## Aplikace matematiky

## Jindřich Nečas; Ivan Hlaváček

Solution of Signorini's contact problem in the deformation theory of plasticity by secant modules method

Aplikace matematiky, Vol. 28 (1983), No. 3, 199-214

Persistent URL: http://dml.cz/dmlcz/104027

## Terms of use:

© Institute of Mathematics AS CR, 1983
Institute of Mathematics of the Czech Academy of Sciences provides access to digitized documents strictly for personal use. Each copy of any part of this document must contain these Terms of use.


This document has been digitized, optimized for electronic delivery and stamped with digital signature within the project DML-CZ: The Czech Digital Mathematics Library http://dml.cz

# SOLUTION OF SIGNORINI'S CONTACT PROBLEM IN THE DEFORMATION THEORY OF PLASTICITY BY SECANT MODULES METHOD 

Jindřich Nečas, Ivan Hlaváčék

(Received September 21, 1982)

A problem of unilateral contact between an elasto-plastic body and a rigid friction less foundation will be solved within the range of the so called deformation theory of plasticity [1], [5]. Thus the famous Signorini's problem in linear elasticity [6] is generalized to non-linear stress-strain relations. The weak solution is defined on the basis of a variational inequality, which in turn is equivalent to the minimum of the potential energy. Then the so-called secant modules (Kačanov) iterative method is introduced, each step of which corresponds to a classical Signorini's problem in elastostatics. Thus a finite element analysis of the latter is available [7].
On an abstract level, we prove the convergence of the secant modules method to the exact solution. Special effort is devoted to some cases when rigid admissible displacements exist.

## 1. INTRODUCTION

Let us consider a bounded domain $\Omega \subset R^{3}$ with a Lipschitz boundary $\partial \Omega$ and assume that

$$
\partial \Omega=\Gamma_{u} \cup \Gamma_{\tau} \cup \Gamma_{K} \cup \Gamma_{M},
$$

where $\Gamma_{u}, \Gamma_{\tau}, \Gamma_{K}$ are open subsets of $\partial \Omega, \Gamma_{K} \neq \emptyset$ and the surface measure of $\Gamma_{M}$ vanishes.

Let the elasto-plastic body, occupying the domain $\Omega$, be governed by the following Hencky-Mises stress-strain relations

$$
\begin{equation*}
\tau_{i j}=\left(k-\frac{2}{3} \mu(\gamma)\right) \delta_{i j} e_{l l}+2 \mu(\gamma) e_{i j}, \tag{1.1}
\end{equation*}
$$

where $k$ is a (constant) bulk modulus,

$$
\gamma(\mathbf{u}, \mathbf{v})=-\frac{2}{3} \vartheta(\mathbf{u}) \vartheta(\mathbf{v})+2 e_{i j}(\mathbf{u}) e_{i j}(\mathbf{v}),
$$

$$
\begin{aligned}
& \gamma(\mathbf{u}, \boldsymbol{u}) \stackrel{\mathrm{df}}{=} \gamma(\boldsymbol{u})=\gamma, \quad \vartheta(\boldsymbol{u})=\operatorname{div} \mathbf{u}, \\
& e_{i j}(\boldsymbol{u})=\frac{1}{2}\left(\frac{\partial u_{i}}{\partial x_{j}}+\frac{\partial u_{j}}{\partial x_{i}}\right)
\end{aligned}
$$

and a repeated index implies summation over the range 1,2,3. Assumptions on the function $\mu$ will be presented in Section 3.

Finally, let the functions $\boldsymbol{u}^{0} \in\left[W^{1,2}(\Omega)\right]^{3}, \boldsymbol{f} \in\left[L^{2}(\Omega)\right]^{3}$ and $\boldsymbol{g} \in\left[L^{2}\left(\Gamma_{\tau}\right)\right]^{3}$ be given.
We are seeking a solution of the non-linear system

$$
\begin{equation*}
-\frac{\partial}{\partial x_{i}}\left[\left(k-\frac{2}{3} \mu(\gamma)\right) \vartheta(\mathbf{u})\right]-2 \frac{\partial}{\partial x_{j}}\left[\mu(\gamma) e_{i j}(\mathbf{u})\right]=f_{i}, \quad i=1,2,3, \tag{1.2}
\end{equation*}
$$

in $\Omega$, such that

$$
\begin{gather*}
u=u^{0} \quad \text { on } \Gamma_{u}  \tag{1.3}\\
\tau_{i j} v_{j}=g_{i} \text { on } \Gamma_{\tau},
\end{gather*}
$$

where $\boldsymbol{v}$ denotes the unit outward normal to $\partial \Omega$. We denote $u_{v}=u_{i} v_{i}, T_{v}=\tau_{i j} v_{i} v_{j}$, $\left(T_{t}\right)_{i}=T_{i}-T_{v} v_{i}$, where $T_{i}=\tau_{i j} v_{j}$, and assume that

$$
\begin{equation*}
u_{v} \leqq 0, \quad T_{v} \leqq 0, \quad u_{v} T_{v}=0 \text { on } \Gamma_{K} . \tag{1.5}
\end{equation*}
$$

The solution of the problem (1.2) till (1.5) leads to minimizing the following functional of potential energy (cf. [1])

$$
\begin{equation*}
\mathscr{L}(\boldsymbol{u})=\frac{1}{2} k \int_{\Omega} \vartheta^{2}(\boldsymbol{u}) \mathrm{d} x+\frac{1}{2} \int_{\Omega}\left(\int_{0}^{\gamma(\boldsymbol{u})} \mu(t) \mathrm{d} t\right) \mathrm{d} x-\int_{\Omega} f_{i} u_{i} \mathrm{~d} x-\int_{\Gamma_{\tau}} g_{i} u_{i} \mathrm{~d} s \tag{1.6}
\end{equation*}
$$

oyer the convex set

$$
\begin{equation*}
K=\left\{\boldsymbol{u} \in\left[W^{1,2}(\Omega)\right]^{3} \mid \mathbf{u}=\mathbf{u}^{0} \text { on } \Gamma_{u}, u_{v} \leqq 0 \text { on } \Gamma_{K}\right\} . \tag{1.7}
\end{equation*}
$$

The latter problem is equivalent to the solution of the variational inequality:

$$
\begin{align*}
\mathbf{u} \in K, & \int_{\Omega}\left[\left(k-\frac{2}{3} \mu(\gamma)\right) \vartheta(\mathbf{u}) \vartheta(\mathbf{v}-\mathbf{u})+2 \mu(\gamma) e_{i j}(\mathbf{u}) e_{i j}(\mathbf{v}-\mathbf{u})-\right.  \tag{1.8}\\
& \left.-f_{i}\left(v_{i}-u_{i}\right)\right] \mathrm{d} x-\int_{\Gamma_{\tau}} g_{i}\left(v_{i}-u_{i}\right) \mathrm{d} s \geqq 0 \quad \forall \mathbf{v} \in K .
\end{align*}
$$

Method of secant modules (or Kačanov method, see [1] - chapter 8 and 11.5) consists in solving a sequence of the following variational inequalities:

$$
\begin{align*}
& \mathbf{u}_{n+1} \in K, \quad \int_{\Omega}\left[\left(k-\frac{2}{3} \mu\left(\gamma\left(\mathbf{u}_{n}\right)\right) \vartheta\left(\mathbf{u}_{n+1}\right) \vartheta\left(\mathbf{v}-\mathbf{u}_{n+1}\right)+\right.\right.  \tag{1.9}\\
+ & \left.2 \mu\left(\gamma\left(\mathbf{u}_{n}\right)\right) e_{i j}\left(\mathbf{u}_{n+1}\right) e_{i j}\left(\mathbf{v}-\mathbf{u}_{n+1}\right)-f_{i}\left(v_{i}-\left(u_{n+1}\right)_{i}\right)\right] \mathrm{d} x-
\end{align*}
$$

$$
-\int_{\Gamma_{\tau}} g_{i}\left(v_{i}-\left(u_{n+1}\right)_{i}\right) \mathrm{d} s \geqq 0, \quad n=1,2, \ldots
$$

Under certain assumptions on the function $\mu$ we shall prove convergence of the method. We use an abstract approach, parallel to that of [1]. The problem is transferred to the solution of a sequence of variational inequalities with variable coefficients, in general.

## 2. ABSTRACT FORMULATION

Let a functional $\Phi$ be given on a Hilbert space $H$. Assume that $\Phi$ has the second Gateaux differential $D^{2} \Phi(u, h, k)$ and the mapping $u \mapsto D^{2} \Phi(u, h, k)$ is continuous on every line segment.

Assume further that

$$
\begin{equation*}
D^{2} \Phi(u, h, h) \geqq m\|h\|^{2}, \quad m=\text { cost. }>0 \tag{2.1}
\end{equation*}
$$

Let a bilinear form $B(u ; x, y)$ be given, symmetric in $x, y$ and such that

$$
\begin{gather*}
B(u ; x, x) \geqq c_{1}\|x\|^{2}, \quad c_{1}=\text { const. }>0  \tag{2.2}\\
|B(u ; x, y)| \leqq c_{2}\|x\|\|y\|  \tag{2.3}\\
B(u ; u, v)=D \Phi(u, v),  \tag{2.4}\\
\frac{1}{2} B(x ; y, y)-\frac{1}{2} B(x ; x, x)-\Phi(y)+\Phi(x) \geqq 0 \quad \forall x, y \in H . \tag{2.5}
\end{gather*}
$$

Moreover, let $K$ be a closed convex subset of $H$.
Theorem 2.1. Let the assumptions (2.1) till (2.5) be satisfied and let an element $\varphi \in H$ be given.

Then the problem: find $u \in K$ such that

$$
\begin{equation*}
D \Phi(u, v-u) \geqq(\varphi, v-u) \quad \forall v \in K \tag{2.7}
\end{equation*}
$$

has a unique solution.
Let $u_{n} \in K, n=1,2, \ldots$, be such that

$$
\begin{equation*}
B\left(u_{n} ; u_{n+1}, v-u_{n+1}\right) \geqq\left(\varphi, v-u_{n+1}\right) \quad \forall v \in K . \tag{2.8}
\end{equation*}
$$

Then

$$
\lim _{n \rightarrow \infty} u_{n}=u
$$

holds in the space $H$.
Proof. The existence and uniqueness of a solution of the problem (2.7) is easy to verify.

Let us introduce the notation

$$
\begin{equation*}
\pi_{n}(v)=\Phi\left(u_{n}\right)-(\varphi, v)+\frac{1}{2} B\left(u_{n} ; v, v\right)-\frac{1}{2} B\left(u_{n} ; u_{n}, u_{n}\right) . \tag{2.9}
\end{equation*}
$$

By virtue of (2.5) we may write

$$
\begin{gather*}
\pi_{n}\left(u_{n+1}\right)=\Phi\left(u_{n}\right)-\left(\varphi, u_{n+1}\right)+\frac{1}{2} B\left(u_{n} ; u_{n+1}, u_{n+1}\right)- \\
-\frac{1}{2} B\left(u_{n} ; u_{n}, u_{n}\right)-\Phi\left(u_{n+1}\right)+\Phi\left(u_{n+1}\right) \geqq \\
\geqq \geqq\left(u_{n+1}\right)-\left(\varphi, u_{n+1}\right) \stackrel{\text { df }}{=} \psi\left(u_{n+1}\right) .
\end{gather*}
$$

We have defined

$$
\begin{equation*}
\psi(v)=\Phi(v)-(\varphi, v) . \tag{2.10}
\end{equation*}
$$

Using (2.8) we obtain

$$
\begin{equation*}
\frac{1}{2} B\left(u_{n} ; u_{n+1}, u_{n+1}\right)-\left(\varphi, u_{n+1}\right) \leqq \frac{1}{2} B\left(u_{n} ; v, v\right)-(\varphi, v) \quad \forall v \in K, \tag{2.11}
\end{equation*}
$$

consequently,

$$
\begin{equation*}
\pi_{n}\left(u_{n}\right)=\psi\left(u_{n}\right) \geqq \pi_{n}\left(u_{n+1}\right) . \tag{2.12}
\end{equation*}
$$

From (2.9') and (2.12) it follows that

$$
\begin{equation*}
\psi\left(u_{n}\right) \geqq \psi\left(u_{n+1}\right) \tag{2.13}
\end{equation*}
$$

Assumption (2.1) yields the coerciveness of $\psi$ :

$$
\begin{equation*}
\psi(v) \geqq c_{3}\|v\|^{2}-c_{4} \quad \forall v \in H \tag{2.14}
\end{equation*}
$$

Therefore, using (2.13) and (2.14) we obtain

$$
\lim _{n \rightarrow \infty} \psi\left(u_{n}\right)=c<-\infty .
$$

We have

$$
\begin{gather*}
c_{1}\left\|u_{n+1}-u_{n}\right\|^{2} \leqq B\left(u_{n} ; u_{n+1}-u_{n}, u_{n+1}-u_{n}\right)=  \tag{2.15}\\
=B\left(u_{n} ; u_{n}, u_{n}\right)+B\left(u_{n} ; u_{n+1}, u_{n+1}-2 u_{n}\right)
\end{gather*}
$$

on the other hand, we may write

$$
\begin{gather*}
2 \psi\left(u_{n}\right)-2 \pi_{n}\left(u_{n+1}\right) \geqq B\left(u_{n} ; u_{n}, u_{n}\right)-B\left(u_{n} ; u_{n+1}, u_{n+1}\right)+  \tag{2.16}\\
+2 B\left(u_{n} ; u_{n+1}-u_{n}\right),
\end{gather*}
$$

using (2.8). Therefore (2.16) and (2.15) yield that $2 \psi\left(u_{n}\right)-2 \pi_{n}\left(u_{n+1}\right) \geqq B\left(u_{n} ; u_{n}, u_{n}\right)+B\left(u_{n} ; u_{n+1}, u_{n+1}-2 u_{n}\right) \geqq c_{1}\left\|u_{n+1}-u_{n}\right\|^{2}$.
Using (2.12), (2.13), (2.9') and the convergence of $\psi\left(u_{n}\right)$, we obtain

$$
\lim _{n \rightarrow \infty}\left\|u_{n+1}-u_{n}\right\|=0
$$

Moreover, we have

$$
\begin{gather*}
\frac{1}{2} m\left\|u_{n}-u\right\|^{2} \leqq D \Phi\left(u_{n}, u_{n}-u\right)-D \Phi\left(u, u_{n}-u\right)=  \tag{2.17}\\
=B\left(u_{n} ; u_{n}, u_{n}-u\right)-D \Phi\left(u, u_{n}-u\right) \leqq B\left(u_{n} ; u_{n}, u_{n}-u\right)+\left(\varphi, u-u_{n}\right),
\end{gather*}
$$

by virtue of (2.7).
We also may write

$$
\begin{gather*}
B\left(u_{n} ; u_{n}, u_{n}-u\right)+\left(\varphi, u-u_{n}\right)=B\left(u_{n} ; u_{n}-u_{n+1}, u_{n}-u\right)+  \tag{2.18}\\
+B\left(u_{n} ; u_{n+1}, u_{n}-u\right)+\left(\varphi, u-u_{n}\right)=B\left(u_{n} ; u_{n}-u_{n+1}, u_{n}-u\right)+ \\
+B\left(u_{n} ; u_{n+1}, u_{n}-u_{n+1}\right)+\left(\varphi, u_{n+1}-u_{n}\right)+ \\
+B\left(u_{n} ; u_{n+1}, u_{n+1}-u\right)+\left(\varphi, u-u_{n+1}\right) \leqq \\
\leqq B\left(u_{n} ; u_{n}-u_{n+1}, u_{n}-u\right)+B\left(u_{n} ; u_{n+1}, u_{n}-u_{n+1}\right)+\left(\varphi, u_{n+1}-u_{n}\right)
\end{gather*}
$$

according to (2.8).
Using (2.18), (2.17), the boundedness of $u_{n}$, (2.3) and (2.16'), we obtain $u_{n} \rightarrow u$.
Q.E.D.

Moreover, let us consider the semi-coercive case, corresponding to the original problem and $\Gamma_{u}=\emptyset$.

Let $P G H$ be a subspace of $H$ such that $\operatorname{dim} P<\infty$. Let $H=P \oplus Q$ be the orthogonal decomposition and assume that

$$
\Phi(v), \quad D \Phi(v, h), \quad D^{2} \Phi(v, h, k) \text { and } B(u ; x, y)
$$

are independent of an addition of $p \in P$ in all variables: for example, $\Phi(v+p)=$ $=\Phi(v) \forall p \in P$, etc.

Assume that the only element $p \in P \cap K$ such that also $-p \in P \cap K$ is $p=0$. Let $\varphi \in H$ be such that

$$
\begin{equation*}
(\varphi, p)<0 \quad \forall p \in P \cap K \doteq\{0\} . \tag{2.19}
\end{equation*}
$$

Assume that $K$ is a closed convex cone with the vertex at the origin.
Lemma 2.1. Let the conditions (2.1), (2.2), (2.3) be fulfilled for elements $h, x$, $y \in Q$ and let (2.4), (2.19) hold.

Then the functionals $\psi(v)$ and

$$
\omega(x)=\frac{1}{2} B(v ; x, x)-(\varphi, x)
$$

are coercive, weakly lower semi-continuous in K. Consequently, solutions of the variational inequalities (2.7), (2.8) exist.
If e.g. $\bar{u}$ and $u$ are two solutions of (2.7), then $\bar{u}=u+p$, where $p \in P, u+p \in K$, $(\varphi, p)=0$. Each such $\bar{u}$ represents a solution of (2.7). A parallel assertion holds for solutions of (2.8).

Proof. To prove the coerciveness of $\psi$, it is sufficient to show that positive constants $c_{5}, c_{6}$ exist such that

$$
\begin{equation*}
\psi(v) \geqq c_{5}\|v\|-c_{6} \quad \forall v \in K . \tag{2.21}
\end{equation*}
$$

The latter inequality is equivalent to the following

$$
\begin{equation*}
\lim _{\substack{\|v\| \rightarrow \infty \\ v \in K}} \inf \frac{\psi(v)}{\|v\|} \geqq c_{5}>0 \tag{2.22}
\end{equation*}
$$

Assume that (2.22) is false. Then there exist $v_{n} \in K$ such that for $n \rightarrow \infty$

$$
\left\|v_{n}\right\| \rightarrow \infty, \quad \lim \frac{\psi\left(v_{n}\right)}{\left\|v_{n}\right\|}=c_{7} \leqq 0 .
$$

From (2.1) (for $h \in Q$ ) and (2.4) we obtain

$$
\begin{equation*}
\Phi(v) \geqq c_{8}\left\|\Pi_{Q} v\right\|^{2}-c_{9}, \tag{2.23}
\end{equation*}
$$

where $\Pi_{Q}$ stands for the projector of $H$ onto $Q$ and $c_{8}>0$. Consequently, we have

$$
\begin{equation*}
\psi(v) \geqq c_{8}\left\|\Pi_{Q} v\right\|^{2}-c_{9}-(\varphi, v) . \tag{2.24}
\end{equation*}
$$

Setting $v_{n}^{\prime}=v_{n}\| \| v_{n} \|$, we may write

$$
\frac{\psi\left(v_{n}\right)}{\left\|v_{n}\right\|} \geqq c_{8}\left\|v_{n}\right\|\left\|\Pi_{Q} v_{n}^{\prime}\right\|^{2}-\frac{c_{9}}{\left\|v_{n}\right\|}-\left(\varphi, v_{n}^{\prime}\right) .
$$

Therefore it must hold that

$$
\left\|\Pi_{Q} v_{n}^{\prime}\right\| \rightarrow 0
$$

We can assume that $v_{n}^{\prime} \rightarrow v^{\prime}$ in $H$ and therefore $v^{\prime} \in P \cap K,\left\|v^{\prime}\right\|=1$. We thus obtain

$$
\lim \frac{\psi\left(v_{n}\right)}{\left\|v_{n}\right\|} \geqq-\left(\varphi, v^{\prime}\right)>0
$$

in accordance with (2.19), which is a contradiction. Consequently, (2.21) is valid.
The rest of the existence proof is easy, since $\psi$ is convex in virtue of the assumptions on $\Phi$.

Let $u$ and $\bar{u}$ be two solutions of (2.7). Then

$$
\begin{gathered}
D \psi(u, \bar{u}-u) \geqq 0, \quad D \psi(\bar{u}, u-\bar{u}) \geqq 0, \\
m\left\|\Pi_{Q} \bar{u}-I I_{Q} u\right\|^{2} \leqq D \Phi(\bar{u}, \bar{u}-u)-D \Phi(u, \bar{u}-u) \leqq 0
\end{gathered}
$$

Consequently, $\bar{u}-u \in P$. Let us denote $\bar{u}-u=p$. We have

$$
D \Phi(u, p)-(\varphi, p) \geqq 0, \quad-D \Phi(u, p)+(\varphi, p) \geqq 0 ;
$$

since $D \Phi(u, p)=0$ by assumption, we are led to the conclusion that $(\varphi, p)=0$.

It is easy to verify that $u+p$ is a solution, provided $u$ is a solution and $p \in P$ fulfils $u+p \in K,(\varphi, p)=0$.

The analysis of the problem (2.8) could be carried out in a parallel way. Q.E.D.
Theorem 2.2. Let the assumptions of Lemma 2.1 be fulfilled. Moreover, let (2.5) hold and for any $h, k \in H$,

$$
\begin{equation*}
w_{n} \rightarrow w \Rightarrow B\left(w_{n} ; h, k\right) \rightarrow B(w ; h, k) . \tag{2.25}
\end{equation*}
$$

Let $K$ be a closed convex cone with the vertex at the origin, let $u_{n}$ and $u$ be as in Theorem 2.1.
Then

$$
\Pi_{Q} u_{n} \rightarrow \Pi_{Q} u
$$

and if $\lim _{k \rightarrow \infty} u_{n_{k}} \rightarrow v$, then $v$ is a solution of (2.7); we have $\left\|u_{n}\right\| \leqq c<\infty$.
Proof. As previously, we deduce that positive constants $c_{10}, c_{11}$ exist such that

$$
\begin{equation*}
\frac{1}{2} B(v ; w, w)-(\varphi, w) \geqq c_{10}\|w\|-c_{11}, \quad \forall w \in K \tag{2.26}
\end{equation*}
$$

holds uniformly with respect to $v$ and

$$
\begin{equation*}
\psi(w) \geqq c_{10}\|w\|-c_{11} . \tag{2.27}
\end{equation*}
$$

Hence Lemma 2.1 implies the existence of a sequence $\left\{u_{n}\right\}$. There exists a constant $c_{12}$ such that $\left\|u_{n}\right\| \leqq c_{12} \forall n$. Indeed, this is a consequence of (2.8), if we insert $v=0$ and use (2.2), (2.26).

Now the proof follows the same lines as the proof of Theorem 2.1 till (2.15), where we obtain

$$
\begin{equation*}
c_{1}\left\|\Pi_{Q} u_{n+1}-\Pi_{Q} u_{n}\right\|^{2} \leqq B\left(u_{n} ; u_{n}, u_{n}\right)+B\left(u_{n} ; u_{n+1}, u_{n+1}-2 u_{n}\right) \tag{2.28}
\end{equation*}
$$

Consequently,

$$
\left\|\Pi_{Q} u_{n+1}-\Pi_{Q} u_{n}\right\| \rightarrow 0 .
$$

Next, we may write

$$
\begin{gather*}
\frac{1}{2} m\left\|\Pi_{Q} u_{n}-\Pi_{Q} u\right\|^{2} \leqq D \Phi\left(u_{n}, u_{n}-u\right)-D \Phi\left(u, u_{n}-u\right)=  \tag{2.29}\\
=B\left(u_{n} ; u_{n}, u_{n}-u\right)-D \Phi\left(u, u_{n}-u\right)=D \Phi\left(u, u_{n+1}-u_{n}\right)+ \\
+B\left(u_{n} ; u_{n+1}, u_{n+1}-u\right)+B\left(u_{n} ; u_{n+1}, u_{n}-u_{n+1}\right)+ \\
+B\left(u_{n} ; u_{n}-u_{n+1}, u_{n}-u\right)-D \Phi\left(u, u_{n+1}-u\right) \leqq \\
\leqq D \Phi\left(u, u_{n+1}-u_{n}\right)+B\left(u_{n} ; u_{n+1}, u_{n+1}-u\right)+ \\
+B\left(u_{n} ; u_{n+1}, u_{n}-u_{n+1}\right)+B\left(u_{n} ; u_{n}-u_{n+1}, u_{n}-u\right)+
\end{gather*}
$$

$$
\begin{aligned}
+\left(\varphi, u-u_{n+1}\right) & \leqq D \Phi\left(u, u_{n+1}-u_{n}\right)+B\left(u_{n} ; u_{n+1}, u_{n}-u_{n+1}\right)+ \\
& +B\left(u_{n} ; u_{n}-u_{n+1}, u_{n}-u\right) \rightarrow 0
\end{aligned}
$$

where we have used the boundedness of $\left\|u_{n}\right\|$ and the convergence ( $2.28^{\prime}$ ).
Suppose now that a subsequence $u_{n_{k}} \rightarrow v$. Then we have for all $w \in K$

$$
B\left(u_{n_{k}-1} ; u_{n_{k}}, w-u_{n_{k}}\right) \geqq\left(\varphi, w-u_{n_{k}}\right),
$$

consequently,

$$
B\left(u_{n_{k}-1} ; v, w-v\right) \geqq(\varphi, w-v)+\varepsilon_{n_{k}}(w),
$$

where

$$
\varepsilon_{n_{k}}(w) \rightarrow 0 .
$$

By the assumption (2.25) and using (2.28') we obtain

$$
\begin{gathered}
B\left(u_{n_{k}-1} ; v, w-v\right)=B\left(u_{n_{k}}+u_{n_{k}-1}-u_{n_{k}} ; v, w-v\right)= \\
=B\left(u_{n_{k}}+\Pi_{Q} u_{n_{k}-1}-\Pi_{Q} u_{n_{k}} ; v, w-v\right) \rightarrow B(v ; v, w-v)= \\
=D \Phi(v, w-v) .
\end{gathered}
$$

Q.E.D.

## 3. APPLICATION TO AN ELASTO-PLASTIC BODY

We assume that the function $\mu$ is continuously differentiable in $[0, \infty)$ and satisfies the following conditions

$$
\begin{gather*}
0<\mu_{0} \leqq \mu(\gamma) \leqq \frac{3}{2} k,  \tag{3.1}\\
0<\alpha \leqq \mu(\gamma)+2 \gamma \frac{\mathrm{~d} \mu}{\mathrm{~d} \gamma}(\gamma) \leqq \beta<\infty . \tag{3.2}
\end{gather*}
$$

Then the inequalities (2.1) till (2.4) and (2.25) in the sense of Theorems 2.1 and 2.2 are fulfilled. For the details see [1] - chapter 8 and 11.5. Obviously, we put

$$
\begin{gathered}
B(\mathbf{v} ; \boldsymbol{w}, \mathbf{u})=\int_{\Omega}\left[\left(k-\frac{2}{3} \mu(\gamma(\mathbf{v}))\right) \vartheta(\mathbf{w}) \vartheta(\mathbf{u})+2 \mu(\gamma(\mathbf{v})) e_{i j}(\mathbf{w}) e_{i j}(\mathbf{u})\right] \mathrm{d} x, \\
\Phi=\mathscr{L}(\operatorname{see}(1.6)), \\
(\varphi, \mathbf{v})=\int_{\Omega} f_{i} v_{i} \mathrm{~d} x+\int_{\Gamma_{\tau}} g_{i} v_{i} \mathrm{~d} s, \\
P=\left\{\boldsymbol{p} \in\left[W^{1,2}(\Omega)\right]^{3} \mid e_{i j}(\boldsymbol{p})=0 \text { a.e. }\right\}=\left\{\boldsymbol{p}=\boldsymbol{a}+\boldsymbol{b} \times \mathbf{x}_{\}}\right\},
\end{gathered}
$$

where $\boldsymbol{a}, \boldsymbol{b} \in R^{3}$ are arbitrary constant vectors.
Let us recall a result from [1] - $11 \cdot 5$ (see also the references in [1]).

Theorem 3.1. If $\mathrm{d} \mu / \mathrm{d} \gamma \leqq 0$, then the condition (5.2) is fulfilled.
Proof. Setting

$$
M(\gamma)=\int_{0}^{\gamma} \mu(t) \mathrm{d} t,
$$

the condition (2.5) takes the following form

$$
\begin{equation*}
\int_{\Omega}[\mu(\gamma(\boldsymbol{u}))(\gamma(\boldsymbol{v})-\gamma(\boldsymbol{u}))-(M(\gamma(\boldsymbol{v}))-M(\gamma(\boldsymbol{u})))] \mathrm{d} x \geqq 0 . \tag{3.3}
\end{equation*}
$$

Consequently, (3.3) is satisfied if the function $M(\gamma)$ is concave.
Q.E.D.

Remark 4.1. The above conclusion can be verified also in two-dimensional problems of elastoplastic bodies.

## 4. SOME FURTHER SEMICOERCIVE CASES

In the present section we consider two-dimensional problems and the cases when $\Gamma_{u}=\emptyset$ but on a part $\Gamma_{0}$ of the boundary of the domain $\Omega \subset R^{2}$ the conditions of the so called bilateral contact, i.e.

$$
\begin{equation*}
u_{v}=0, \quad T_{t}=0 \quad \text { on } \quad \Gamma_{0} \tag{4.0}
\end{equation*}
$$

are prescribed. The latter conditions hold for example on the axis of symmetry.
Then the space of virtual rigid displacements has the dimension one and we can formulate uniquely solvable original and approximate contact problems. Besides, we shall prove that the solutions of the approximate problems (2.8) converge to the solution of the original problem (2.7).

First of all we study the cases when the whole problem can be solved in a subspace $Q \subset H$. We start again with an abstract analysis.

### 4.1. Solution of the Signorini problem in a subspace

Let $P \underset{\rightarrow}{G} H$ be a subspace of a Hilbert space $H, H=P \oplus Q$ the orthogonal decomposition of $H, \operatorname{dim} P<\infty$.

Assume that $\Phi(v), D \Phi(v, h), D^{2} \Phi(v, h, k)$ and $B(u ; x, y)$ are independent of an addition of $p \in P$ in all variables.

Let an element $\varphi \in Q$ and a convex cone $K$ with its vertex at the origin be given such that

$$
\begin{equation*}
P \subset K \tag{4.1}
\end{equation*}
$$

Lemma 4.1. Let the conditions (2.1), (2.2) (2.3) hold for $h, x, y \in Q$ and let (2.4) be satisfied.

Then the functionals

$$
\psi(v) \quad \text { and } \quad \omega(x)=\frac{1}{2} B(v ; x, x)-(\varphi, x)
$$

are coercive and weakly lower semicontinuous on $Q$.
There exists a unique solution $\hat{u} \in K \cap Q$ and $\hat{u}_{n+1} \in K \cap Q$ of the inequality

$$
D \Phi(\hat{u}, v-\hat{u}) \geqq(\varphi, v-\hat{u}) \quad \forall v \in K \cap Q
$$

and

$$
B\left(\hat{u}_{n} ; \hat{u}_{n+1}, v-\hat{u}_{n+1}\right) \geqq\left(\varphi, v-\hat{u}_{n+1}\right) \quad \forall v \in K \cap Q,
$$

respectively. Any solution of (2.7) and (2.8) can be written in the form

$$
u=\hat{u}+p \quad \text { and } \quad u_{n+1}=\hat{u}_{n+1}+p
$$

where $\hat{u}$ and $\hat{u}_{n+1}$ are the solutions of (2.7') and (2.8'), respectively, and $p \in P$. If $\hat{u}$ and $\hat{u}_{n+1}$ are solutions of (2.7') and (2.8), respectively, then $u=\hat{u}+p$ and $u_{n+1}=\hat{u}_{n+1}+p$, where $p$ is any element of $P$, represent solutions of (2.7) and (2.8), respectively.

Proof. From (2.1) it follows that $D^{2} \psi=D^{2} \Phi$ is positive definite on $Q$ and therefore $\psi$ is coercive on $Q . \psi$ is also strictly convex and differentiable, $K \cap Q$ convex and closed. Hence a unique solution of (2.7') exists.

Similar conclusions are valid for the functional $\omega(x)$, as follows from (2.2).
Since $\varphi \in Q$, we have

$$
\begin{equation*}
\psi(v)=\psi(v+p) \quad \forall p \in P \tag{4.2}
\end{equation*}
$$

The assumption (4.1) implies

$$
\begin{equation*}
K \cap Q=\Pi_{Q}(K) \tag{4.3}
\end{equation*}
$$

where $\Pi_{Q}$ denotes the projector of $H$ onto $Q$. In fact, let $v \in K$. Then

$$
\Pi_{Q} v=v-\Pi_{P} v=v+\left(-\Pi_{P} v\right) \in K
$$

consequently $\Pi_{Q}(K) \subset K \cap Q$. The converse inclusion is obvious:

$$
K \cap Q=\Pi_{Q}(K \cap Q) \subset \Pi_{Q}(K)
$$

Let $u$ be a solution of (2.7). Using (4.2), we may write

$$
\psi\left(\Pi_{Q} v\right)=\psi\left(\Pi_{Q} v+\Pi_{P} v\right)=\psi(v) \quad \forall v \in H
$$

Since $\Pi_{Q} u \in \Pi_{Q}(K)=K \cap Q$ and we have

$$
\psi\left(\Pi_{Q} u\right)=\psi(u) \leqq \psi(v)=\psi\left(\Pi_{Q} v\right) \quad \forall v \in K
$$

$\Pi_{Q} u=\hat{u}$ is a solution of $\left(2.7^{\prime}\right), u=\hat{u}+p, p \in P$.

In a parallel way we may prove that $\Pi_{Q} u_{n+1}=\hat{u}_{n+1}$ is a solution of $\left(2.8^{\prime}\right)$, hence $u_{n+1}=\hat{u}_{n+1}+p, p \in P$.

Let $\hat{u}$ be a solution of $\left(2.7^{\prime}\right)$. Then for $u=\hat{u}+p, p \in P$ we have

$$
\begin{equation*}
\psi(u)=\psi(\hat{u}) \leqq \psi(z) \quad \forall z \in K \cap Q \tag{4.4}
\end{equation*}
$$

Let $v \in K$. Then $\Pi_{Q} v \in \Pi_{Q}(K)=K \cap Q$ and

$$
\begin{equation*}
\psi\left(\Pi_{Q} v\right)=\psi(v) \tag{4.5}
\end{equation*}
$$

Combining (4.4) and (4.5) we obtain

$$
\psi(u) \leqq \psi(v) \quad \forall v \in K .
$$

By the assumption (4.1) $p \in P \subset K$, consequently $u=\hat{u}+p \in K$ and $u$ is a solution of (2.7).

The same argument is applicable to the functional $\omega$.
Theorem 4.1. Let the assumptions of Lemma 4.1 be fulfilled. Moreover, let (2.5) hold and let for all $h, k \in H$ the condition (2.25) be satisfied.

Denote by $\hat{u}$ and $\hat{u}_{n+1}$ the solutions of (2.7') and (2.8'), respectively. Then

$$
\lim _{n \rightarrow \infty} \hat{u}_{n}=\hat{u} .
$$

Proof. By the assumption (2.2) we have

$$
\frac{1}{2} B(v ; w, w)-(\varphi, w) \geqq \frac{1}{2} c_{1}\|w\|^{2}-c_{2}\|w\| \quad \forall v \in Q,
$$

with $c_{1}, c_{2}$ independent of $v$. Furthermore, we may write (by virtue of (2.1))

$$
\psi(v) \geqq c_{3}\|w\|^{2}-c_{4}\|w\| \quad \forall w \in Q .
$$

Lemma 4.1 implies existence of a sequence $\hat{u}_{n} \in Q \cap K$ and $c_{0}=$ const such that

$$
\left\|u_{n}\right\| \leqq c_{0} \quad \forall n .
$$

The proof then proceeds like that of Theorem 2.1 with the only change - the space $H$ is replaced everywhere by the subspace $Q$.

Application. Let $\Omega \subset R^{2}$ be a bounded domain with a Lipschitz boundary $\partial \Omega$ and let

$$
\partial \Omega=\Gamma_{0} \cup \Gamma_{\tau} \cup \Gamma_{K} \cup \Gamma_{M},
$$

where $\Gamma_{\mathrm{C}}$ and $\Gamma_{\mathrm{K}}$ have a positive length, whereas $\Gamma_{M}$ has zero length. Let the conditions (4.0) hold on $\Gamma_{0}$. We define

$$
K=\left\{\boldsymbol{u} \in\left[W^{1,2}(\Omega)\right]^{2} \mid u_{v}=0 \text { on } \Gamma_{0}, u_{v} \leqq 0 \text { on } \Gamma_{K}\right\}
$$

$$
\begin{gathered}
H=V=\left\{v \in\left[W^{1,2}(\Omega)\right]^{2} \mid v_{v}=0 \text { on } \Gamma_{0}\right\}, \\
\mathscr{R}=\left\{\mathbf{v} \in\left[W^{1,2}(\Omega)\right]^{2} \mid v_{1}=a_{1}-b x_{2}, v_{2}=a_{2}+b x_{1}\right\},
\end{gathered}
$$

where $a_{1}, a_{2}, b$ are arbitrary real constants;

$$
P=\{\boldsymbol{p} \in \mathscr{R} \cap K \mid-\boldsymbol{p} \in \mathscr{R} \cap K\}=\left\{\boldsymbol{p} \in \mathscr{R} \mid p_{\mathrm{v}}=0 \text { on } \Gamma_{0} \cup \Gamma_{K}\right\} .
$$

The same bilinear form $B$ and the functional $\psi$ will be chosen as in Section 3. Only the coefficient $(-2 / 3)$ has to be replaced by $(-1)$ and $3 k / 2$ in the formula (3.1) by $k$.

Obviously, the condition $P \subset K$ is fulfilled. Assume that $\Gamma_{\mathrm{C}}$ and $\Gamma_{K}$ consist of straight segments parallel with the $x_{1}$-axis. Then

$$
P=\left\{\boldsymbol{p}=\left(p_{1}, p_{2}\right) \mid p_{1}=a_{1}=\text { const }, p_{2}=0\right\} ;
$$

$\varphi \in Q$ if and only if

$$
V_{1} \equiv \int_{\Omega} f_{1} \mathrm{~d} x+\int_{\Gamma_{\tau}} g_{1} \mathrm{~d} s=0
$$

holds for the resultant of the external forces.
The space $V$ will be decomposed by means of some suitable inner product. We may choose for instance

$$
\begin{gathered}
(\mathbf{u}, \mathbf{v})_{v}=\int_{\Omega} e_{i j}(\mathbf{u}) e_{i j}(\mathbf{v}) \mathrm{d} x+p(\mathbf{u}) p(\mathbf{v}) \\
p(\mathbf{v})=\int_{\Gamma_{1}} v_{1} \mathrm{~d} s, \Gamma_{1} \subset \bar{\Omega}, \Gamma_{1} \text { has a positive length. }
\end{gathered}
$$

Then

$$
Q=V \ominus P=\{\mathbf{v} \in V \mid p(\mathbf{v})=0\}
$$

### 4.2. Solution of more general problems with unilateral contact

Lemma 4.2. Let the assumptions of Lemma 2.1 be satisfied. Moreover, let us assume that

$$
\begin{equation*}
(\varphi, p) \neq 0 \quad \forall p \in P \dot{-}\{0\} . \tag{4.5}
\end{equation*}
$$

Then there exist unique solutions of the problems (2.7) and (2.8).
Proof. Like at the beginning of the proof of Lemma 2.1 we derive for any two solutions $\bar{u}$ and $u$ of the inequality (2.7) that

$$
\bar{u}-u=p \in P, \quad(\varphi, p)=0
$$

By means of (4.5) we conclude that $p=0$ and there exists at most one solution. The argument for the inequality (2.8) is quite analogous.

The proof of coerciveness of $\psi$ and $\omega$ on $K$ follows the same lines as that of Lemma 2.1. Both functionals are convex and differentiable, hence they are weakly lower semicontinuous. Consequently, the solutions exist.

Theorem 4.2. Let the assumptions of Lemma 4.2. be fulfilled. Moreover, let (2.5) and (2.25) hold. Denote by $u$ and $u_{n+1}$ the solutions of the problem (2.7) and (2.8), respectively.

Then

$$
\lim _{n \rightarrow \infty} u_{n}=u
$$

Proof. Following the proof of Theorem 2.2, we arrive at the conclusion (cf. (2.29)) that

$$
\begin{equation*}
\left\|\Pi_{Q} u_{n}-\Pi_{Q} u\right\| \rightarrow 0, \quad n \rightarrow \infty . \tag{4.6}
\end{equation*}
$$

Besides, we derive the boundedness of norms $\left\|u_{n}\right\|$. Hence a subsequence $\left\{u_{m}\right\}$ exists such that

$$
\begin{equation*}
u_{m} \rightarrow u^{*} \quad(\text { weakly in } H), \quad m \rightarrow \infty \tag{4.7}
\end{equation*}
$$

Since $K$ is weakly closed, $u^{*} \in K$. It follows from (4.7) that

$$
\begin{equation*}
\Pi_{Q} u_{m} \rightarrow \Pi_{Q} u^{*} \tag{4.8}
\end{equation*}
$$

On the other hand, from (4.6) we obtain that

$$
\Pi_{Q} u_{m} \rightarrow \Pi_{Q} u
$$

consequently, $\Pi_{Q} u^{*}=\Pi_{Q} u$ and the convergence (4.8) is even strong. Moreover, by virtue of (4.7), we have

$$
\begin{equation*}
\Pi_{p} u_{m} \rightarrow \Pi_{F} u^{*} \tag{4.9}
\end{equation*}
$$

(the subspace $P$ being finite-dimensional).
Combining (4.8) and (4.9) we obtain the convergence

$$
\left\|u_{m}-u^{*}\right\| \rightarrow 0
$$

In the end of the proof of Theorem 2.2, however, we have shown that the limit element $u^{*}$ solves the inequality (2.7). The uniqueness of the solution implies $u^{*}=u$ and the whole sequence $\left\{u_{n}\right\}$ converges to $u$ strongly.

Application. The assumption (4.5) can be satisfied if and only if $\operatorname{dim}(\mathscr{R} \cap V)=1$. Indeed, let $H=V, K, \mathscr{R}, B, \psi$ be defined as above, $P=\mathscr{R} \cap V$. If $\Gamma_{0}=\emptyset$, then
$\mathscr{R}=\mathscr{R} \cap V, \operatorname{dim} \mathscr{R}=3$ and

$$
(\varphi, \mathbf{p})=a_{1} V_{1}+a_{2} V_{2}+b M, \quad \mathbf{p} \in \mathscr{R},
$$

where $V_{i}$ are the components of the external forces resultant and $M$ is the moment resultant. The condition (4.5) does not hold, since $(\varphi, \mathbf{p})=0$ for each vector $\left(a_{1}, a_{2}, b\right) \in R^{3}$ orthogonal to the vector $\left(V_{1}, V_{2}, M\right)$. The case $\operatorname{dim}(\mathscr{R} \cap V)=2$ is not possible.
Let $\Gamma_{9}$ consists of straight segments parallel with the $x_{1}$-axis. Then obviously

$$
\mathscr{R} \cap V=\left\{\boldsymbol{p} \in \mathscr{R} \mid p_{1}=a_{1}=\text { const., } p_{2}=0\right\}, \quad \operatorname{dim}(\mathscr{R} \cap V)=1 .
$$

The condition (4.5) is fulfilled if and only if the component $V_{1}$ of the force resultant does not vanish.

Next, let $\Gamma_{\mathrm{K}}$ be such that (see Fig. 1)


Fig. 1.

$$
P \cap K=\mathscr{R} \cap K=\left\{\boldsymbol{p} \in \mathscr{R} \mid p_{1}=a_{1} \leqq 0, p_{2}=0\right\} .
$$

Then the condition (2.19) is satisfied exactly if $V_{1}$ is positive.

Lemm 4.3. Let $P \neq\{0\}$, let $P$ be a subspace of $H$ as in Section 2, $\operatorname{dim} P<\infty$. Let (4.5) and

$$
\begin{equation*}
P \cap K=\{0\} \tag{4.10}
\end{equation*}
$$

hold. Assume that (2.1), (2.2), (2.3) hold for $h, x, y \in Q$ and (2.4) is valid.
Then there exist unique solutions of (2.7) and (2.8).
Proof. The assumption (2.1) yields

$$
\begin{equation*}
\Phi(v) \geqq c_{3}-c_{4}\|v\|+c_{5}\left\|\Pi_{Q} v\right\|^{2} \quad \forall v \in H \tag{4.11}
\end{equation*}
$$

On the other hand, we have

$$
\begin{equation*}
\left\|\Pi_{Q} v\right\| \geqq c_{6}\|v\| \quad \forall v \in K \tag{4.12}
\end{equation*}
$$

In fact, $\left\|\Pi_{Q} v\right\|$ is a seminorm in $H$ such that the assumptions of Theorem 2.2 in [4] are satisfied.

By combining (4.11) and (4.12) the coerciveness of $\psi$ on $K$ follows. Since $\psi$ is weakly lower semicontinuous, we obtain the existence of a solution of (2.7). The uniqueness is a consequence of (4.5) as in Lemma 4.2.

The argument for $\omega$ and (2.8) is analogous.
Theorem 4.3. Let the assumptions of Lemma 4.3 and (2.5), (2.25) be satisfied. Denote by $u$ and $u_{n+1}$ the solutions of (2.7) and (2.8), respectively.

Then

$$
\lim _{n \rightarrow \infty} u_{n}=u
$$

Proof is the same as that of Theorem 4.2.


Fig. 2.
Applications. Define $H=V, K, \mathscr{R}, B, \psi$ and $P=\mathscr{R} \cap V$ as above. Let $\Gamma_{0}$ consist of straight segments parallel with the $x_{1}$-axis. Then the condition (4.10) is satisfied, if $\Gamma_{K}$ has a proper form (see Fig. 2), i.e. if the component $v_{1}$ of the normal changes the sign. The condition (4.5) is again equivalent to $V_{1} \neq 0$.

## References

[I] J. Nečas, I. Hlaváček: Mathematical Theory of Elastic and Elasto-Plastic Bodies. Elsevier, Amsterdam 1981.
[2] J. Haslinger, I. Hlaváček: Contact between elastic bodies. Apl. mat. 25 (1980), 324-348, 26 (1981), 263-290, 321-344.
[3] I. Hlaváčék, J. Nečas: On inequalities of Korn's type. Arch. Ratl. Mech. Anal., 36 (1970), 305-334.
[4] J. Nečas: On regularity of solutions to nonlinear variational inequalities for second-order elliptic systems. Rend. di Matematica 2 (1975), vol. 8, Ser. VI., 481-498.
[5] L. M. Kačanov: Mechanika plastičeskich sred, Moskva 1948.
[6] G. Fichera: Boundary value problems of elasticity with unilateral constraints. In: S. Flüge (ed.): Encycl. of Physics, vol. VIa/2, Springer-Verlag, Berlin, 1972.
[7] I. Hlaváček, J. Lovíšek: A finite element analysis for the Signorini problem in plane elastostatics. Apl. mat. 22, (1977) 215-228, 25 (1980), 273-285.

## Souhrn

# ŘEŠENÍ SIGNORINIHO KONTAKTNÍHO PROBLÉMU <br> V DEFORMAČNÍ TEORII PLASTICITY METODOU SEČNÝCH MODULU゚ 

Jindřich Nečas, Ivan Hlaváček

Řeši se úloha jednostranného kontaktu mezi pružně plastickým tělesem a dokonale hladkou tuhou podporou v mezích tzv. deformační teorie plasticity. Řešení je formulováno pomocí variační nerovnice, ekvivalentní s principem minima potenciální energie. Metodou sečných modulů (Kačanova) je sestrojen iterační algoritmus, jehož každý krok odpovídá klasické Signoriniho úloze v teorii pružnosti. Dokazuje se konvergence této metody k přesnému řešení a studují se také některé úlohy, kdy existují přípustná pole posunutí tuhého tělesa.

Author's addresses: Doc. RNDr. Jindřich Nečas, DrSc., MFF UK, Malostranské nám. 25. 11800 Praha 1; Ing. Ivan Hlaváček, CSc., Matematický ústav ČSAV, Žitná 25, 11567 Praha 1,

