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# ON THE MODIFIED LOGARITHMIC POTENTIAL 

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Introduction. Let $K$ be a simple oriented path of finite length in the complex plane $R^{2}$. Given a continuous real-valued function $F$ on $K$, consider the corresponding integral of the Cauchy type

$$
\begin{equation*}
\int_{K} \frac{F(\xi)}{\xi-z} \mathrm{~d} \xi \tag{1}
\end{equation*}
$$

as well as its real part

$$
\begin{equation*}
P_{K} F(z)=\operatorname{Re} \int_{K} \frac{F(\xi)}{\xi-z} \mathrm{~d} \xi \tag{2}
\end{equation*}
$$

(called the modified logaritmic potential with density $F$ ) and imaginary part

$$
\begin{equation*}
W_{K} F(z)=\operatorname{Im} \int_{K} \frac{F(\xi)}{\xi-z} \mathrm{~d} \xi \tag{3}
\end{equation*}
$$

(which is the double layer logarithmic potential with density $F$ ). Investigation of the behavior of $P_{K} F(z)$ and $W_{K} F(z)$ as $z \notin K$ approaches $K$ is of importance for a number of applications (see [1], [9]). Under additional assumptions on $K$ (like smoothness and Ljapunov condition) and $F$ (like Hölder continuity) the integral

$$
\int_{K} \frac{F(\xi)-F(\eta)}{\xi-z} \mathrm{~d} \xi
$$

is well known to possess angular limits as $z$ tends to a fixed point $\eta \in K$. Necessary and sufficient conditions on $K$ guaranteeing the existence of angular limits of $W_{K} F(z)$ at $\eta$ for arbitrary continuous $F$ have been established in [3]. The present paper deals with angular limits of $P_{K} F$. We fix a bounded lower-semicontinuous function $Q \geqq 0$ and consider the class $\Omega_{Q}(\eta)$ of all continuous real-valued functions $F$ satisfying

$$
\begin{equation*}
F(\xi)-F(\eta)=o(Q(\xi)) \quad \text { as } \quad \xi \rightarrow \eta \tag{4}
\end{equation*}
$$

Our main objective is to determine necessary and sufficient conditions on $K$ (whose end-point and initial-point are denoted by $\beta$ and $\alpha$, respectively) guaranteeing, for any $F \in \Omega_{Q}(\eta)$, the existence of angular limits of

$$
P_{K} F(z)-F(\eta) \log \frac{|\beta-z|}{|\alpha-z|}
$$

at $\eta$. For this purpose it is useful to associate with $K$ the following simple geometric quantities generalizing those introduced in [2]. Let us form the sum

$$
U_{K}^{\varrho}(\varrho, \eta)=\sum_{\xi} Q(\xi), \quad \xi \in K \cap\{\xi ;|\xi-\eta|=\varrho\}
$$

counting, with the weight $Q(\xi)$, the points $\xi$ in the intersection of $K$ and the circumference of center $\eta$ and radius $\varrho$. Then $U_{K}^{\varrho}(\varrho, \eta)$ is a Lebesgue measurable function of the variable $\varrho>0$ and we may put

$$
U_{K}^{\varrho}(\eta)=\int_{0}^{\infty} \varrho^{-1} U_{K}^{\varrho}(\varrho, \eta) \mathrm{d} \varrho .
$$

Consider also, for each $\gamma \in\langle 0,2 \pi)$ and $r>0$, the segment $S_{r}^{\eta}(\eta)=\left\{\eta+\varrho e^{i \gamma}\right.$; $0<\varrho<r\}$ and introduce the sum

$$
V_{K r}^{Q}(\gamma, \eta)=\sum_{\xi}|\xi-\eta| Q(\xi), \quad \xi \in K \cap S_{r}^{\gamma}(\eta),
$$

counting, with the weight $|\xi-\eta| Q(\xi)$, the points $\xi$ in the intersection of $K$ and $S_{r}^{\gamma}(\eta)$. Since $V_{K r}^{Q}(\gamma, \eta)$ is a Lebesgue measurable function of the variable $\gamma \in\langle 0,2 \pi)$, we are justified to define

$$
V_{K r}^{Q}(\eta)=\int_{0}^{2 \pi} V_{K r}^{Q}(\gamma, \eta) \mathrm{d} \gamma .
$$

With this notation we are now in position to formulate the following typical corollary of main results (some of whose have been announced without proofs in [4], [7]) established below.

Theorem. Let $S \subset R^{2} \backslash K$ be a connected set whose closure meets $K$ at $\eta$ only. Suppose that the contingent of $S$ at $\eta$ (in the sense of [11], chap. IX, § $2-$ see also theorem 9 below) together with its reflection in $\eta$ is disjoint from the contingent of $K$ at $\eta$.

If

$$
\underset{\substack{z \rightarrow \eta \\ z \in S}}{\lim \sup }\left|P_{K} F(z)-F(\eta) \log \frac{|\beta-z|}{|\alpha-z|}\right|<\infty
$$

for any $F \in \Omega_{Q}(\eta)$, then

$$
\begin{equation*}
U_{K}^{Q}(\eta)+\sup _{r>0} r^{-1} V_{K r}^{Q}(\eta)<\infty . \tag{5}
\end{equation*}
$$

Conversely, suppose that (5) holds. Let $\theta(r) \geqq 0$ be a bounded continuous nondecreasing function of the variable $r \geqq 0, \theta \neq 0$. If $F$ is a bounded Baire function on $K$ satisfying

$$
F(\xi)-F(\eta)=O(\theta(|\xi-\eta|) Q(\xi)) \text { as } \quad \xi \rightarrow \eta
$$

then the integral

$$
\int_{K} \frac{F(\xi)-F(\eta)}{|\xi-\eta|} \mathrm{d}|\xi-\eta|=P_{K}^{0} F(\eta)
$$

converges and for $z \in S$ the following estimate holds

$$
\begin{equation*}
P_{K} F(z)-P_{K}^{0} F(\eta)-F(\eta) \log \frac{|\beta-z|}{|\alpha-z|}=O\left(|z-\eta| \int_{|z-\eta|}^{\infty} r^{-2} \theta(r) \mathrm{d} r\right) \tag{6}
\end{equation*}
$$

as $z \rightarrow \eta$.
If $F$ satisfies

$$
F(\check{\xi})-F(\eta)=o(\theta(|\xi-\eta|) Q(\xi)) \quad \text { as } \quad \xi \rightarrow \eta
$$

then the right-hand side in (6) can be replaced by

$$
o\left(|z-\eta| \int_{|z-\eta|}^{\infty} r^{-2} \theta(r) \mathrm{d} r\right)
$$

or

$$
O(|z-\eta|)
$$

according as the integral $\int_{0}^{\infty} r^{-2} \theta(r) \mathrm{d} r$ diverges or converges.

1. Notation. If $f$ is a complex- or real-valued function defined on an interval $J \subset R^{1}$ then, for each set $G \subset J$ which is open in $J$, $\operatorname{var} f(G)$ will denote the variation of $f$ on $G$; thus $\operatorname{var} f(\emptyset)=0$ and, for $G \neq \emptyset, \operatorname{var} f(G)$ is the least upper bound of all the sums

$$
\sum_{j=1}^{n}\left|f\left(b_{j}\right)-f\left(a_{j}\right)\right|
$$

where $\left\langle a_{1}, b_{1}\right\rangle, \ldots,\left\langle a_{n}, b_{n}\right\rangle$ are non-overlapping compact intervals contained in $G$. For any $M \subset J$ we let

$$
\operatorname{var} f(M)=\inf _{G} \operatorname{var} f(G),
$$

where $G$ runs over all sets $G \subset J$ that are open in $J$ and contain $M$. If necessary, we shall also use the more explicit notation like $\operatorname{var}_{t}[f(t) ; M]$ to denote $\operatorname{var} f(M)$. As it is well known, var $f(\ldots)$ is a Carathéodory outer measure; its restriction to
var $f$-measurable subsets of $J$ is a measure. The integral (over $M \subset J$ ) of an extended real-valued function $F$ with respect to this measure will be denoted by the symbols

$$
\int_{M} F \mathrm{~d} \operatorname{var} f, \quad \int_{M} F(t) \mathrm{d} \operatorname{var} f(t),
$$

We shall say that $f$ has locally finite variation on $J$ provided var $f(I)<\infty$ for every compact interval $I \subset J$. For such $f$ the integral

$$
\int_{M} F \mathrm{~d} f\left(=\int_{M} F(u) \mathrm{d} f(u)\right)
$$

is always to be understood in the sense of Lebesgue-Stieltjes.
We shall now recall several known basic lemmas to be used below.
2. Lemma. Let $f$ be a continuous real-valued function of bounded variation on $\langle a, b\rangle$ and let $p$ be a function on $f(\langle a, b\rangle)$. Suppose that $p$ has a continuous derivative on $f(\langle a, b\rangle)$ and put $h=p \circ f(=$ the composite of $f$ and $p)$. Then $h$ has bounded variation on $\langle a, b\rangle$ and, for each lower-semicontinuous (extended real-valued) function $F \geqq 0$ on $\langle a, b\rangle$,

$$
\int_{a}^{b} F \mathrm{~d} \operatorname{var} h=\int_{a}^{b}\left|p^{\prime}(f(t))\right| F(t) \mathrm{d} \operatorname{var} f(t) .
$$

3. Lemma. Let $f, g$ be continuous functions having locally finite variation on an interval $J$. Then, for each lower-semicontinuous function $F \geqq 0$ on $J$,

$$
\int_{J} F \mathrm{~d} \operatorname{var}(f . g) \leqq \int_{J} F|f| \mathrm{d} \operatorname{var} g+\int_{J} F|g| \mathrm{d} \operatorname{var} f .
$$

4. Lemma. Let $f$ be a continuous real-valued function having locally finite variation on an interval $J$. Suppose that $F \geqq 0$ is a lower-semicontinuous function on $J$ and denote, for each $u \in R^{1}$, by $\sigma(u ; F)$ the sum

$$
\sum_{t} F(t), \quad f(t)=u,
$$

which is extended over all $t \in J$ with $f(t)=u$ (so that $\sigma(u ; F)=0$ provided $u \notin f(J)$ and $\sigma(u ; F)=+\infty$ whenever $F(t)>0$ for uncountably many $t \in J$ with $f(t)=u)$. Then $\sigma(u ; F)$ is a Lebesgue measurable function of the variable $u \in R^{1}$ and

$$
\int_{-\infty}^{+\infty} \sigma(u ; F) \mathrm{d} u=\int_{J} F \mathrm{~d} \operatorname{var} f .
$$

5. Lemma. Let $f$ be a continuous (real-or complex-valued) function having locally finite variation on an interval $J$. Then $\operatorname{var} f(M)=0$ for each $M \subset J$ with countable $f(M)$.

For continuous $F$, elementary proofs of lemmas 2-4 may be found in [5] (see theorems 6.22, 6.21, 6.17); their extension to the case of a lower-semicontinuous $F$ is immediate since such an $F$ is a limit of a non-decreasing sequence of continuous functions. Let us note here that proof of lemma 4 is based on Banach's theorem on variation of a continuous function (see also [8], [10]). For the proof of lemma 5 (which is, in fact, an easy consequence of lemma 4) see, e.g., [3], lemma 3.4.
6. Notation. In what follows we shall always assume that $\psi$ is a continuous complexvalued function of bounded variation on $\langle a, b\rangle$ and $q \geqq 0$ is a bounded lowersemicontinuous function on $\langle a, b\rangle . R^{2}$ will denote the Euclidean plane whose points will be identified with complex numbers. Given $z \in R^{2}$ and $\varrho>0$ we put

$$
\begin{aligned}
\Omega_{\varrho}(z) & =\left\{\xi \in R^{2} ;|\xi-z|<\varrho\right\}, \\
u_{\psi}^{q}(\varrho, z) & =\sum_{t} q(t), \quad|\psi(t)-z|=\varrho,
\end{aligned}
$$

the last sum being extended over all $t \in\langle a, b\rangle$ with $|\psi(t)-z|=\varrho$; similarly, put for any $\gamma \in\langle 0,2 \pi\rangle$

$$
v_{\psi_{\varrho}}^{q}(\gamma, z)=\sum_{t} q(t)|\psi(t)-z|,
$$

where now the sum is extended over all $t \in\langle a, b\rangle$ satisfying

$$
\varrho>|\psi(t)-z|>0, \quad \psi(t)-z=|\psi(t)-z| \mathrm{e}^{i \gamma}
$$

If $f$ is a function in $R^{1}$, then spt $f$ will denote the support of $f$.
7. Lemma. For fixed $z \in R^{2}, u_{\psi}^{q}(\varrho, z)$ is a Lebesgue measurable function of the variable $\varrho$ and, for fixed $\varrho>0, v_{\psi, \ell}^{q}(\gamma, z)$ is a Lebesgue measurable function of the variable $\gamma \in\langle 0,2 \pi\rangle$. The integrals

$$
\begin{equation*}
u_{\psi}^{q}(z)=\int_{0}^{\infty} \varrho^{-1} u_{\psi}^{q}(\varrho, z) \mathrm{d} \varrho, \quad v_{\psi \varrho}^{q}(z)=\int_{0}^{2 \pi} v_{\psi \varrho}^{q}(\gamma, z) \mathrm{d} \gamma \tag{7}
\end{equation*}
$$

are lower-semicontinuous functions of the variable $z \in R^{2}$.
Proof. Given $z \in R^{2}$ and $\varrho>0$ we denote by $\mathscr{S}_{\varrho}(z)$ the system of all components of $\{t \in\langle a, b\rangle ; 0<|\psi(t)-z|<\varrho\}$. With each $J \in \mathscr{S}_{\varrho}(z)$ we associate a continuous argument $\vartheta_{z}(t ; J)$ of $\psi(t)-z$ on $J$. Denote by $F_{J}$ the restriction of

$$
\begin{equation*}
F(t)=q(t)|\psi(t)-z| \tag{8}
\end{equation*}
$$

to $J$.

Employing lemma 4 we get

$$
\int_{J} F(t) \mathrm{d} \operatorname{var}_{t} \vartheta_{z}(t ; J)=\sum_{n=-\infty}^{+\infty} \int_{0}^{2 \pi} \sigma\left(\gamma+2 n \pi ; F_{J}\right) \mathrm{d} \gamma .
$$

Let us agree to write briefly $\sum_{J} \ldots$ for the sum extended over all $J \in \mathscr{S}_{\varrho}(z)$. Since

$$
\sum_{J} \sum_{n=-\infty}^{+\infty} \sigma\left(\gamma+2 n \pi ; F_{J}\right)=v_{\psi \varrho}^{q}(\gamma, z), \quad \gamma \in\langle 0,2 \pi\rangle,
$$

we conclude that $v_{\psi e}^{q}(\gamma, z)$ is a Lebesgue measurable function of the variable $\gamma \in$ $\in\langle 0,2 \pi\rangle$ and

$$
\begin{equation*}
v_{\psi e}^{q}(z)=\sum_{J} \int_{J} F(t) \mathrm{d}^{\operatorname{var}_{t}} \vartheta_{z}(t ; J) \tag{9}
\end{equation*}
$$

Let $\mathscr{F}_{\varrho}(z)$ be the class of all continuous real-valued functions $f$ in $R^{1}$ such that $|f| \leqq 1$ and

$$
\operatorname{spt} f \subset\{t \in\langle a, b\rangle ; 0<|\psi(t)-z|<\varrho\} .
$$

We shall first observe that

$$
\begin{equation*}
v_{\psi e}^{q}(z)=\sup \left\{\operatorname{Im} \int_{a}^{b} F(t) \frac{f(t)}{\psi(t)-z} \mathrm{~d} \psi(t) ; f \in \mathscr{F}_{e}(z)\right\} \tag{10}
\end{equation*}
$$

where, of course $f(t) /(\psi(t)-z)$ means 0 outside spt $f \subset\{t ; \psi(t) \neq z\}$. Indeed, if $f \in \mathscr{F}_{e}(z)$, then there is a finite number of components $J_{1}, \ldots, J_{s} \in \mathscr{S}_{e}(z)$ such that

$$
\operatorname{spt} f \subset \bigcup_{j=1}^{s} J_{j}
$$

Fix $t_{j} \in J_{j}$. One easily verifies that, for $t \in J_{j}$,

$$
\operatorname{Im} \int_{t_{j}}^{t} \frac{\mathrm{~d} \psi(u)}{\psi(u)-z}
$$

differs only by an additive constant from $\vartheta_{z}\left(t ; J_{j}\right)$ (see 7.43 in [5]). Consequently,

$$
\operatorname{Im} \int_{a}^{b} F(t) f(t) \frac{\mathrm{d} \psi(t)}{\psi(t)-z}=\sum_{j=1}^{s} \int_{J_{j}} F(t) f(t) \mathrm{d}_{t} \vartheta_{z}\left(t ; J_{j}\right) \leqq v_{\psi_{e}}^{q}(z)
$$

on account of (9). Fix now an arbitrary $k<v_{\psi e}^{q}(z)$. Then there is a finite number of components $J_{1}, \ldots, J_{n} \in \mathscr{S}_{\varrho}(z)$ such that

$$
\sum_{j=1}^{n} \int_{J_{j}} F(t) \mathrm{dvar}_{t} \vartheta_{z}\left(t ; J_{j}\right)>k
$$

For each $j(=1, \ldots, n)$ choose a $k_{j}<\int_{J_{j}} F(t) \mathrm{d}_{\operatorname{var}}^{t} \vartheta_{z}\left(t ; J_{j}\right)$ and an $f_{j} \in \mathscr{F}_{g}(z)$ with spt $f_{j} \subset J_{j}$ such that

$$
\sum_{j=1}^{n} k_{j}=k
$$

and

$$
\int_{J_{j}} f_{j}(t) F(t) \mathrm{d}_{t} \vartheta_{z}\left(t ; J_{j}\right)>k_{j}, \quad 1 \leqq j \leqq n
$$

Defining $f=\sum_{j=1}^{n} f_{j}$ we get $f \in \mathscr{F}_{e}(z)$ and

$$
\operatorname{Im} \int_{a}^{b} F(t) f(t) \frac{\mathrm{d} \psi(t)}{\psi(t)-z}=\sum_{j=1}^{n} \int_{J_{j}} f_{j}(t) F(t) \mathrm{d}_{t} \vartheta_{z}\left(t ; J_{j}\right)>\sum_{j=1}^{n} k_{j}=k .
$$

Thus (10) is established.
Given $f \in \mathscr{F}_{e}(z)$, there is an $\varepsilon>0$ such that $f \in \mathscr{F}_{e}(\xi)$ for any $\xi \in \Omega_{\varepsilon}(z)$. Since

$$
\int_{a}^{b} F(t) f(t) \frac{\mathrm{d} \psi(t)}{\psi(t)-\xi}
$$

is a continuous function of $\xi \in \Omega_{\varepsilon}(z)$, we conclude from (10) that $v_{\psi e}^{q}(\ldots)$ is lowersemicontinuous at $z$.

Let now $J$ run over $\mathscr{S}_{\infty}(z)$. Employing lemma 4 one easily obtains that $u_{\psi}^{q}(\varrho, z)$ is a Lebesgue measurable function of the variable $\varrho$ and

$$
\begin{equation*}
u_{\psi}^{q}(z)=\sum_{J} \int_{J} \frac{q(t)}{|\psi(t)-z|} \mathrm{d}_{\operatorname{var}}^{t}|\psi(t)-z| . \tag{11}
\end{equation*}
$$

Hence

$$
u_{\psi}^{q}(z)=\sup \left\{\operatorname{Re} \int_{a}^{b} f(t) q(t) \frac{\mathrm{d} \psi(t)}{\psi(t)-z} ; \quad f \in \mathscr{F}_{\infty}(z)\right\}
$$

and the lower-semicontinuity of $u_{\psi}^{q}(\ldots)$ at $z$ follows.
8. Notation. If $f$ is a bounded Baire function on $\langle a, b\rangle$ we define for $z \in R^{2} \psi \psi(\langle a, \mathrm{~b}\rangle)$

$$
p_{\psi} f(z)=\int_{a}^{b} \frac{f(t) \mid}{\psi(t)-z \mid} \mathrm{d}_{t}|\psi(t)-z|\left(=\operatorname{Re} \int_{a}^{b} \frac{f(t)}{\psi(t)-z} \mathrm{~d} \psi(t)\right) .
$$

Given $S \subset R^{2}$ and $\eta \in R^{2}$ we denote by

$$
S \odot \eta=S \cup\{2 \eta-\xi ; \xi \in S\}
$$

the union of $S$ and its reflection in $\eta$.
9. Theorem. Let $S \subset R^{2} \backslash \psi(\langle a, b\rangle)$ be a connected set whose closure meets $\psi(\langle a, b\rangle)$ at $\eta$ only.

Suppose that

$$
\begin{equation*}
\lim \sup _{\substack{z \rightarrow \eta \\ z \in S}}\left|p_{\psi} f(z)\right|<\infty \tag{12}
\end{equation*}
$$

for each continuous function $f$ on $\langle a, b\rangle$ satisfying

$$
\begin{equation*}
f(t)=o(q(t)) \quad \text { as } \quad \psi(t) \rightarrow \eta . \tag{13}
\end{equation*}
$$

Then

$$
u_{\psi}^{q}(\eta)<\infty .
$$

If, besides that, the contingent $\left.{ }^{1}\right)$ of $\psi(\langle a, b\rangle)$ at $\eta$ does not meet the contingent of $S \odot \eta$ at $\eta$, then

$$
\sup _{r>0} r^{-1} v_{\psi r}^{q}(\eta)<\infty .
$$

Proof. Consider the class $\mathscr{C}_{q}$ of all continuous functions $f$ on $\langle a, b\rangle$ vanishing on $\{t \in\langle a, b\rangle ; \psi(t)=\eta\}$ and satisfying (13) as well as

$$
\begin{equation*}
|f| \leqq c_{f} q \tag{14}
\end{equation*}
$$

for suitable constant $c_{f}$ (depending on $f$ ). Defining $\|f\|$ as the least upper bound of all $c_{f}$ satisfying (14) we get a norm on $\mathscr{C}_{q}$ which turns $\mathscr{C}_{q}$ into a Banach space. Note that, for $f \in \mathscr{C}_{q}$ and $z \notin \psi(\langle a, b\rangle)$,

$$
\begin{equation*}
\left|p_{\psi} f(z)\right| \leqq\|f\| \int_{a}^{b} q \mathrm{~d} \operatorname{var} \psi / \operatorname{dist}(\psi(\langle a, b\rangle), z), \tag{15}
\end{equation*}
$$

where $\operatorname{dist}(\psi(\langle a, b\rangle), z)=\inf \{|\psi(t)-z| ; t \in\langle a, b\rangle\}$. Combining (15) with the assumption (12) we conclude that

$$
\begin{equation*}
f \in \mathscr{C}_{q} \Rightarrow \sup _{z \in S}\left|p_{\psi} f(z)\right|<\infty . \tag{16}
\end{equation*}
$$

With each $z \in S$ we associate the functional $L_{z}$ defined by

$$
L_{z}(f)=p_{\psi} f(z), \quad f \in \mathscr{C}_{q} .
$$

Clearly, each $L_{z}$ is a bounded linear functional on $\mathscr{C}_{q}$ whose norm is given by

$$
\begin{equation*}
\left\|L_{z}\right\|=\int_{a}^{b} \frac{q(t)}{|\psi(t)-z|} \mathrm{dvar}_{t}|\psi(t)-z| \tag{17}
\end{equation*}
$$

[^0]Since, in view of (16), all the functionals in $\left\{L_{z}\right\}_{z \in S}$ are pointwise bounded on $\mathscr{C}_{q}$, we conclude by the principle of uniform boundedness that

$$
\begin{equation*}
\sup _{z \in S}\left\|L_{z}\right\|=c<\infty \tag{18}
\end{equation*}
$$

According to (11) we have for $z \in S\left(\subset R^{2} \backslash \psi(\langle a, b\rangle)\right)$

$$
\begin{equation*}
u_{\psi}^{q}(z)=\int_{a}^{b} \frac{q(t)}{|\psi(t)-z|} \mathrm{d}_{\operatorname{var}_{t}}|\psi(t)-z| . \tag{19}
\end{equation*}
$$

Combining this with (17), (18) we arrive at

$$
\begin{equation*}
\sup _{z \in S} u_{\psi}^{q}(z)=c<\infty, \tag{20}
\end{equation*}
$$

which implies

$$
\begin{equation*}
u_{\psi}^{q}(\eta) \leqq c \tag{21}
\end{equation*}
$$

by the lower-semicontinuity of $u_{\psi}^{q}(\ldots)$ established in lemma 7. Suppose now that the contingent of $\psi(\langle a, b\rangle)$ at $\eta$ is disjoint from the contigent of $S \odot \eta$ at $\eta$. Given $r>0$ denote by $\mathscr{S}_{r}(\eta)$ the system of all components of $\{t \in\langle a, b\rangle ; 0<|\psi(t)-\eta|<r\}$. With each $I \in \mathscr{S}_{\infty}(\eta)$ and $z \in R^{2} \backslash\{\eta\}$ we associate a continuous argument $\omega_{z}(t ; I)$ of $(\psi(t)-\eta) /(z-\eta)$ on $I$. It is easily seen that there is an $R_{0}>0$ such that for $t \in I \in$ $\in \mathscr{S}_{\infty}(\eta)$ and $z \in S$

$$
\begin{equation*}
\left(|z-\eta|<R_{0},|\psi(t)-\eta|<R_{0}\right) \Rightarrow\left|\sin \omega_{z}(t ; I)\right| \geqq R_{0} . \tag{22}
\end{equation*}
$$

We may assume $R_{0}$ to be small enough to guarantee

$$
S \cap\left\{z ;|z-\eta|=R_{0}\right\} \neq \emptyset
$$

(note that $S$ is connected and $\eta$ belongs to the closure of $S$ ). Consider now an arbitrary $r$ with $0<r \leqq R_{0}$ and choose a $z \in S$ with $|z-\eta|=r$. Consider a $J \in \mathscr{S}_{r}(\eta)$. There is a uniquely determined $I_{J} \in \mathscr{S}_{\infty}(\eta)$ containing $J$ and we put

$$
\omega_{z}(t ; J)=\omega_{z}\left(t ; I_{J}\right), \quad t \in J
$$

For the sake of brevity, we shall also write

$$
\varrho_{\xi}(t)=|\psi(t)-\xi|, \quad t \in\langle a, b\rangle, \quad \xi \in R^{2} .
$$

With this notation we have for $t \in J\left(\in \mathscr{S}_{\mathbf{r}}(\eta)\right)$

$$
\begin{gathered}
\varrho_{z}(t) \leqq \varrho_{\eta}(t)+|z-\eta| \leqq 2 r, \\
\varrho_{z}^{2}(t)=\varrho_{\eta}^{2}(t)+r^{2}-2 r \varrho_{\eta}(t) \cos \omega_{z}(t ; J),
\end{gathered}
$$

whence

$$
\begin{aligned}
& \int_{J} q(t) \varrho_{z}^{-2}(t) \mathrm{d} \operatorname{var}_{t} \varrho_{z}^{2}(t) \geqq(2 r)^{-2} \int_{J} q(t) \mathrm{dvar}_{t}\left[\varrho_{\eta}^{2}(t)-2 r \varrho_{\eta}(t) \cos \omega_{z}(t ; J)\right] \geqq \\
& \geqq(2 r)^{-1} \int_{J} q(t) \mathrm{d}_{\mathrm{var}}^{t}\left[\begin{array}{l}
\eta \\
\left.(t) \cos \omega_{z}(t ; J)\right]- \\
\end{array}\right.
\end{aligned}
$$

$$
\begin{aligned}
& -\frac{1}{4} \int_{J} q(t) \varrho_{\eta}^{-2}(t) \mathrm{d} \operatorname{var}_{t} \varrho_{\eta}^{2}(t)=(\text { see lemma } 2)= \\
& =(2 r)^{-1} \int_{J} q(t) \mathrm{d}_{\operatorname{var}}^{t}\left[\varrho_{\eta}(t) \cos \omega_{z}(t ; J)\right]-\frac{1}{2} \int_{J} q(t) \varrho_{\eta}^{-1}(t) \mathrm{d}_{\operatorname{var}}^{t} \varrho_{\eta}(t) .
\end{aligned}
$$

Hence we get by (19), (20) and lemma 2 letting $J$ run over $\mathscr{S}_{r}(\eta)$

$$
\begin{aligned}
& c \geqq u_{\psi}^{q}(z)=\int_{a}^{b} q(t) \varrho_{z}^{-1}(t){\mathrm{d} \operatorname{var}_{t} \varrho_{z}(t)=} \\
& =\frac{1}{2} \int_{a}^{b} q(t) \varrho_{z}^{-2}(t) \mathrm{d} \operatorname{var}_{t} \varrho_{z}^{2}(t) \geqq(4 r)^{-1} \sum_{J} \int_{J} q(t) \mathrm{d}^{\operatorname{var}}{ }_{t}\left[\varrho_{\eta}(t) \cos \omega_{z}(t ; J)\right]- \\
& -\frac{1}{4} \sum_{J} \int_{J} q(t) \varrho_{\eta}^{-1}(t) \mathrm{d}^{\operatorname{var}} \varrho_{t} \varrho_{\eta}(t)=
\end{aligned}
$$

$=($ employ (11) with $z$ replaced by $\eta)=(4 r)^{-1} \sum_{J} \int_{J} q(t) \mathrm{d} \mathrm{var}_{t}\left[\varrho_{\eta}(t) \cos \omega_{z}(t ; J)\right]-$

$$
\begin{gathered}
-\frac{1}{4} u_{\psi}^{q}(\eta) \geqq(\operatorname{see}(21)) \geqq \\
\geqq \frac{1}{4}\left\{r^{-1} \sum_{J} \int_{J} q(t) \mathrm{d}^{\operatorname{var}}\left[\varrho_{\eta}(t) \cos \omega_{z}(t ; J)\right]-c\right\}
\end{gathered}
$$

Consequently,

$$
\begin{equation*}
r^{-1} \sum_{J} \int_{J} q(t) \mathrm{d} \operatorname{var}_{t}\left[\varrho_{\eta}(t) \cos \omega_{z}(t ; J)\right] \leqq 5 c . \tag{23}
\end{equation*}
$$

Using lemmas 2 and 3 we obtain

$$
\begin{aligned}
& \sum_{J} \int_{J} \varrho_{\eta}(t) q(t) \mathrm{d} \operatorname{var}_{t}\left[\cos \omega_{z}(t ; J)\right]= \\
= & \sum_{J} \int_{J} \varrho_{\eta}(t) q(t) \mathrm{dvar}_{t} \frac{\varrho_{\eta}(t) \cos \omega_{z}(t ; J)}{\varrho_{\eta}(t)} \leqq
\end{aligned}
$$

$$
\begin{gathered}
\leqq \sum_{J} \int_{J} q(t) \mathrm{d} \operatorname{var}_{t}\left[\varrho_{\eta}(t) \cos \omega_{z}(t ; J)\right]+\sum_{J} \int_{J} q(t)\left|\cos \omega_{z}(t ; J)\right| \mathrm{d} \operatorname{var}_{t} \varrho_{\eta}(t) \leqq \\
\leqq(\operatorname{see}(23)) \leqq 5 c r+\sum_{J} r \int_{J} \varrho_{\eta}^{-1}(t) q(t){\mathrm{d} \operatorname{var}_{t} \varrho_{\eta}(t)=}^{=5 c r+r u_{\psi}^{q}(\eta) \leqq(\operatorname{see}(21)) \leqq 6 c r .}
\end{gathered}
$$

On the other hand, lemma 2 together with (22) yield for any $J \in \mathscr{S}_{r}(\eta)$

$$
\begin{aligned}
& \int_{J} \varrho_{\eta}(t) q(t) \mathrm{d}_{\operatorname{var}}^{t}\left[\cos \omega_{z}(t ; J)\right]=\int_{J} \varrho_{\eta}(t) q(t)\left|\sin \omega_{z}(t ; J)\right| \mathrm{dvar}_{t} \omega_{z}(t ; J) \geqq \\
& \geqq R_{0} \int_{J} \varrho_{\eta}(t) q(t) \mathrm{d} \operatorname{var}_{t} \omega_{z}(t ; J),
\end{aligned}
$$

whence we conclude

$$
\sum_{J} \int_{J} \varrho_{\eta}(t) q(t) \mathrm{d} \operatorname{var}_{t} \omega_{z}(t ; J) \leqq 6 c r R_{0}^{-1}
$$

Noting that $\omega_{z}(i ; J)$ differs only by an additive constant from a continuous argument of $\psi(t)-\eta$ on $J$ we have by (9), (8)

$$
\sum_{J} \int_{J} \varrho_{\eta}(t) q(t) \mathrm{d}_{\operatorname{var}}^{t} \omega_{z}(t ; J)=v_{\psi r}^{q}(\eta)
$$

We have thus shown that

$$
0<r \leqq R_{0} \Rightarrow r^{-1} v_{\psi r}^{q}(\eta) \leqq 6 c R_{0}^{-1}
$$

Consider now an arbitrary $r>R_{0}$. It follows easily from (10), (8) that

$$
r^{-1} v_{\psi r}^{q}(\eta) \leqq R_{0}^{-1} \int_{a}^{b} q(t) \mathrm{d} \operatorname{var} \psi(t)
$$

so that

$$
\sup _{r>0} r^{-1} v_{\psi r}^{q}(\eta) \leqq R_{0}^{-1} \max \left\{6 c, \int_{a}^{b} q \mathrm{~d} \operatorname{var} \psi\right\}
$$

and the proof is complete.
10. Remark. If $S \subset R^{2} \backslash \psi(\langle a, b\rangle)$ is a connected set whose closure meets $\psi(\langle a, b\rangle)$ at a single point $\eta$ such that the contingent of $S \odot \eta$ at $\eta$ is disjoint from the contingent of $\psi(\langle a, b\rangle)$ at $\eta$, then

$$
\begin{equation*}
\sup _{z \in S} u_{\psi}^{q}(z)<\infty \tag{24}
\end{equation*}
$$

implies

$$
\begin{equation*}
u_{\psi}^{q}(\eta)+\sup _{r>0} r^{-1} v_{\psi r r}^{q}(\eta)<\infty . \tag{25}
\end{equation*}
$$

This has been established in the course of the above proof. The converse of this assertion is also valid as shown in proposition 12 which will be needed below. Before going into its proof we shall recall the following known lemma (which, as shown in [6], may be used as a basis for development of the Lebesgue theory of integration).
11. Lemma. Let $\mu \geqq 0$ be a measure defined on Borel subsets of an interval I and suppose that $F \geqq 0$ is an extended real-valued Baire function on I. Given $\tau>0$ let

$$
F_{\tau}=\{t \in I ; F(t)>\tau\} .
$$

Then

$$
\int_{I} F \mathrm{~d} \mu=\int_{0}^{\infty} \mu\left(F_{\tau}\right) \mathrm{d} \tau .
$$

Now we are in position to prove the following
12. Proposition. Let $S \subset R^{2} \backslash \psi(\langle a, b\rangle)$ be a set whose closure meets $\psi(\langle a, b\rangle)$ at a single point $\eta$. Suppose that the contingent of $\psi(\langle a, b\rangle)$ at $\eta$ is disjoint from the contingent of $S \odot \eta$ at $\eta$.

If

$$
u_{\psi}^{q}(\eta)+\sup _{r>0} r^{-1} v_{\psi r}^{q}(\eta)<\infty
$$

then

$$
\sup _{z \in S} u_{\psi}^{q}(z)<\infty .
$$

Proof. In accordance with the notation introduced earlier we shall write $\mathscr{S}_{\infty}(\eta)$ for the system of all components of $\{t \in\langle a, b\rangle ;|\psi(t)-\eta|>0\}$ and, for each $I \in \mathscr{S}_{\infty}(\eta)$, we fix a continuous argument $\vartheta_{\eta}(t ; I)$ of $\psi(t)-\eta$ on $I$; given $z \in R^{2} \backslash\{\eta\}$ we denote by $\omega_{z}(t ; I)$ a continuous argument of $[\psi(t)-\eta] / z-\eta$ (so that $\vartheta_{\eta}(\ldots ; I)-$ $-\omega_{z}(\ldots ; I)$ is constant on $\left.I\right)$. For the sake of brevity we put for $z \in R^{2}$

$$
\begin{equation*}
\varrho_{z}(t)=|\psi(t)-z|, \quad t \in\langle a, b\rangle . \tag{26}
\end{equation*}
$$

We agree to use $I$ as a generic notation for elements of $\mathscr{S}_{\infty}(\eta)$ and set for $r>0$

$$
I_{r}=\left\{t \in I ; \varrho_{\eta}(t)<r\right\}, \quad I^{r}=I \backslash I_{r} .
$$

Let

$$
k=\sup _{r>0} r^{-1} v_{\psi r}^{q}(\eta)
$$

In view of (9), (8) we have for any $r>0$

$$
\begin{equation*}
v_{\psi r}^{q}(\eta)=\sum_{I} \int_{I_{r}} \varrho_{\eta}(t) q(t) \mathrm{d}_{\operatorname{var}}^{t} \vartheta_{\eta}(t ; I) \leqq k r . \tag{27}
\end{equation*}
$$

Consider now the sum

$$
\begin{equation*}
s(r)=\sum_{I} \int_{I^{r}} \varrho_{\eta}^{-1}(t) q(t) \mathrm{d} \operatorname{var}_{t} \vartheta_{\eta}(t ; I) \tag{28}
\end{equation*}
$$

Put for $\tau>0$

$$
I_{\tau}^{r}=\left\{t \in I ; r \leqq \varrho_{\eta}(t)<\tau^{-1 / 2}\right\},
$$

so that

$$
I_{\tau}^{r}=\left\{t \in I^{r} ; \varrho_{\eta}^{-2}(t)>\tau\right\} .
$$

## Defining

$$
\mu(B)=\int_{B} \varrho_{\eta}(t) q(t) \mathrm{d} \operatorname{var}_{t} \vartheta_{\eta}(t ; I)
$$

for Borel sets $B \subset I$ and employing lemma 11 one easily obtains

$$
\int_{I^{r}} \varrho_{\eta}^{-1}(t) q(t){\mathrm{d} \operatorname{var}_{t} \vartheta_{\eta}(t ; I)=\int_{I^{r}} \varrho_{\eta}^{-2}(t) \mathrm{d} \mu(t)=\int_{0}^{r^{-2}} \mu\left(I_{\tau}^{r}\right) \mathrm{d} \tau . . . . . . . . . .}
$$

Noting that $I_{\tau}^{r} \subset I_{x}$ with $x=\tau^{-1 / 2}$ we conclude from (27) that

$$
\sum_{I} \mu\left(I_{\tau}^{r}\right) \leqq \sum_{I} \mu\left(I_{x}\right) \leqq k x=k \tau^{-1 / 2}
$$

whence

$$
s(r) \leqq \int_{0}^{r^{-2}} k \tau^{-1 / 2} \mathrm{~d} \tau=2 k r^{-1}
$$

Consequently,

$$
\begin{equation*}
\sup _{r>0} r s(r) \leqq 2 k \tag{29}
\end{equation*}
$$

Denote by $T$ the union of $\mathscr{S}_{\infty}(\eta)$. Since

$$
\psi(\langle a, b\rangle \backslash T)=\{\eta\},
$$

we infer from lemma 5

$$
\operatorname{var} \psi(\langle a, b\rangle \backslash T)=0
$$

Let now $z$ be an arbitrary point in $R^{2}$. Since var $\psi=\operatorname{var}(\psi-z)$ dominates var $\varrho_{z}$, we have also

$$
\begin{equation*}
\operatorname{var} \varrho_{z}(\langle a, b\rangle \backslash T)=0 \tag{30}
\end{equation*}
$$

and (19) yields

$$
\begin{equation*}
u_{\psi}^{q}(z)=\sum_{I} \int_{I} \varrho_{z}^{-1}(t) q(t){\mathrm{d} \operatorname{var}_{t}} \varrho_{z}(t) \tag{31}
\end{equation*}
$$

In view of our assumptions concerning the contingents of $S \odot \eta$ and $\psi(\langle a, b\rangle)$ at $\eta$, there is an $R_{0}>0$ such that (22) holds for $z \in S$ and $t \in I\left(\in \mathscr{S}_{\infty}(\eta)\right)$. Let $z \in S$, put $|z-\eta|=r$ and assume $r<R_{0}$. Then

$$
\begin{equation*}
\frac{\varrho_{n}(t)}{\varrho_{z}(t)} \leqq \frac{1}{\sin \omega_{z}(t ; I)} \leqq R_{0}^{-1}, \tag{32}
\end{equation*}
$$

$$
\begin{equation*}
\frac{r}{\varrho_{z}(t)} \leqq \frac{1}{\sin \omega_{z}(t ; I)} \leqq R_{0}^{-1} \tag{33}
\end{equation*}
$$

and

$$
\begin{equation*}
\varrho_{z}^{2}(t)=r^{2}+\varrho_{\eta}^{2}(t)-2 r \varrho_{\eta}(t) \cos \omega_{z}(t ; I) . \tag{34}
\end{equation*}
$$

Hence we obtain using lemmas 2,3

Making use of (31), (27), (28) and (29) we get

$$
u_{\psi}^{q}(z) \leqq R_{0}^{-2}\left[2 u_{\psi}^{q}(\eta)+3 k\right] .
$$

$$
\begin{aligned}
& \int_{I} q(t) \varrho_{z}^{-1}(t) \mathrm{d}_{\operatorname{var}_{t}} \varrho_{z}(t)=\frac{1}{2} \int_{I} q(t) \varrho_{z}^{-2}(t) \mathrm{d}_{\operatorname{var}_{t}}\left[\varrho_{\eta}^{2}(t)-2 r \varrho_{\eta}(t) \cos \omega_{z}(t ; I)\right] \leqq \\
& \leqq \frac{1}{2} \int_{I} q(t) \varrho_{z}^{-2}(t){\mathrm{d} v \operatorname{var}_{t} \varrho_{\eta}^{2}(t)+r \int_{I} q(t) \varrho_{z}^{-2}(t) \mathrm{d} \operatorname{var}_{t}\left[\varrho_{\eta}(t) \cos \omega_{z}(t ; I)\right] \leqq} \\
& \leqq \frac{1}{2} R_{0}^{-2} \int_{I} q(t) \varrho_{\eta}^{-2}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}^{2}(t)+r \int_{I_{r}} q(t) \varrho_{z}^{-2}(t) \varrho_{\eta}(t) \mathrm{d} \operatorname{var} \omega_{z}(t ; I)+ \\
& +r \int_{I_{r}} q(t) \varrho_{z}^{-2}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}(t)+ \\
& +r \int_{I^{r}} q(t) \varrho_{z}^{-2}(t) \varrho_{\eta}(t) \mathrm{d} \operatorname{var} \omega_{z}(t ; I)+r \int_{I^{r}} q(t) \varrho_{z}^{-2}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}(t) \leqq \\
& \leqq R_{0}^{-2}\left[\int_{I} q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}(t)+r^{-1} \int_{I_{r}} q(t) \varrho_{\eta}(t) \mathrm{d} \operatorname{var} \omega_{z}(t ; I)+\right. \\
& +\int_{I_{r}} q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}(t)+r \int_{I^{r}} q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \omega_{z}(t ; I)+ \\
& \left.+\int_{I_{r}} q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}(t)\right]=R_{0}^{-2}\left[2 \int_{I} q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}(t)+\right. \\
& \left.+r^{-1} \int_{I_{r}} q(t) \varrho_{\eta}(t) \mathrm{d} \operatorname{var} \omega_{z}(t ; I)+r \int_{I_{r}} q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \omega_{z}(t ; I)\right] .
\end{aligned}
$$

Since $\eta$ is the only point in $\psi(\langle a, b\rangle)$ belonging to the closure of $S$, we have

$$
\inf \left\{\varrho_{z}(t) ;|z-\eta| \geqq R_{0}, t \in\langle a, b\rangle\right\}=\delta>0,
$$

whence it follows for any $z \in S$ with $|z-\eta| \geqq R_{0}$

$$
u_{\psi}^{q}(z)=\int_{a}^{b} q(t) \varrho_{z}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{z}(t) \leqq \delta^{-1} \int_{a}^{b} q(t) \mathrm{d} \operatorname{var} \psi(t) .
$$

We conclude that

$$
\sup _{z \in S} u_{\psi}^{q}(z) \leqq \max \left\{\delta^{-1} \int_{a}^{b} q \mathrm{~d} \operatorname{var} \psi, R_{0}^{-2}\left[2 u_{\psi}^{q}(\eta)+3 k\right]\right\} .
$$

13. Remark. The above proposition together with remark 10 form an alternative to the inequalities concerning so-called cyclic and radial variation as established in [2].

Now we are able to show that the converse of theorem 9 is also valid. We shall derive a more precise result.
14. Theorem. Let $S \subset R^{2} \backslash \psi(\langle a, b\rangle)$ be a set whose closure meets $\psi(\langle a, b\rangle)$ at a single point $\eta$. Let the contingent of $S \odot \eta$ at $\eta$ be disjoint from the contingent of $\psi(\langle a, b\rangle)$ at $\eta$ and assume (25). Let $\theta(r) \geqq 0$ be a continuous non-decreasing function of the variable $r \geqq 0, \theta$ 丰 0 . If $x \in R^{1}$ and $f$ is a continuous function on $\langle a, b\rangle$ satisfying

$$
\begin{equation*}
|f(t)-x|=O(\theta(|\psi(t)-\eta|) q(t)) \quad \text { as } \quad \psi(t) \rightarrow \eta, \tag{35}
\end{equation*}
$$

then for $z \in S$

$$
\begin{gather*}
p_{\psi} f(z)-x \log \frac{|\psi(b)-z|}{|\psi(a)-z|}-\int_{a}^{b} \frac{f(t)-x}{|\psi(t)-\eta|} \mathrm{d}_{t}|\psi(t)-\eta|=  \tag{36}\\
=O\left(|z-\eta| \int_{|z-\eta|}^{\infty} \theta(x) x^{-2} \mathrm{~d} x\right) \text { as } z \rightarrow \eta .
\end{gather*}
$$

If $f$ satisfies (35) with $O$ replaced by $o$, then the right-hand side in (36) can be replaced by

$$
\circ\left(|z-\eta| \int_{|z-\eta|}^{\infty} \theta(x) x^{-2} \mathrm{~d} x\right)+O(|z-\eta|) .
$$

Proof. For the sake of brevity, we put

$$
f_{\chi}(t)=f(t)-\chi, \quad t \in\langle a, b\rangle,
$$

and adopt the notation introduced in (26). Then

$$
\begin{equation*}
p_{\psi} f(z)=x \log \frac{|\psi(b)-z|}{|\psi(a)-z|}+\int_{a}^{b} f_{\chi}(t) \varrho_{z}^{-1}(t) \mathrm{d} \varrho_{z}(t), \quad z \in S . \tag{37}
\end{equation*}
$$

Fix now constants $\varepsilon>0, R_{0}>0$ such that, for $t \in\langle a, b\rangle$,

$$
\begin{equation*}
0<\varrho_{\eta}(t)<R_{0} \Rightarrow\left|f_{\varkappa}(t)\right| \leqq \varepsilon \theta\left(\varrho_{\eta}(t)\right) q(t) . \tag{38}
\end{equation*}
$$

Let $z \in S,|z-\eta|=r$, and put

$$
M_{r}=\left\{t \in\langle a, b\rangle ; 0<\varrho_{\eta}(t)<r\right\} .
$$

Then

$$
\begin{equation*}
\left|\int_{M_{r}} f_{\chi}(t) \varrho_{z}^{-1}(t) \mathrm{d} \varrho_{z}(t)\right| \leqq \varepsilon \theta(r) \int_{M_{r}} \varrho_{z}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{z}(t) \leqq \varepsilon \theta(r) u_{\psi}^{q}(z) \tag{39}
\end{equation*}
$$

by (19). Similarly, (11) implies

$$
\begin{equation*}
\left|\int_{M_{r}} f_{\varkappa}(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \varrho_{\eta}(t)\right| \leqq \varepsilon \theta(r) u_{\psi}^{q}(\eta) . \tag{40}
\end{equation*}
$$

Next consider the set

$$
L^{r}=M_{R_{0}} \backslash M_{r}=\left\{t \in\langle a, b\rangle ; r \leqq \varrho_{\eta}(t)<R_{0}\right\}
$$

We have

$$
\begin{equation*}
\int_{L^{r}} f_{\chi}(t) \varrho_{z}^{-1}(t) \mathrm{d} \varrho_{z}(t)-\int_{L^{r}} f_{\varkappa}(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \varrho_{\eta}(t)=A_{1}+A_{2}, \tag{41}
\end{equation*}
$$

where

$$
\begin{aligned}
& A_{1}=\int_{L^{r}} f_{\chi}(t)\left[\varrho_{z}^{-1}(t)-\varrho_{\eta}^{-1}(t)\right] \mathrm{d} \varrho_{z}(t), \\
& A_{2}=\int_{L^{r}} f_{\varkappa}(t) \varrho_{\eta}^{-1}(t) \mathrm{d}_{t}\left[\varrho_{z}(t)-\varrho_{\eta}(t)\right] .
\end{aligned}
$$

Defining

$$
\begin{equation*}
A=\int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{z}^{-1}(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{z}(t) \tag{42}
\end{equation*}
$$

we get

$$
\begin{equation*}
\left|A_{1}\right| \leqq \varepsilon \int_{\cdot_{L}} \theta\left(\varrho_{\eta}(t)\right) q(t)\left|\varrho_{z}(t)-\varrho_{\eta}(t)\right| \varrho_{z}^{-1}(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{z}(t) \leqq \varepsilon r A . \tag{43}
\end{equation*}
$$

Define the measure $v$ on Borel sets $B \subset\langle a, b\rangle$ by

$$
\begin{equation*}
v(B)=\int_{B \cap L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{z}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{z}(t), \tag{44}
\end{equation*}
$$

so that

$$
A=\int_{a}^{b} \varrho_{\eta}^{-1}(t) \mathrm{d} v(t) .
$$

Applying lemma 11 and noting that

$$
L_{u}^{r}=\left\{t \in L^{r} ; \varrho_{\eta}^{-1}(t)>u\right\}=\left\{t \in\langle a, b\rangle ; r \leqq \varrho_{\eta}(t)<u^{-1}\right\}
$$

equals $\emptyset$ or $L^{r}$ according as $u>r^{-1}$ or $u \leqq R_{0}^{-1}$, one easily obtains

$$
A=v\left(L^{r}\right) R_{0}^{-1}+\int_{R_{0}-1}^{r^{-1}} v\left(L_{u}^{r}\right) \mathrm{d} u .
$$

Introducing the variable $\tau=u^{-1}$ we have for $R_{0}^{-1}<u \leqq r^{-1}$

$$
L_{u}^{r}=M_{\tau} \backslash M_{r}
$$

and the last integral transforms into

$$
\int_{r}^{R_{0}} \tau^{-2} \nu\left(M_{\tau} \backslash M_{r}\right) \mathrm{d} \tau
$$

Using the estimates (compare (44), (19))

$$
\begin{gathered}
v\left(M_{\tau} \backslash M_{r}\right) \leqq v\left(M_{\tau}\right) \leqq \theta(\tau) u_{\psi}^{q}(z), \\
v\left(L^{r}\right) \leqq v\left(M_{R_{0}}\right) \leqq \theta\left(R_{0}\right) u_{\psi}^{q}(z),
\end{gathered}
$$

we get

$$
\begin{equation*}
A \leqq u_{\psi}^{q}(z)\left[\theta\left(R_{0}\right) R_{0}^{-1}+\int_{r}^{R_{0}} \theta(\tau) \tau^{-2} \mathrm{~d} \tau\right] \tag{45}
\end{equation*}
$$

Next consider

$$
\begin{gathered}
\left|A_{2}\right| \leqq \varepsilon \int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var}_{t} \frac{\varrho_{z}^{2}(t)-\varrho_{\eta}^{2}(t)}{\varrho_{z}(t)+\varrho_{\eta}(t)} \leqq \\
\leqq \varepsilon \int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t) \frac{\varrho_{z}^{2}(t)-\varrho_{\eta}^{2}(t)}{\left[\varrho_{z}(t)+\varrho_{\eta}(t)\right]^{2}} \mathrm{~d}_{\operatorname{var}_{t}\left[\varrho_{z}(t)+\varrho_{\eta}(t)\right]+}^{+\varepsilon \int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t)\left[\varrho_{z}(t)+\varrho_{\eta}(t)\right]^{-1} \mathrm{~d} \operatorname{var}_{t}\left[\varrho_{z}^{2}(t)-\varrho_{\eta}^{2}(t)\right] .}
\end{gathered}
$$

We have thus

$$
\begin{equation*}
\left|A_{2}\right| \leqq \varepsilon\left(r C_{1}+r C_{2}+C_{3}\right), \tag{46}
\end{equation*}
$$

where

$$
\begin{aligned}
C_{1} & =\int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t)\left[\varrho_{z}(t)+\varrho_{\eta}(t)\right]^{-1} \mathrm{~d} \operatorname{var} \varrho_{z}(t), \\
C_{2} & =\int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t)\left[\varrho_{z}(t)+\varrho_{\eta}(t)\right]^{-1} \mathrm{~d} \operatorname{var} \varrho_{\eta}(t), \\
C_{3} & =\int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t)\left[\varrho_{z}(t)+\varrho_{\eta}(t)\right]^{-1} \mathrm{~d} \operatorname{var}\left[\varrho_{z}^{2}(t)-\varrho_{\eta}^{2}(t)\right] .
\end{aligned}
$$

As before, we associate with each component $I$ of $\left\{t \in\langle a, b\rangle ; \varrho_{\eta}(t)>0\right\}=M_{\infty}$ a continuous argument $\omega_{z}(t ; I)$ of $[\psi(t)-\eta] / z-\eta$ on $I$. We may clearly assume that $R_{0}>0$ has been chosen small enough to guarantee (22) for $z \in S$ and $t \in I \in$ $\in \mathscr{S}_{\infty}(\eta)\left(=\right.$ the system of all components of $\left.M_{\infty}\right)$. Noting that, for $t \in L$,

$$
\varrho_{z}(t) \leqq r+\varrho_{\eta}(t) \leqq 2 \varrho_{\eta}(t)
$$

and assuming $|z-\eta|=r<R_{0}$ we infer from (32)

$$
\begin{equation*}
t \in L^{r} \Rightarrow \frac{1}{2} \leqq \varrho_{\eta}(t) \varrho_{z}^{-1}(t) \leqq R_{0}^{-1} \tag{47}
\end{equation*}
$$

This permits us to derive the following estimates

$$
\begin{align*}
& \left.C_{1} \leqq \int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t) \varrho_{z}^{-1}(t)\left(1+\frac{1}{2}\right)^{-1} \mathrm{~d} \text { var } \varrho_{z}(t) \leqq \frac{2}{3} A \quad \text { (compare }(45)\right),  \tag{48}\\
& C_{2} \leqq \int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-2}(t)\left(1+R_{0}\right)^{-1} \mathrm{~d} \operatorname{var} \varrho_{\eta}(t)=\left(1+R_{0}\right)^{-1} \bar{A}
\end{align*}
$$

where

$$
\begin{equation*}
\bar{A}=\int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-2}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}(t) \tag{49}
\end{equation*}
$$

Defining the measure $\bar{v}$ on Borel sets $B \subset\langle a, b\rangle$ by

$$
\bar{v}(B)=\int_{B \cap L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var} \varrho_{\eta}(t)
$$

and repeating the argument used above for the estimate of $A$ we obtain

$$
\begin{gathered}
\bar{A}=\int_{a}^{b} \varrho_{\eta}^{-1}(t) \mathrm{d} \bar{v}(t)=\bar{v}\left(L^{r}\right) R_{0}^{-1}+\int_{r}^{R_{0}} \tau^{-2} \bar{v}\left(M_{\tau} \backslash M_{r}\right) \mathrm{d} \tau, \\
\bar{v}\left(M_{\tau} \backslash M_{r}\right) \leqq \theta(\tau) u_{\psi}^{q}(\eta), \quad \bar{v}\left(L^{L}\right) \leqq \theta\left(R_{0}\right) u_{\psi}^{q}(\eta),
\end{gathered}
$$

whence

$$
\begin{equation*}
\bar{A} \leqq u_{\psi}^{q}(\eta)\left[\theta\left(R_{0}\right) R_{0}^{-1}+\int_{r}^{R_{0}} \tau^{-2} \theta(\tau) \mathrm{d} \tau\right], \tag{50}
\end{equation*}
$$

so that

$$
\begin{equation*}
C_{2} \leqq\left(1+R_{0}\right)^{-1} u_{\psi}^{q}(\eta)\left[\theta\left(R_{0}\right) R_{0}^{-1}+\int_{r}^{R_{0}} \tau^{-2} \theta(\tau) \mathrm{d} \tau\right] . \tag{51}
\end{equation*}
$$

Let now $J$ range over the system $\mathscr{S}_{R_{0}}(\eta)$ of all components of $M_{R_{0}}=\{t \in\langle a, b\rangle$; $\left.\varrho_{\eta}(t)<R_{0}\right\}$ and put

$$
J^{r}=L^{r} \cap J
$$

Each $J$ is contained in a uniquely determined $I_{J} \in \mathscr{S}_{\infty}(\eta)$ and we put $\omega_{z}(t ; J)=$ $=\omega_{z}\left(t ; I_{J}\right), t \in J$. Employing (47) and (34) we get by lemmas 2,3

$$
\begin{aligned}
& C_{3} \leqq \int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-2}(t)\left(1+R_{0}\right)^{-1} \mathrm{~d} \operatorname{var}_{t}\left[\varrho_{z}^{2}(t)-\varrho_{\eta}^{2}(t)\right] \leqq \\
& \leqq \sum_{J} 2 r\left[\int_{J r} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-2}(t) \mathrm{d}_{\operatorname{var}_{t}} \varrho_{\eta}(t)+\int_{J r} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var}_{t} \omega_{z}(t ; J)\right]= \\
& =2 r \int_{L^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-2}(t) \mathrm{d}_{\operatorname{var}_{t}} \varrho_{\eta}(t)+ \\
& +2 \sum_{J} r \int_{J r} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t) \mathrm{d}_{\operatorname{var}_{t}} \omega_{z}(t ; J) .
\end{aligned}
$$

Recalling (49) we may write

$$
\begin{equation*}
C_{3} \leqq 2 r \bar{A}+2 \sum_{J} r \int_{J^{r}} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t) \mathrm{d}_{\operatorname{var}_{t}} \omega_{z}(t ; J) \tag{52}
\end{equation*}
$$

Define now the measures $\mu, \mu_{1}$ on Borel sets $B \subset J$ by

$$
\begin{aligned}
\mu(B) & =\int_{B} \varrho_{\eta}(t) q(t){\mathrm{d} \operatorname{var}_{t} \omega_{z}(t ; J)} \\
\mu_{1}(B) & =\int_{B} \theta\left(\varrho_{\eta}(t)\right) \varrho_{\eta}(t) q(t) \mathrm{d}^{\operatorname{var}_{t}} \omega_{z}(t ; J)
\end{aligned}
$$

so that

$$
\begin{equation*}
\int_{J r} \theta\left(\varrho_{\eta}(t)\right) q(t) \varrho_{\eta}^{-1}(t) \mathrm{d} \operatorname{var}_{t} \omega_{z}(t ; J)=\int_{J r} \varrho_{\eta}^{-2}(t) \mathrm{d} \mu_{1}(t) . \tag{53}
\end{equation*}
$$

Put for $\tau>0$

$$
J_{\tau}^{r}=\left\{t \in J^{r} ; \varrho_{\eta}^{-2}(t)>\tau\right\}
$$

and observe that $J_{\tau}^{r}=\emptyset$ for $\tau>r^{-2}$ and

$$
J_{\tau}^{r}=\left\{t \in J ; r \leqq \varrho_{\eta}(t)<\tau^{-1 / 2}\right\} \quad \text { for } \quad 0<\tau \leqq r^{-2} .
$$

Hence

$$
\mu_{1}\left(J_{\tau}^{r}\right) \leqq \theta\left(\tau^{-1 / 2}\right) \mu\left(J_{\tau}^{r}\right)
$$

and we get by lemma 2

$$
\begin{equation*}
\int_{J^{r}} \varrho_{\eta}^{-2}(t) \mathrm{d} \mu_{1}(t)=\int_{0}^{r^{-2}} \mu_{1}\left(J_{\tau}^{r}\right) \mathrm{d} \tau \leqq \int_{0}^{r^{-2}} \theta\left(\tau^{-1 / 2}\right) \mu\left(J_{\tau}^{r}\right) \mathrm{d} \tau . \tag{54}
\end{equation*}
$$

Since $\omega_{z}(t ; J)$ differs only by an additive constant from a continuous argument $\vartheta_{\eta}(t ; J)$ of $\psi(t)-\eta$ on $J$ and

$$
J_{\tau}^{r} \subset\left\{t \in J ; \varrho_{\eta}(t)<x\right\}=J_{x} \quad \text { with } \quad x=\tau^{-1 / 2}
$$

we obtain by (27)

$$
\sum_{J} \mu\left(J_{\tau}^{r}\right) \leqq k \tau^{-1 / 2}
$$

(recall that

$$
\left.k=\sup _{u>0} u^{-1} v_{\psi \psi u}^{q}(\eta)\right),
$$

which together with (54), (53), (52) and (50) implies

$$
\begin{gathered}
C_{3} \leqq 2 r \bar{A}+2 k r \int_{0}^{r^{-2}} \theta\left(\tau^{-1 / 2}\right) \tau^{-1 / 2} \mathrm{~d} \tau \leqq \\
\leqq 2 r\left[u_{\psi}^{q}(\eta) \theta\left(R_{0}\right) R_{0}^{-1}+\left(u_{\psi}^{q}(\eta)+k\right) \int_{r}^{\infty} \theta(x) x^{-2} \mathrm{~d} x\right] .
\end{gathered}
$$

Combining this with (51), (48), (45), (46) and writing

$$
U=\sup _{z \in \bar{S}} u_{\psi}^{q}(z)
$$

(cf. proposition 12 and note also that $u_{\psi}^{q}(\eta) \leqq U$ in view of the lower-semicontinuity of $u_{\psi}^{q}(\ldots)$ established in lemma 7 ) we arrive at

$$
\left|A_{2}\right| \leqq \varepsilon r\left[4 U \theta\left(R_{0}\right) R_{0}^{-1}+(3 U+k) \int_{r}^{\infty} \theta(x) x^{-2} \mathrm{~d} x\right] .
$$

Adding this to the estimate (43) (see also (45)) we get by virtue of (41)

$$
\begin{gather*}
\left|\int_{L^{r}} f_{\varkappa} \varrho_{z}^{-1} \mathrm{~d} \varrho_{z}-\int_{L^{r}} f_{\chi} \varrho_{\eta}^{-1} \mathrm{~d} \varrho_{\eta}\right| \leqq  \tag{55}\\
\leqq \varepsilon r\left[5 U \theta\left(R_{0}\right) R_{0}^{-1}+(4 U+k) \int_{r}^{\infty} \theta(x) x^{-2} \mathrm{~d} x\right] .
\end{gather*}
$$

Finally, consider the set

$$
Z=\left\{t \in\langle a, b\rangle ; \varrho_{\eta}(t) \geqq R_{0}\right\}
$$

and note that

$$
\int_{z} f_{\chi} \varrho_{z}^{-1} \mathrm{~d} \varrho_{z}-\int_{z} f_{\chi} \varrho_{\eta}^{-1} \mathrm{~d} \varrho_{\eta}=\operatorname{Re} \int_{z} f_{\chi}(t)\left[\frac{1}{\psi(t)-z}-\frac{1}{\psi(t)-\eta}\right] \mathrm{d} \psi(t)
$$

Writting

$$
m=\sup \left\{\left|f_{\chi}(t)\right| ; t \in\langle a, b\rangle\right\}
$$

one concludes easily that

$$
\begin{equation*}
\left|\int_{Z} f_{\chi} \varrho_{Z}^{-1} \mathrm{~d} \varrho_{z}-\int_{Z} f_{\chi} \varrho_{\eta}^{-1} \mathrm{~d} \varrho_{\eta}\right| \leqq \frac{m r \operatorname{var} \psi(\langle a, b\rangle)}{R_{0}\left(R_{0}-r\right)} . \tag{56}
\end{equation*}
$$

Let $T=\left\{t \in\langle a, b\rangle ; \varrho_{\eta}(t)>0\right\}$. On account of (30)

$$
\int_{a}^{b} f_{\varkappa} \varrho_{z}^{-1} \mathrm{~d} \varrho_{z}=\int_{T} f_{\chi} \varrho_{z}^{-1} \mathrm{~d} \varrho_{z},
$$

which together with (37) yields for

$$
D(z)=p_{\psi} f(z)-\varkappa \log \frac{|\psi(b)-z|}{|\psi(a)-z|}-\int_{a}^{b} f_{\varkappa} \varrho_{\eta}^{-1} \mathrm{~d} \varrho_{\eta}
$$

the estimate

$$
\begin{aligned}
|D(z)| \leqq\left|\int_{M_{r}} f_{\chi} \varrho_{z}^{-1} \mathrm{~d} \varrho_{z}\right| & +\left|\int_{M_{r}} f_{\varkappa} \varrho_{\eta}^{-1} \mathrm{~d} \varrho_{\eta}\right|+\left|\int_{L^{r}} f_{\varkappa} \varrho_{z}^{-1} \mathrm{~d} \varrho_{z}-\int_{L^{r}} f_{\varkappa} \varrho_{\eta}^{-1} \mathrm{~d} \varrho_{\eta}\right|+ \\
& +\left|\int_{Z} f_{\chi} \varrho_{z}^{-1} \mathrm{~d} \varrho_{z}-\int_{Z} f_{\chi} \varrho_{\eta}^{-1} \mathrm{~d} \varrho_{\eta}\right| .
\end{aligned}
$$

Noting that

$$
\theta(r) r^{-1} \leqq \int_{r}^{\infty} \theta(x) x^{-2} \mathrm{~d} x
$$

we infer from (39), (40), (55), (56) for arbitrary $z \in S$ with $|z-\eta|=r<R_{0}$

$$
|D(z)| \leqq \varepsilon r(11 U+k) \int_{r}^{\infty} \theta(x) x^{-2} \mathrm{~d} x+m r \text { var } \psi(\langle a, b\rangle) R_{0}^{-1}\left(R_{0}-r\right)^{-1}
$$

which completes the proof, because $\varepsilon>0, R_{0}>0$ are arbitrary constants fulfilling (38) and (22).
15. Remark. Suppose now that $\psi$ is simple in the sense that for $t_{1}, t_{2} \in\langle a, b\rangle$

$$
0<\left|t_{1}-t_{2}\right|<b-a \Rightarrow \psi\left(t_{1}\right) \neq \psi\left(t_{2}\right)
$$

and put

$$
\psi(\langle a, b\rangle)=K
$$

(in the introduction, the same symbol is used to denote the oriented curve described by $\psi$ ).

Let $Q \geqq 0$ be a bounded lower-semicontinuous function on $K$ and put

$$
q(t)=Q(\psi(t)), \quad t \in\langle a, b\rangle
$$

Defining $U_{K}^{\varrho}(\varrho, \eta)$ as in the introduction we have for $\varrho \notin\{|\psi(b)-\eta|,|\psi(a)-\eta|\}$

$$
U_{K}^{Q}(\varrho, \eta)=u_{\psi}^{q}(\varrho, \eta),
$$

whence

$$
U_{\bar{K}}^{Q}(\eta)=u_{\psi}^{q}(\eta) .
$$

Similarly,

$$
V_{K r}^{Q}(\eta)=v_{\psi r}^{q}(\eta) .
$$

Now it is easy to see that theorems 14,9 imply the theorem stated in the introduction.

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[^0]:    ${ }^{1}$ ) cf. [11], chap. IX, § 2; let us recall that a half-line $H \subset R^{2}$ issuing at $\eta$ belongs to the contingent of $K \subset R^{2}$ at $\eta$ provided there are points $z_{n} \in K \backslash\{\eta\}(n=1,2, \ldots)$ tending to $\eta$ such that the half-lines $\left\{\eta+r\left(z_{n}-\eta\right) ; r \geqq 0\right\}$ converge (in the natural sense) to $H$.

