq k \in \mathbb{R}$ we define

$$
\widetilde{k A}: T_{k X} E_{1} \rightarrow T_{k X} E_{2} \quad \text { by } \quad \widetilde{k A}(Y)=k \tilde{A}\left(\frac{1}{k} Y\right)
$$

while for $k=0$ we prescribe $\widetilde{0 A}$ to be $T \varphi: T_{0} E_{1 x} \rightarrow T_{0} E_{2 x}$. In coordinates, if $A_{1}=\left(x^{i}, X_{1}^{i}, \varphi^{a}, \Phi_{1}^{a}\right)$ and $A_{2}=\left(x^{i}, X_{2}^{i}, \varphi^{a}, \Phi_{2}^{a}\right)$, then

$$
\begin{equation*}
A_{1}+A_{2}=\left(x^{i}, X_{1}^{i}+X_{2}^{i}, \varphi^{a}, \Phi_{1}^{a}+\Phi_{2}^{a}\right), k A_{1}=\left(x^{j}, k X_{1}^{i}, \varphi^{a}, k \Phi_{1}^{a}\right) \tag{5}
\end{equation*}
$$

This proves that each $\pi^{-1}(\varphi)$ is a vector space.
In general, consider another pair $E_{3} \rightarrow N, E_{4} \rightarrow N$ of fiber bundles over the same base and subset $\varepsilon \subset \mathscr{F}\left(E_{1}, E_{2}\right)$.

Definition 3. A map $f: \varepsilon \rightarrow \mathscr{F}\left(E_{3}, E_{4}\right)$ is called smooth, if $f \circ g: Q \rightarrow$ $\mathscr{F}\left(E_{3}, E_{4}\right)$ is smooth for every smooth map $g: Q \rightarrow \varepsilon$.

Definition 4. A vector field on $\mathscr{F}\left(E_{1}, E_{2}\right)$ is a smooth map $A: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow$ $T \mathscr{F}\left(E_{1}, E_{2}\right)$ satisfying $\pi \circ A=\mathrm{id}$. We say that $A$ is projectable, if there exists a classical smooth vector field $A^{0}: M \rightarrow T M$ such that $A^{0} \circ \pi=T p \circ A$.

Write $V \mathscr{F}\left(E_{1}, E_{2}\right)$ for the kernel of $T p: T \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T M$, which will be called the vertical tangent bundle of $\mathscr{F}\left(E_{1}, E_{2}\right)$. Then we have an exact sequence

$$
\begin{equation*}
0 \rightarrow V \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right) \underset{M}{\times} T M \rightarrow 0 \tag{6}
\end{equation*}
$$

Consider a linear splitting $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \times T M \rightarrow T \mathscr{F}\left(E_{1}, E_{2}\right)$, i.e. $\pi \circ \Gamma=$ $p r_{1}, T p \circ \Gamma=p r_{2}$ and $\Gamma(\varphi,-): T_{x} M \rightarrow M_{\varphi}^{M} \mathscr{F}\left(E_{1}, E_{2}\right)$ is a linear map for each $\varphi \in \mathscr{F}\left(E_{1}, E_{2}\right), x=\pi(\varphi)$. Then for every vector field $X: M \rightarrow T M$ we have defined its $\Gamma$-lift $\Gamma X: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T \mathscr{F}\left(E_{1}, E_{2}\right)$. We say that $\Gamma$ is smooth, if $\Gamma X$ is smooth for every classical smooth vector field $X: M \rightarrow T M$.

Definition 5. A connection (in tangent form) on $\mathscr{F}\left(E_{1}, E_{2}\right)$ is a smooth linear splitting $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \underset{M}{\times T M} \rightarrow T \mathscr{F}\left(E_{1}, E_{2}\right)$.

Remark 2. If $E_{1}$ is the trivial fibering $M \rightarrow M$, then $\mathscr{F}\left(E_{1}, E_{2}\right)=E_{2}$ and we obtain the standard connection on $E_{2} \rightarrow M$.

## 2. Jet prolongations of $\mathscr{F}\left(E_{1}, E_{2}\right)$

The simplest way how to define the $r$-th jet prolongation of $\mathscr{F}\left(E_{1}, E_{2}\right)$ is based on the concept of the fiber $r$-jet, [9], [10]. In general, given a fiber bundle $E \rightarrow M$ and a manifold $N$, two maps $f, g: E \rightarrow N$ are said to determine the same fiber $r$-jet $j_{x}^{r} f=j_{x}^{r} g$ at $x \in M$, if $j_{y}^{r} f=j_{y}^{r} g$ for all $y \in E_{x}$. Every smooth section $s$ of $\mathscr{F}\left(E_{1}, E_{2}\right)$ determines the associated base-preserving morphism $\tilde{s}: E_{1} \rightarrow E_{2}$, $\tilde{s}(y)=s\left(p_{1} y\right)(y)$.

Definition 6. Two sections $s_{1}, s_{2}: M \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ determine the same $r$-jet $j_{x}^{r} s_{1}=j_{x}^{r} s_{2}$ at $x \in M$, if $j_{x}^{r} \tilde{s}_{1}=j_{x}^{r} \tilde{s}_{2}$. The set $J^{r} \mathscr{F}\left(E_{1}, E_{2}\right)$ of all $r$-jets of the local sections of $\mathscr{F}\left(E_{1}, E_{2}\right)$ is called the $r$-jet prolongation of $\mathscr{F}\left(E_{1}, E_{2}\right)$.

However, it will be useful to discuss another approach as well. Since $\tilde{s}: E_{1} \rightarrow E_{2}$ is a base-preserving morphism, we can construct its $r$-th jet prolongation $J^{r} \tilde{s}: J^{r} E_{1} \rightarrow$ $J^{r} E_{2}$. Write $J_{x}^{r} \tilde{s}=J^{r} \tilde{s} \mid J_{x}^{r} E_{1}, x \in M$. By direct evaluation, one easily verifies.

Proposition 1. We have $j_{x}^{r} s_{1}=j_{x}^{r} s_{2}$ iff $J_{x}^{r} \tilde{s}_{1}=J_{x}^{r} \tilde{s}_{2}$.
Let $z^{a}=f^{a}\left(x^{i}, y^{p}\right)$ be the coordinate expression of $\tilde{s}$. Then the additional coordinate expression of $J_{x}^{1} \tilde{s}$ is

$$
\begin{equation*}
z_{i}^{a}=\frac{\partial f^{a}}{\partial x^{i}}+\frac{\partial f^{a}}{\partial y^{p}} y_{i}^{p} \tag{7}
\end{equation*}
$$

where $y_{i}^{p}$ or $z_{i}^{a}$ are the induced coordinates on $J^{1} E_{1}$ or $J^{1} E_{2}$. For $x=0$, the functions $\varphi^{n}\left(y^{p}\right):=f^{a}\left(0, y^{p}\right)$ are the coordinates of the target $s(0)$ of $j_{0}^{1} s_{1}$ and $J_{0}^{1} \tilde{s}$ has the form

$$
\begin{equation*}
z_{i}^{a}=\frac{\partial \varphi^{a}(y)}{\partial y^{p}} y_{i}^{p}+\varphi_{i}^{a}(y), \quad \varphi_{i}^{a}(y)=\frac{\partial f^{a}(0, y)}{\partial x^{i}} \tag{8}
\end{equation*}
$$

It is well-known that $J_{x}^{1} E_{1}$ or $J_{x}^{1} E_{2}$ is an affine bundle over $E_{1 x}$ or $E_{2 x}$, whose derived vector bundle is $V_{x} E_{1} \otimes T_{x}^{*} M$ or $V_{x} E_{2} \otimes T_{x}^{*} M$, respectively. Obviously, (8) is an affine bundle morphism over $\varphi$ with the derived linear morphism $T \varphi \otimes \mathrm{id}_{T_{*}^{*} M}$. Similarly to $\S 1$, we denote by $\widetilde{j_{x}^{1}} s$ the associated map $J_{x}^{1} \tilde{s}: J_{x}^{1} E_{1} \rightarrow J_{x}^{1} E_{2}$. Analogously to Lemma 1, one can prove

Lemma 2. Let $S: J_{x}^{1} E_{1} \rightarrow J_{x}^{1} E_{2}$ be an affine bundle morphism over $\varphi: E_{1 x} \rightarrow$ $E_{2 x}$ with the derived linear morphism $T \underset{\sim}{\varphi} \otimes \mathrm{id}_{T_{x}^{*} M}$. Then there exists a local section $s$ of $\mathscr{F}\left(E_{1}, E_{2}\right)$ such that $s(x)=\varphi$ and $\widetilde{j_{x}^{1}} s=S$.

By (8), every $X=\left.\frac{\partial}{\partial t}\right|_{0} f \in T_{x} M$ and every $S=j_{x}^{1} s$ define a vector

$$
\begin{equation*}
S(X)=\left.\frac{\partial}{\partial t}\right|_{0}(s \circ f) \in T_{s(x)} \mathscr{F}\left(E_{1}, E_{2}\right) \tag{9}
\end{equation*}
$$

such that $T p(S(X))=X$.
Definition 7. A connection in the jet form on $\mathscr{F}\left(E_{1}, E_{2}\right)$ is a smooth section $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ of the target jet projection.

Proposition 2. The map (9) establishes a bijection between the jet form and the tangent form of connections on $\mathscr{F}\left(E_{1}, E_{2}\right)$.

Proof. Using (8) we find directly that (9) defines a bijection between the linear splittings $T_{x} M \rightarrow T_{\varphi} \mathscr{F}\left(E_{1}, E_{2}\right)$ of $T p$ and the elements of $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)_{\varphi}$. Assume the jet form of $\Gamma$ is smooth and $f: Q \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ is a smooth map, so that $\Gamma \circ f$ : $Q \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ is smooth. For every smooth vector field $X: M \rightarrow T M$, the map $(\Gamma \circ f)(X \circ p \circ f)$ is also smooth, so that the tangent form of $\Gamma$ is smooth. Conversely, take a local basis $X_{1}, \ldots, X_{m}$ of vector fields on $T M$. Then $\left(\Gamma X_{1}\right) \circ f, \ldots,\left(\Gamma X_{m}\right) \circ f$ are smooth maps $Q \rightarrow T \mathscr{F}\left(E_{1}, E_{2}\right)$. By (8) we deduce that $\Gamma \circ f: Q \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ is smooth.

To define the curvature of a connection of $\mathscr{F}\left(E_{1}, E_{2}\right)$ in $\S 5$, we shall use the second semiholonomic prolongation of $\mathscr{F}\left(E_{1}, E_{2}\right)$. We recall that $J^{1}\left(J^{1} E_{1} \rightarrow M\right):=\tilde{J}^{2} E_{1}$ is the classical second nonholonomic prolongation of $E_{1} \rightarrow M$. If $x^{i}, y^{p}, y_{i}^{p}$ are the above local coordinates of $J^{1} E_{1}$, then the induced coordinates on $\tilde{J}^{2} E_{1}$ are $y_{0 i}^{p}=\frac{\partial y^{p}}{\partial x^{i}}$ and $y_{i j}^{p}=\frac{\partial y_{i}^{p}}{\partial x^{j}}$. We have the target jet projection $\beta_{1}: \tilde{J}^{2} E_{1} \rightarrow J^{1} E_{1}$ and the induced map $J^{1} \beta: \tilde{J}^{2} E_{1} \rightarrow J^{1} E_{1}$ of the target jet projection $\beta: J^{1} E_{1} \rightarrow E_{1}$. An element $Y \in \tilde{J}^{2} E_{1}$ is said to be semiholonomic if $\beta_{1}(Y)=J^{1} \beta(Y)$. In coordinates this is characterized by $y_{i}^{p}=y_{0 i}^{p}$. All semiholonomic elements form a subbundle $\bar{J}^{2} E_{1} \subset \tilde{J}^{2} E_{1}$, and the second holonomic prolongation $J^{2} E$ is a subbundle of $\bar{J}^{2} E$.

Since we have interpreted $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ as a subset of $\mathscr{F}\left(J^{1} E_{1}, J^{1} E_{2}\right)$, we have defined $j_{x}^{1} \sigma$ for a local smooth section $\sigma$ of $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow M$ by $j_{x}^{1} \tilde{\sigma}$. In this way we
introduce the second nonholonomic prolongation $\tilde{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$ of $\mathscr{F}\left(E_{1}, E_{2}\right)$. An element $j_{x}^{1} \sigma$ is said to be semiholonomic, if $\sigma(x)=j_{x}^{1}(\beta \circ \sigma)$, where $\beta: J^{1} \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow$ $\mathscr{F}\left(E_{1}, E_{2}\right)$ is the target jet projection. This defines $\bar{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right) \subset \tilde{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$. The inclusion $J^{2} \mathscr{F}\left(E_{1}, E_{2}\right) \subset \bar{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$ is given by $j_{x}^{2} s \mapsto j_{x}^{1}\left(j^{1} s\right)$.

Analogously to the first order case, $j_{x}^{1} \sigma$ determines a map $\tilde{j_{x}^{1}} \sigma: \tilde{J}_{x}^{2} E_{1} \rightarrow \tilde{J}_{x}^{2} E_{2}$. In coordinates, if $\sigma=\left(f^{a}(x, y), f_{i}^{a}(x, y)\right)$, then $\tilde{s}$ is of the form

$$
\begin{equation*}
z^{a}=f^{a}(x, y), \quad z_{i}^{a}=\frac{\partial f^{a}(x, y)}{\partial y^{p}} y_{i}^{p}+f_{i}^{a}(x, y) \tag{10}
\end{equation*}
$$

Hence

$$
\begin{equation*}
\varphi^{a}(y)=f^{a}(0, y), \quad \varphi_{i}^{a}=f_{i}^{a}(0, y), \varphi_{0 i}^{a}=\frac{\partial f^{a}(0, y)}{\partial x^{i}}, \quad \varphi_{i j}^{a}=\frac{\partial f_{i}^{a}(0, y)}{\partial x^{j}} \tag{11}
\end{equation*}
$$

are the coordinates of $j_{0}^{1} \sigma$. From (10) we obtain the coordinate expression of $\tilde{j_{x}^{1}} \sigma$ in the form $z^{a}=\varphi^{a}(y)$ and

$$
\begin{gather*}
z_{i}^{a}=\frac{\partial \varphi^{a}}{\partial y^{p}} y_{i}^{p}+\varphi_{i}^{a}, \quad z_{0 i}^{a}=\frac{\partial \varphi^{a}}{\partial y^{p}} y_{0 i}^{p}+\varphi_{0 i}^{a},  \tag{12}\\
z_{i j}^{a}=\varphi_{i j}^{a}+\frac{\partial \varphi_{i}^{a}}{\partial y^{p}} y_{0 j}^{p}+\frac{\partial \varphi_{0 j}^{a}}{\partial y^{p}} y_{i}^{p}+\frac{\partial^{2} \varphi^{a}}{\partial y^{p} \partial y^{q}} y_{i}^{p} y_{0 j}^{q}+\frac{\partial \varphi^{a}}{\partial y^{p}} y_{i j}^{p} .
\end{gather*}
$$

Using (12) we deduce directly the following assertion.
Proposition 3. $j_{x}^{1} \sigma$ is semiholonomic or holonomic iff $\tilde{j_{x}^{1}} \sigma$ maps $\bar{J}_{x}^{2} E_{1}$ into $\bar{J}_{x}^{2} E_{2}$ or $J_{x}^{2} E_{1}$ into $J_{x}^{2} E_{2}$, respectively.

In coordinates, an element of $\bar{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$ is characterized by $\varphi_{i}^{a}=\varphi_{0 i}^{a}$ and the additional condition for a holonomic element is $\varphi_{i j}^{a}=\varphi_{j i}^{a}$.

We remark that the higher order nonholonomic and semiholonomic prolongations of $\mathscr{F}\left(E_{1}, E_{2}\right)$ can be defined in a quite similar way.

## 3. The finite order case

Since both vector fields from $\S 1$ and the connections from $\S 2$ are defined on a functional bundle, they represent a kind of differential operators. We are going to describe the simplest case of finite order operators.

Definition 8. A projectable vector field $A: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T \mathscr{F}\left(E_{1}, E_{2}\right)$ over $A^{0}$ : $M \rightarrow T M$ is of order $r$, if the condition $j_{y}^{r} \varphi=j_{y}^{r} \psi, \varphi, \psi \in C^{\infty}\left(E_{1 x}, E_{2 x}\right), y \in E_{1 x}$ implies that the restrictions of $\widetilde{A(\varphi)}$ and $\widetilde{A(\psi)}$ over $y$ coincide, i.e.

$$
\begin{equation*}
\widetilde{A(\varphi)}\left|\left(T_{A^{0}(x)} E_{1}\right)_{y}=\widetilde{A(\psi)}\right|\left(T_{A^{0}(x)} E_{1}\right)_{y} \tag{13}
\end{equation*}
$$

Let $S\left(T E_{1}, T E_{2}\right)$ be the set of all affine morphism $\left(T_{X} E_{1}\right)_{y} \rightarrow\left(T_{X} E_{2}\right)_{z}, p_{1} y=$ $p_{2} z=\pi_{M} X$, where $\pi_{M}: T M \rightarrow M$ is the bundle projection. This is a fibered manifold over $E_{1} \times E_{M} \times T M$. Write

$$
\mathscr{F} J^{r}\left(E_{1}, E_{2}\right)=\bigcup_{x \in M} J^{r}\left(E_{1 x}, E_{2 x}\right)
$$

This is a classical manifold as well.
A projectable $r$-th order vector field $A: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T \mathscr{F}\left(E_{1}, E_{2}\right)$ over $A^{0}$ defines the associated map $\mathscr{A}: \mathscr{F} J^{r}\left(E_{1}, E_{2}\right) \rightarrow S\left(T E_{1}, T E_{2}\right)$ by

$$
\begin{equation*}
\mathscr{A}\left(j_{y}^{r} \varphi\right)=A(\varphi) \mid\left(T_{A^{0}(x)} E_{1}\right)_{y} \tag{14}
\end{equation*}
$$

Proposition 4. The associated map of a projectable $r$-th order vector field on $\mathscr{F}\left(E_{1}, E_{2}\right)$ is a classical $C^{\infty}$-map.

Proof. This follows from the fact that $A$ is smooth in the sense of Definition 3 quite analogously to [6].

The local coordinates on $\mathscr{F} J^{r}\left(E_{1}, E_{2}\right)$ induced by $x^{i}, y^{p}$ and $z^{a}$ are $z_{\alpha}^{a}, 1 \leqslant|\alpha| \leqslant r$, where $\alpha$ is a multiindex, the range of which is the fiber dimension of $E_{1}$. Hence the coordinate form of $\mathscr{A}$ is $X^{i}\left(x^{j}\right)$ and

$$
\begin{equation*}
\Phi^{a}=\Phi^{a}\left(x^{i}, y^{p}, z_{\alpha}^{a}\right), \quad 0 \leqslant|\alpha| \leqslant r \tag{15}
\end{equation*}
$$

The derived linear map of each element of $S\left(T E_{1}, T E_{2}\right)$ is identified with an element of $\mathscr{F} J^{1}\left(E_{1}, E_{2}\right)$. This defines a map $D: S\left(T E_{1}, T E_{2}\right) \rightarrow \mathscr{F} J^{1}\left(E_{1}, E_{2}\right)$ and the following diagram commutes:

where $\beta_{r}$ is the jet projection. Conversely, let $\mathscr{A}: \mathscr{F} J^{r}\left(E_{1}, E_{2}\right) \rightarrow S\left(T E_{1}, T E_{2}\right)$ be a smooth map with an underlying vector field $A^{0}: M \rightarrow T M$ such that (16) commutes. Then the rule

$$
\begin{equation*}
\widetilde{A(\varphi)}=\bigcup_{y \in E_{1 x}} \mathscr{A}\left(j_{y}^{r} \varphi\right) \tag{17}
\end{equation*}
$$

defines a projectable $r$-th order vector field $A$ on $\mathscr{F}\left(E_{1}, E_{2}\right)$.
Since $T \mathscr{F}\left(E_{1}, E_{2}\right)$ is a subset of $\mathscr{F}\left(T E_{1}, T E_{2}\right)$, we can define the second tangent bundle $T\left(T \mathscr{F}\left(E_{1}, E_{2}\right)\right)$. This will be described in more detail in $\S 6$. Here we restrict ourselves to a general remark, which is related to our study of the order of connections.

Definition 9. A vector field $A: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ is called differentiable if the formula

$$
\begin{equation*}
T A\left(\left.\frac{\partial}{\partial t}\right|_{0} f\right)=\left.\frac{\partial}{\partial t}\right|_{0} A \circ f \tag{18}
\end{equation*}
$$

defines a smooth map $T A: T \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T T \mathscr{F}\left(E_{1}, E_{2}\right)$.
From (16) we easily deduce (see the coordinate formula in §6) the following assertion.

Proposition 5. Every $r$-th order vector field on $\mathscr{F}\left(E_{1}, E_{2}\right)$ is differentiable.
Definition 10. A connection $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ is of order $r$ if the condition $j_{y}^{r} \varphi=j_{y}^{r} \psi, \varphi, \psi \in C^{\infty}\left(E_{1 x}, E_{2 x}\right), y \in E_{1 x}$, implies

$$
\begin{equation*}
\widetilde{\Gamma(\varphi)}\left|J_{y}^{1} E_{1}=\widetilde{\Gamma(\Psi)}\right| J_{y}^{1} E_{1} . \tag{19}
\end{equation*}
$$

Let $S\left(J^{1} E_{1}, J^{1} E_{2}\right)$ be the set of all affine maps $\left(J^{1} E_{1}\right)_{y} \rightarrow\left(J^{1} E_{2}\right)_{z}$ with the derived linear map of the form

$$
\begin{equation*}
B \otimes \mathrm{id}_{T_{x}^{*} M} \quad B \in \mathscr{L} \in\left(V_{y} E_{1}, V_{z} E_{2}\right) . \tag{20}
\end{equation*}
$$

An $r$-th order connection $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ defines the associated map $\mathscr{G}: \mathscr{F} J^{r}\left(E_{1}, E_{2}\right) \rightarrow S\left(J^{1} E_{1}, J^{1} E_{2}\right)$ by

$$
\begin{equation*}
\mathscr{G}\left(j_{y}^{r} \varphi\right)=\widetilde{\Gamma(\varphi)} \mid J_{y}^{1} E_{1} . \tag{21}
\end{equation*}
$$

The coordinate form of $\mathscr{G}$ is

$$
\begin{equation*}
\Phi_{i}^{a}=\Phi_{i}^{a}\left(x^{i}, y^{p}, z_{\alpha}^{a}\right), \quad 0 \leqslant|\alpha| \leqslant r . \tag{22}
\end{equation*}
$$

Analogously to Proposition 4, one proves

Proposition 6. The associated map of an $r$-th order connection $\mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow$ $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ is a classical $C^{\infty}$-map.

Let $D: S\left(J^{1} E_{1}, J^{1} E_{2}\right) \rightarrow \mathscr{F} J^{1}\left(E_{1}, E_{2}\right)$ be the map defined by (20). Then the following diagram commutes


Conversely, let $\mathscr{G}: \mathscr{F} J^{r}\left(E_{1}, E_{2}\right) \rightarrow S\left(J^{1} E_{1}, J^{1} E_{2}\right)$ be a smooth morphism over the identity of $E_{1} \times E_{2}$ such that (23) commutes. Then the rule

$$
\begin{equation*}
\widetilde{\Gamma(\varphi)}=\bigcup_{y \in E_{1 x}} \mathscr{G}\left(j_{y}^{r} \varphi\right) \tag{24}
\end{equation*}
$$

defines an $r$-th order connection on $\mathscr{F}\left(E_{1}, E_{2}\right)$.
Analogously to Definition 9, we introduce
Definition 11. A connection $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ is called differentiable if the formula

$$
\begin{equation*}
J^{1} \Gamma\left(j_{x}^{1} s\right)=j_{x}^{1}(\Gamma \circ s) \tag{25}
\end{equation*}
$$

defines a smooth map $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \tilde{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$.
Proposition 7. Every $r$-th order connection is differentiable.
Proof. We deduce from (22) the coordinate form of $J^{1} \Gamma$ in some coordinates $x^{i}, \varphi^{a}, \psi_{i}^{a}$ on $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ and $x^{i}, \varphi^{a}, \varphi_{i}^{a}, \varphi_{0 i}^{a}, \varphi_{i j}^{a}$ on $\tilde{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$. Take a section $\sigma$

$$
\begin{equation*}
z^{a}=\Psi^{a}\left(x^{i}, y^{p}\right) \tag{26}
\end{equation*}
$$

so that $\varphi^{a}=\psi^{a}(0, y)$ and $\psi_{i}^{a}=\frac{\partial \psi^{a}(0, y)}{\partial x^{i}}$. Then we obtain for $\Gamma \circ \sigma$

$$
\begin{equation*}
z_{i}^{a}=\frac{\partial \psi^{a}(x, y)}{\partial y^{p}} y_{i}^{p}+\Phi_{i}^{a}\left(x, y, \partial_{\alpha} \psi^{a}(x, y)\right) \tag{27}
\end{equation*}
$$

Now (26) yields

$$
\begin{equation*}
z_{0 i}^{a}=\frac{\partial \psi^{a}(0, y)}{\partial y^{p}} y_{0 i}^{p}+\frac{\partial \psi^{a}(0, y)}{\partial x^{i}}, \quad \text { i.e. } \quad \varphi_{0 i}^{a}=\psi_{i}^{a} \tag{28}
\end{equation*}
$$

and (27) implies

$$
\begin{align*}
z_{i j}^{a}= & \frac{\partial \psi_{j}^{a}}{\partial y^{p}} y_{i}^{p}+\frac{\partial^{2} \varphi^{a}}{\partial y^{p} \partial y^{q}} y_{i}^{p} y_{0 j}^{q}+\frac{\partial \varphi^{a}}{\partial y^{p}} y_{i j}^{p}+\frac{\partial \varphi_{i}^{a}}{\partial y^{p}} y_{0 j}^{p} \\
& +\frac{\partial \Phi_{i}^{a}}{\partial x^{j}}+\frac{\partial \Phi_{i}^{a}}{\partial z^{b}} \partial_{j} \psi^{b}+\ldots+\frac{\partial \Phi_{i}^{a}}{\partial z_{\alpha}^{b}} \partial_{\alpha} \partial_{j} \psi^{b} \tag{29}
\end{align*}
$$

In particular, (29) shows that $J^{1} \Gamma$ is well-defined and smooth.

Following Virsik, [17], if $\Gamma$ is differentiable and $\Delta$ is another connection $\mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$, we define a section

$$
\begin{equation*}
\Gamma * \Delta=J^{1} \Gamma \circ \Delta: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \tilde{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right) \tag{30}
\end{equation*}
$$

The order of such a section can be introduced similarly to Definition 10.

Proposition 8. If $\Gamma$ and $\Delta$ are connections of orders $r$ and $s$, respectively, then $\Gamma * \Delta$ has the order $r+s$.

Proof. We substitute the associated map of $\Delta$ into (28) and (29).
To obtain an explicit formula for the associated map of $\Gamma * \Delta$, we introduce the following concept. Having a smooth function $f: \mathscr{F} J^{r}\left(E_{1}, E_{2}\right) \rightarrow \mathbb{R}$, we define its formal differential $D f$ by

$$
\begin{equation*}
D f: \mathscr{F} J^{r+1}\left(E_{1}, E_{2}\right) \rightarrow V^{*} E_{1}, D f\left(j_{y}^{r+1} \varphi\right)=d_{y} f\left(j^{r} \varphi\right) \tag{31}
\end{equation*}
$$

Then every vertical vector field $\mu$ on $V^{*} E_{1}$ determines $\langle D f, \mu\rangle: \mathscr{F} J^{r+1}\left(E_{1}, E_{2}\right) \rightarrow \mathbb{R}$. For the coordinate vector fields $\frac{\partial}{\partial y^{n}}$ we obtain the formal derivatives

$$
\begin{equation*}
D_{p} f=\frac{\partial f}{\partial y^{p}}+\frac{\partial f}{\partial z^{a}} z_{p}^{a}+\ldots+\frac{\partial f}{\partial z_{\alpha}^{a}} z_{\alpha+p}^{a} \tag{32}
\end{equation*}
$$

By itcration, we introduce $D_{\beta} f: \mathscr{F} J^{r+|\beta|}\left(E_{1}, E_{2}\right) \rightarrow \mathbb{R}$. Let $\Psi_{i}^{a}\left(x^{i}, y^{p}, z_{\beta}^{a}\right), 0 \leqslant|\beta| \leqslant$ $s$, be associated map of $\Delta$. Then the coordinate form of the main term of (29) is

$$
\begin{equation*}
\varphi_{i j}^{a}=\frac{\partial \Phi_{i}^{a}}{\partial x^{j}}+\frac{\partial \Phi_{i}^{a}}{\partial z^{b}} \Psi_{j}^{b}+\frac{\partial \Phi_{i}^{a}}{\partial z_{p}^{b}} D_{p} \Psi_{j}^{b}+\ldots+\frac{\partial \Phi_{i}^{a}}{\partial z_{\alpha}^{b}} D_{\alpha} \Psi_{j}^{b} . \tag{33}
\end{equation*}
$$

Remark 3. In both cases of connections in the jet form and of projectable vector fields we have a situation somewhat similar to the vertical prolongation operators on classical fibered manifolds studied by Kosmann-Schwarzbach, [11], and the second author, [8]. In [10] Slovák deduced that every vertical prolongation operator is differentiable in the sense of our Definitions 9 and 11. However, his proof is based on quite sophisticated procedures in mathematical analysis, so that we have the feeling that such a problem in our setting is beyond the scope of the present paper.

## 4. Ehresmann prolongation in the classical case

We describe some properties of connections on a classical fibered manifold $p$ : $E \rightarrow M$ in a way which can be generalized to $\mathscr{F}\left(E_{1}, E_{2}\right)$. Given $A \in J_{y}^{1} E$ and $B \in T_{x} M, x=p y$, we denote by $A(B) \in T_{y} E$ the $A$-lift of $B$. We show that every $A \in \tilde{J}_{y}^{2} E$ induces similarly a lifting $\lambda A: T T_{x} M \rightarrow T T_{y} E$. If $A=J_{x}^{1} \sigma$ and $B=\left.\frac{\partial}{\partial t}\right|_{0} f(t) \in T T_{x} M$, then we construct $\sigma(\pi(f(t)))(f(t)): \mathbb{R} \rightarrow T E$ and set

$$
\begin{equation*}
\lambda A(B)=\left.\frac{\partial}{\partial t}\right|_{0} \sigma(\pi(f(t)))(f(t)) \tag{34}
\end{equation*}
$$

where $\pi: T M \rightarrow M$ is the bundle projection. Given some local fiber coordinates $x^{i}$, $y^{p}$ on $E$, we have the induced coordinates $y_{i}^{p}, y_{0 i}^{p}, y_{i j}^{p}$ on $\tilde{J}^{2} E$, the induced coordinates $X^{i}, Y^{p}$ on $T E$ and the additional coordinates on TEE denoted by a dot. Then one finds easily the following coordinate form of (34):

$$
\begin{equation*}
Y^{p}=y_{i}^{p} X^{i}, \quad \dot{y}^{p}=y_{0 i}^{p} \dot{x}^{i}, \quad \dot{Y}^{p}=y_{i j}^{p} X^{i} \dot{x}^{j}+y_{i}^{p} \dot{X}^{i} . \tag{35}
\end{equation*}
$$

Let $\kappa$ be the canonical involution of the second tangent bundle. If $A \in \bar{J}_{y}^{2} E$, then $\kappa_{E} \circ \lambda A \circ \kappa_{M}: T T_{x} M \rightarrow T T_{y} E$ is the lifting of another element $\kappa A \in \bar{J}_{y}^{2} E$, [15]. In coordinates, $y_{j i}^{p}(\kappa A)=y_{i j}^{p}(A)$. Hence $A$ is holonomic iff $\kappa A=A$. Since $\bar{J}_{y}^{2} E \rightarrow J^{1} E$ is an affine bundle with the derived vector bundle $V E \otimes\left(\otimes^{2} T^{*} M\right)$, the points $\kappa A$ and $A$ determine a vector $\Delta(A):=\overrightarrow{(\kappa A) A} \in V_{y} E \otimes \Lambda^{2} T_{x}^{*} M$, which is called the deviation (or difference tensor) of $A,[7]$, [12]. The coordinates of $\Delta(A)$ are $y_{i j}^{p}-y_{j i}^{p}$. If $X_{1}, X_{2} \in T_{x} M$, then we have $\Delta(A)\left(X_{1}, X_{2}\right) \in V_{y} E$.

Let $\pi_{1}=\pi_{T M}=T T M \rightarrow T M$ and $\pi_{2}=T \pi_{T M}=T T M \rightarrow T M$ be the canonical projections. Consider $C, D \in T T_{x} M$ satisfying

$$
\begin{equation*}
\pi_{1}(C)=\pi_{2}(D) \quad \text { and } \quad \pi_{1}(D)=\pi_{2}(C) \tag{36}
\end{equation*}
$$

Since $\kappa$ exchanges the two projections, $C$ and $\kappa D$ are in the same fiber of $T T M$ with respect to $\pi_{1}$ and satisfy $\pi_{2}(C-\kappa D)=0$. Hence $C-\kappa D$ is a tangent vector to a fiber of $T M$ and such a vector can be identified with an element of $T_{x} M$, which will be denoted by $C-D$ and called the strong difference of $C$ and $D$. In coordinates, if

$$
\begin{equation*}
C \equiv\left(a^{i}, b^{i}, c^{i}\right), D \equiv\left(b^{i}, a^{i}, d^{i}\right) \quad \text { then } \quad C \doteq D \equiv\left(c^{i}-d^{i}\right) \tag{37}
\end{equation*}
$$

In [8] it is deduced the the bracket [ $X, Y$ ] of two vector fields $X, Y: M \rightarrow T M$ can be expressed by

$$
\begin{equation*}
[X, Y]=T Y \circ X \perp T X \circ Y \tag{38}
\end{equation*}
$$

Lemma 3. Let $C, D \in T T_{x} M$ satisfy the condition (36) for the strong difference and $A \in \bar{J}_{y}^{2} E$. Then $\lambda A(C), \lambda A(D)$ also satisfy (36) and

$$
\Delta A\left(\pi_{1} C, \pi_{2} C\right)=(\lambda A(C)-\lambda A(D))-\beta_{1}(A)(C \doteq D)
$$

where $\beta_{1}: \bar{J}_{1}^{2} E \rightarrow J^{1} E$ is the jet projection.
Proof. By (35) and (37) we have $\lambda A(C)=\left(y_{i}^{p} a^{i}, y_{i}^{p} b^{i}, y_{i j}^{p} a^{i} b^{j}+y_{i}^{p} c^{i}\right), \lambda A(D)=$ $\left(y_{i}^{p} b^{i}, y_{i}^{p} a^{i}, y_{i j}^{p}{ }^{i} a^{j}+y_{i}^{p} d^{i}\right)$. This implies our claim.

According to Remark 2, two connections $\Gamma, \Delta: E \rightarrow J^{1} E$ determine $\Gamma * \Delta=$ $J^{1} \Gamma \circ \Delta: E \rightarrow \tilde{J}^{2} E$. For $\Gamma=\Delta$ the values of $\Gamma * \Gamma$ lie in $\bar{J}_{y}^{2} E$. In this case we obtain a construction closely related to an idea by Ehresmann, [2].

Definition 12. The map $\tilde{\Gamma}=J^{1} \Gamma \circ \Gamma: E \rightarrow \bar{J}^{2} E$ is the Ehresmann prolongation of $\Gamma$. The composition

$$
\begin{equation*}
C \Gamma:=-\Delta \circ \tilde{\Gamma}: E \rightarrow V E \otimes \Lambda^{2} T^{*} M \tag{39}
\end{equation*}
$$

is the curvature of $\Gamma$.
To deduce that $C \Gamma$ coincides with the standard curvature of $\Gamma$, we need a property of the lifting map

$$
\lambda \tilde{\Gamma}: \underset{M}{E \times T M} T T T E .
$$

Consider two vector fields $X, Y: M \rightarrow T M$, so that $T X \circ Y: M \rightarrow T T M$.

Lemma 4. We have

$$
\lambda \tilde{\Gamma}(T X \circ Y)=(T \Gamma X) \circ \Gamma Y: E \rightarrow T T E
$$

Proof. We have $\tilde{\Gamma}(y)=j_{x}^{1}(\Gamma \circ s), j_{x}^{1} s=\Gamma(y)$. If $Y(x)=\left.\frac{\partial}{\partial t}\right|_{0} f(t)$, then

$$
T X(Y(x))=\left.\frac{\partial}{\partial t}\right|_{0}(X \circ f)
$$

By (34),

$$
\lambda \tilde{\Gamma}(T X(Y(x)))=\left.\frac{\partial}{\partial t}\right|_{0} \Gamma(s(f(t)))(X(f(t)))=(T \Gamma X \circ \Gamma Y)(y)
$$

Proposition 9. For every vector fields $X, Y: M \rightarrow T M$, we have

$$
C \Gamma(X, Y)=[\Gamma X, \Gamma Y]-\Gamma([X, Y])
$$

Proof. Consider $T X \circ Y, T Y \circ X: M \rightarrow T T M$. By Lemma 4 we obtain

$$
\lambda \tilde{\Gamma}(T X \circ Y)=T \Gamma X \circ \Gamma Y \quad \text { and } \quad \lambda \tilde{\Gamma}(T Y \circ X)=T \Gamma Y \circ \Gamma X
$$

Then Lemma 3 and (38) imply

$$
\begin{aligned}
\Delta \circ \tilde{\Gamma}(X, Y) & =(\lambda \tilde{\Gamma}(T X \circ Y) \dot{-} \lambda \tilde{\Gamma}(T Y \circ X))-\Gamma(T X \circ Y \doteq T Y \circ X)= \\
& =-[\Gamma X, \Gamma Y]+\Gamma([X, Y])
\end{aligned}
$$

## 5. The curvature of a connection on $\mathscr{F}\left(E_{1}, E_{2}\right)$

The deviation of an element $j_{x}^{1} \sigma \in \tilde{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$ can be defined by means of the associated map $\tilde{j_{x}^{1}} \sigma: \bar{J}_{x}^{2} E_{1} \rightarrow \bar{J}_{x}^{2} E_{2}$. In the semiholonomic case we have $\varphi_{i}^{a}=\varphi_{0 i}^{a}$. So if we take a holonomic 2-jet $Y \in J_{x}^{2} E_{1}$, then the right-hand side of the second line in (12) is symmetric except the first term. Hence the deviation $\Delta\left(\widetilde{j_{x}^{1}} \sigma(Y)\right)$ is independent of $y_{i}^{p}$ and $y_{i j}^{p}$. This defines a map $\Delta\left(j_{x}^{1} \sigma\right): E_{1 x} \rightarrow V_{x} E_{2} \odot \Lambda^{2} T_{x}^{*} M$ over $\varphi$, i.e. an element of $\mathscr{F}\left(E_{1}, V E_{2} \otimes \Lambda^{2} T^{*} M\right)$.

Definition 13. $\Delta\left(j_{x}^{1} \sigma\right)$ is called the deviation of $j_{x}^{1} \sigma$. The coordinate form of $\Delta\left(j_{x}^{1} \sigma\right)$ is $\varphi_{i j}^{a}-\varphi_{j i}^{a}$.

Definition 14. For a differentiable connection $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$. the map $\tilde{\Gamma}:=J^{1} \Gamma \circ \Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \bar{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$ is the Ehresmann prolongation of $\Gamma$.

Definition 15. The composition

$$
C \Gamma:=-\Delta \circ \tilde{\Gamma}: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \mathscr{F}\left(E_{1}, V E_{2} \otimes \Lambda^{2} T^{*} M\right)
$$

is the curvature of a differentiable connection $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$.
Clearly, $C \Gamma$ is a section of the canonical projection $\mathscr{F}\left(E_{1}, V E_{2} \otimes \Lambda^{2} T^{*} M\right) \rightarrow$ $\mathscr{F}\left(E_{1}, E_{2}\right)$.

Let $\Gamma$ be an $r$-th order connection with the associated map $\Phi_{i}^{a}\left(x^{i}, y^{p}, z_{\alpha}^{a}\right)$. Then we obtain the associated map of $C \Gamma$ by setting $\Psi_{i}^{a}=\Phi_{i}^{a}$ in (33) and by antisymmetrizing in $i$ and $j$. This implies

Proposition 10. The curvature of an $r$-th order connection has the order $2 r$.

## 6. The bracket formula for curvature

As remarked in $\S 3$, the inclusion $T \mathscr{F}\left(E_{1}, E_{2}\right) \subset \mathscr{F}\left(T E_{1} \rightarrow T M, T E_{2} \rightarrow T M\right)$ defines the second tangent bundle $T\left(T \mathscr{F}\left(E_{1}, E_{2}\right)\right)=T T \mathscr{F}\left(E_{1}, E_{2}\right)$. We have a projection $T T p: T T \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T T M$ and two projections $\pi_{T}, T \pi: T T \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow$ $T \mathscr{F}\left(E_{1}, E_{2}\right)$. In the above coordinates, consider an element $F \in T T \mathscr{F}\left(E_{1}, E_{2}\right)$ tangent to a curve $x^{j}(t), X^{i}(t), f^{a}(y, t)$ and

$$
Z^{\dot{a}}=\frac{\partial f^{a}(y, t)}{\partial y^{p}} Y^{p}+\Phi^{a}(y, t)
$$

Then its associated map $\tilde{F}: T T_{x} E_{1} \rightarrow T T_{x} E_{2}, X=T T p(F)$, is of the form

$$
\begin{gather*}
Z^{a}=\frac{\partial \varphi^{a}}{\partial y^{p}} Y^{p}+\Phi^{a}(y), \dot{z}^{a}=\frac{\partial \varphi^{a}}{\partial y^{p}} \dot{y}^{p}+f^{a}(y)  \tag{40}\\
\dot{Z}^{a}=F^{a}(y)+\frac{\partial \Phi^{a}}{\partial y^{p}} \dot{y}^{p}+\frac{\partial f^{a}}{\partial y^{p}} Y^{p}+\frac{\partial^{2} \varphi^{a}}{\partial y^{p} \partial y^{q}} Y^{p} \dot{y}^{q}+\frac{\partial \varphi^{a}}{\partial y^{p}} \dot{Y}^{p} .
\end{gather*}
$$

So $\varphi^{a}, \Phi^{a}, f^{a}, F^{a}$ are the functional coordinates of $F$, which are completed by the coordinates $x^{i}, X^{i}, \dot{x}^{i}, \dot{X}^{i}$ of $X \in T T M$. The coordinate form of $\pi_{T}$ or $T \pi$ is

$$
\begin{aligned}
& \pi_{T}\left(x^{i}, X^{i}, \dot{x}^{i}, \dot{X}^{i}, \varphi^{a}, \Phi^{a}, f^{a}, F^{a}\right)=\left(x^{i}, X^{i}, \varphi^{a}, \Phi^{a}\right) \\
& T \pi\left(x^{i}, X^{i}, \dot{x}^{i}, \dot{X}^{i}, \varphi^{a}, \Phi^{a}, f^{a}, F^{a}\right)=\left(x^{i}, \dot{x}^{i}, \varphi^{a}, f^{a}\right)
\end{aligned}
$$

Consider the canonical involution $\kappa_{E_{1}}$ or $\kappa_{E_{1}}$ of the second tangent bundle.
Proposition 11. For every $F \in T T \mathscr{F}\left(E_{1}, E_{2}\right)$ over $X \in T T M$ there exists a unique element $\kappa F \in T T \mathscr{F}\left(E_{1}, E_{2}\right)$ such that its associated map $\widetilde{\kappa F}: T T_{\kappa_{M} X} E_{1} \rightarrow$ $T T_{\kappa_{M} X} E_{2}$ is $\widetilde{\kappa F}=\kappa_{E_{2}} \circ \tilde{F} \circ \kappa_{E_{1}}$.

Proof. This follows from (40).
Obviously, the coordinate form of $\kappa$ is

$$
\begin{equation*}
\kappa(x, X, \dot{x}, \dot{X}, \varphi, \Phi, f, F)=(x, \dot{x}, X, \dot{X}, \varphi, f, \Phi, F) \tag{41}
\end{equation*}
$$

Consider $C, \bar{C} \in T T \mathscr{F}\left(E_{1}, E_{2}\right)$ over $X, \bar{X} \in T T M$ satisfying

$$
\begin{equation*}
\pi_{T}(C)=T \pi(\bar{C}) \quad \text { and } \quad \pi_{T}(\bar{C})=T \pi(C) \tag{42}
\end{equation*}
$$

Then we define the strong difference $C \doteq \bar{C} \in T \mathscr{F}\left(E_{1}, E_{2}\right), T p(C \perp \bar{C})=X \doteq \bar{X}$, as follows. For every $B \in\left(T_{X}-\bar{X} E_{1}\right)_{y}$ we take any $Y, \bar{Y} \in\left(T T E_{1}\right)_{y}$ over $X, \bar{X}$
such that $Y \dot{-}=B$. Then one easily verifies that $C(Y), \bar{C}(\bar{Y})$ also satisfy (42), $C(Y) \doteq \bar{C}(\bar{Y})$ depends on $C, \bar{C}$ and $B$ only and represents the associated map of an element $C-\bar{C} \in T \mathscr{F}\left(E_{1}, E_{2}\right)$, whose coordinates are

$$
\begin{equation*}
\left(x^{i}, \dot{X}^{i}-\dot{\bar{X}}^{i}, \varphi^{a}, F^{a}-\bar{F}^{a}\right) . \tag{43}
\end{equation*}
$$

Let $A, B$ be two differentiable vector fields on $\mathscr{F}\left(E_{1}, E_{2}\right)$. Then the maps $T A$ 。 $B, T B \circ A: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T T \mathscr{F}\left(E_{1}, E_{2}\right)$ satisfy the condition (42) at every $\varphi \in$ $\mathscr{F}\left(E_{1}, E_{2}\right)$.

Definition 16. The vector field

$$
[A, B]:=T B \circ A-T A \circ B: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow T \mathscr{F}\left(E_{1}, E_{2}\right)
$$

is called the bracket of $A$ and $B$.
By (38) we immediately deduce
Proposition 12. If $A$ and $B$ are projectable over $A^{0}$ and $B^{0}$, then $[A, B]$ is projectable over $\left[A^{0}, B^{0}\right]$.

Assume $A$ is of order $r$ and $B$ is of order $s$ with the associated maps $X^{i}(x)$, $A^{a}\left(x^{i}, y^{p}, z_{\alpha}^{a}\right),|\alpha| \leqslant r$ and $Y^{i}(x), B^{a}\left(x^{i}, y^{p}, z_{\beta}^{a}\right),|\beta| \leqslant s$, respectively. Analogously to $\S 3$, the fourth component of the associated map of $T A \circ B$ is

$$
\begin{equation*}
\frac{\partial A^{a}}{\partial x^{i}} Y^{i}+\frac{\partial A^{a}}{\partial z^{b}} B^{b}+\frac{\partial A^{a}}{\partial z_{p}^{b}} D_{p} B^{b}+\ldots+\frac{\partial A^{a}}{\partial Z_{\alpha}^{b}} D_{\alpha} B^{b}, \quad|\alpha| \leqslant r \tag{44}
\end{equation*}
$$

while the fourth component of the associated map of $T B \circ A$ is

$$
\begin{equation*}
\frac{\partial B^{a}}{\partial x^{i}} X^{i}+\frac{\partial B^{a}}{\partial z^{b}} A^{b}+\frac{\partial B^{a}}{\partial z_{p}^{b}} D_{p} A^{b}+\ldots+\frac{\partial B^{a}}{\partial z_{\beta}^{b}} D_{\beta} A^{b}, \quad|\beta| \leqslant s . \tag{45}
\end{equation*}
$$

Hence we can summarize by
Proposition 13. The bracket $[A, B]$ has the order $r+s$ and its associated map is $\left[A^{0}, B^{0}\right]$ and the difference (45)-(44).

We are going to generalize Proposition 9 to connections on $\mathscr{F}\left(E_{1}, E_{2}\right)$. First of all we remark that every $A=j_{x}^{1} \sigma \in \tilde{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)_{\varphi}$ defines a lifting $\lambda A: T T_{x} M \rightarrow$ $T T_{\varphi} \mathscr{F}\left(E_{1}, E_{2}\right)$ by

$$
\lambda A\left(\left.\frac{\partial}{\partial t}\right|_{0} f\right)=\left.\frac{\partial}{\partial t}\right|_{0} \sigma\left(\pi_{M}(f(t))(f(t)) .\right.
$$

In coordinates, if $A=\left(x^{i}, \varphi^{a}, \varphi_{i}^{a}, \varphi_{0 i}^{a}, \varphi_{i j}^{a}\right)$ and $B=\left.\frac{\partial}{\partial t}\right|_{0} f=\left(x^{i}, X^{i}, \dot{x}^{i}, \dot{X}^{i}\right)$, then one easily finds the following coordinate form of $\lambda A(B)$ :

$$
\begin{equation*}
\left(x^{i}, \varphi^{a}, \varphi_{i}^{a} X^{i}, \varphi_{0 i}^{a} \dot{x}^{i}, \varphi_{i j}^{a} X^{i} \dot{x}^{j}+\varphi_{i}^{a} \dot{X}^{i}\right) \tag{46}
\end{equation*}
$$

This directly implies the following generalization of Lemma 3.
Lemma 5. Let $C, D \in T T_{x} M$ satisfy the condition (36) for the strong difference and $A \in \bar{J}^{2} \mathscr{F}\left(E_{1}, E_{2}\right)$. Then $\lambda A(C), \lambda A(D)$ satisfy (42) and

$$
\Delta A\left(\pi_{T} C, T \pi C\right)=(\lambda A(C)-\lambda A(D))-\beta_{1}(A)(C \perp D)
$$

Now we need an assumption of technical character (which is fulfilled for every finite order connection).

Definition 17. A differentiable connection $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ is called strongly differentiable, if $\Gamma X$ is a differentiable vector field on $\mathscr{F}\left(E_{1}, E_{2}\right)$ for every smooth vector field $X: M \rightarrow T M$.

Proposition 14. For every strongly differentiable connection $\Gamma$ on $\mathscr{F}\left(E_{1}, E_{2}\right)$ and for all vector fields $X, Y$ on $M$ we have

$$
C \Gamma(X, Y)=[\Gamma X, \Gamma Y]-\Gamma([X, Y])
$$

Proof. In the same way as in Lemma 4 we deduce $\lambda \tilde{\Gamma}(T X \circ Y)=(T \Gamma X) \circ \Gamma Y$. Then we apply Lemma 5.

## 7. The absolute differentiation

Let $A, B \in J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)_{\varphi}$ be two 1-jets with the same target $\varphi$. To deduce that their difference is an element $A-B \in \mathscr{F}\left(E_{1}, V E_{2} \otimes T^{*} M\right)$ over $\varphi$, we consider the associated maps $\tilde{A}, \tilde{B}: J_{x}^{1} E_{1} \rightarrow J_{x}^{1} E_{2}$,

$$
A \equiv z_{i}^{a}=\frac{\partial \varphi^{a}}{\partial y^{p}} y_{i}^{p}+\Phi_{i}^{p}(y), \quad B \equiv z_{i}^{a}=\frac{\partial \varphi^{a}(x, y)}{\partial y^{p}} y_{i}^{p}+\Psi_{i}^{p}(y)
$$

The element $A(Y)-B(Y)$ is independent of the choice of $Y \in J_{x}^{1} E_{1}$, which defines a map $E_{1 x} \rightarrow V E_{2} \otimes T^{*} M$ over $\varphi$. (In this sense $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ is an affine bundle with the derived vector bundle $\mathscr{F}\left(E_{1}, V E_{2} \otimes T^{*} M\right)$ analogously to the classical case.)

Let $s: M \rightarrow \mathscr{F}\left(E_{1}, E_{2}\right)$ be a section and $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ a connection.

Definition 18. The absolute differential

$$
\nabla s: M \rightarrow \mathscr{F}\left(E_{1}, V E_{2} \otimes T^{*} M\right)
$$

is the above difference $\nabla s(x)=j_{x}^{1} s-\Gamma(s(x))$.
If $X: M \rightarrow T M$ is a vector field, we define the absolute derivative of $s$ with respect to $X$ by

$$
\begin{equation*}
\nabla_{X} s=\langle\nabla s, X\rangle: M \rightarrow \mathscr{F}\left(E_{1}, V E_{2}\right) \tag{47}
\end{equation*}
$$

where $\langle$,$\rangle is the extension of the evaluation map T \times T^{*} \rightarrow \mathbb{R}$. Having an $r$-th order connection with the associated map (22) and a section $s$ of the form $z^{a}=\varphi^{a}(x, y)$, then the coordinate form of $\nabla s$ is

$$
\begin{equation*}
\frac{\partial \varphi^{a}(x, y)}{\partial x^{i}}-\Phi_{i}^{a}\left(x^{i}, y^{p}, \partial_{\alpha} \varphi^{a}(x, y)\right) \tag{48}
\end{equation*}
$$

To obtain $\nabla_{x} s$, we contract (48) with the coordinate functions $X^{i}(x)$ of $X$.
Remark 4. In the case $E_{1}=E_{2}:=E$ we have a distinguished section $I$ : $M \rightarrow \mathscr{F}(E, E), I(x)=\mathrm{id}_{E_{x}}$. Analogously to the case of a classical linear connection on $T M$, the absolute differential $\nabla I: M \rightarrow \mathscr{F}\left(E, V E \otimes T^{*} M\right)$ can be called the torsion of a connection $\Gamma$ on $\mathscr{F}(E, E)$. By (48), the coordinate form of the torsion of an $r$-th order connection is $-\Phi_{i}^{p}\left(x^{i}, y^{p}, y^{p}, \delta_{q}^{p}, 0, \ldots, 0\right)$.

It might be instructive to discuss a special case in more detail. Let $E \rightarrow M$ be a vector bundle. Consider the subspace $L E \subset \mathscr{F}(E, E)$ of all linear maps, which is a classical vector bundle over $M$. A connection $\Gamma$ on $L E$ in our sense is a classical general connection on $L E$. Hence our approach leads to the original idea of the torsion of a general connection $\Gamma$ on $L E$. If $w_{q}^{p}$ are the induced fiber coordinates on $L E$, the usual coordinate expression of $\Gamma$ is $\mathrm{d} w_{q}^{p}=F_{q i}^{p}\left(x^{j}, w_{s}^{r}\right) \mathrm{d} x^{i}$. Then $-F_{q i}^{p}\left(x^{j}, \delta_{s}^{r}\right)$ is the coordinate form of the torsion of $\Gamma$. Of course, if we take for $\Gamma$ the tensor product $\Delta \otimes \Delta^{*}$ of a linear connection $\Delta$ on $E$ and of the dual connection $\Delta^{*}$ on $E^{*}$, [10], then the torsion of $\Delta \otimes \Delta^{*}$ vanishes, for $I$ is invariant with respect to $\Delta \otimes \Delta^{*}$.

## 8. The vector bundle case

Assume $p: E_{2} \rightarrow M$ is a vector bundle. Then each fiber of $\mathscr{F}\left(E_{1}, E_{2}\right)$ is a vector space, provided the linear operations on $C^{\infty}\left(E_{1 x}, E_{2 x}\right)$ are defined by extending the linear operations on $E_{2 x}$. In other words, $\mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow M$ is a vector bundle over sets, cf. [4]. Such a vector bundle structure is further extended to $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ by

$$
j_{x}^{1} s_{1}+j_{x}^{1} s_{2}=j_{x}^{1}\left(\tilde{s}_{1}+\tilde{s}_{2}\right), \quad j_{x}^{1}(k s)=j_{x}^{1} k \tilde{s}, \quad k \in \mathbb{R}
$$

with addition and multiplication by reals in $E_{2}$. Hence $J^{1} \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow M$ also is a vector bundle over sets.

Definition 19. A connection $\Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow J^{1} \mathscr{F}\left(E_{1}, E_{2}\right)$ is called linear if $\Gamma$ is a linear morphism over $M$.

In the case of an $r$-th order linear connection, its associated map (22) has the form

$$
\begin{equation*}
\Phi_{i b}^{a}(x, y) z^{b}+\Phi_{i b}^{a q}(x, y) z_{q}^{b}+\ldots+\Phi_{i b}^{a \alpha} z_{\alpha}^{b} \tag{49}
\end{equation*}
$$

If $E_{2}$ is a vector bundle, then $V E_{2}=E_{2} \times E_{2}$, which implies

$$
\mathscr{F}\left(E_{1}, V E_{2} \otimes \Lambda^{2} T^{*} M\right)=\mathscr{F}\left(E_{1}, E_{2}\right) \underset{M}{\times \mathscr{F}}\left(E_{1}, E_{2} \otimes \Lambda^{2} T^{*} M\right) .
$$

In this case, analogously to the classical situation, the curvature will be interpreted as the second component of the map from Definition 15,

$$
C \Gamma: \mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \mathscr{F}\left(E_{1}, E_{2} \otimes \Lambda^{2} T^{*} M\right)
$$

while the first component is the identity.
Proposition 15. For every differentiable linear connection $\Gamma$, the map $C \Gamma$ : $\mathscr{F}\left(E_{1}, E_{2}\right) \rightarrow \mathscr{F}\left(E_{1}, E_{2} \otimes \Lambda^{2} T^{*} M\right)$ is a linear morphism over $M$.

Proof. One easily verifies that in the linear case both $\tilde{\Gamma}$ and $\Delta$ in Definition 15 are linear morphisms over $M$.

Quite similarly, if $E_{2}$ is a vector bundle, then the absolute derivative $\nabla_{X} s$ of a section $s$ with respect to a vector field $X$ on $M$ is identified with the second component of (47), so that it is section of $\mathscr{F}\left(E_{1}, E_{2}\right)$ as well.

We finally remark that several other ideas from the classical theory of connections can be generalized to the case of $\mathscr{F}\left(E_{1}, E_{2}\right)$. The most interesting ones could be the vertical prolongation of $\Gamma$, the connections on $T \mathscr{F}\left(E_{1}, E_{2}\right) \subset \mathscr{F}\left(T E_{1} \rightarrow\right.$ $\left.T M, T E_{2} \rightarrow T M\right)$ or a detailed study of the absolute differentiation in the linear case. Such a research can be based on some general ideas from the theory of classical connections collected in the book [10].

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