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# PARALLEL DISPLACEMENT OF VECTORS ON RHEONOMOUS ANHOLONOMIC MANIFOLD 

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We shall make use of the notations of [2]. Let us suppose that in the stationary space $r-L_{n}(t)$ there is defined the rheonomous anholonomic manifold $r-L_{n}^{m}(t)$ by means of the equations

$$
\begin{equation*}
B_{a}^{\alpha}=B_{a}^{\alpha}\left(x^{\omega}, t\right), \quad{ }_{p} n^{\alpha}={ }_{p} n^{\alpha}\left(x^{\omega}, t\right), \quad x^{\alpha}=x^{\alpha}\left(u^{A}, t\right) \tag{1}
\end{equation*}
$$

where $\left\{x^{\omega}, t\right\} \in \Omega,\left\{u^{A}, t\right\} \in \Lambda \times I($ see [2] , Equations(1,2) and (1,3)). Let the parametric equations

$$
\begin{equation*}
x^{\alpha}=x^{\alpha}(T), \quad t=T, \quad T \in J \tag{2}
\end{equation*}
$$

describe a trajectory which lies on the rheonomous manifold $r-L_{n}^{m}(t)$. Thus for every $T \in J$ there exist $m$ numbers $\mathrm{d} u^{a} / \mathrm{d} T$ such that the equation

$$
\begin{equation*}
\frac{\mathrm{d} x^{\alpha}}{\mathrm{d} T}=\frac{\mathrm{d} u^{a}}{\mathrm{~d} T} B_{a}^{\alpha}+B_{t}^{\alpha} \tag{3}
\end{equation*}
$$

holds.
Let $v^{a}(t), w^{a}(t)$ be the functions of the class $C_{2}$ defined on the interval $J$. We shall suppose that along the trajectory (2) a field of virtual or tangential vectors of the rheonomous manifold $r-L_{n}^{m}(t)$ is defined by means of functions

$$
\begin{equation*}
v^{\alpha}=v^{a} B_{a}^{\alpha}, \quad T \in J \tag{4}
\end{equation*}
$$

or

$$
\begin{equation*}
w^{\alpha}=w^{a} B_{a}^{\alpha}+B_{t}^{\alpha}, \quad T \in J, \tag{5}
\end{equation*}
$$

respectively. The vector field (4) or (5) will be called the $W$-parallel or T-parallel if for every $T \in J$

$$
\begin{equation*}
B_{\alpha}^{a} D_{T} v^{\alpha}=0 \tag{6}
\end{equation*}
$$

or

$$
\begin{equation*}
B_{\alpha}^{a} D_{r} w^{\alpha}=0, \tag{7}
\end{equation*}
$$

respectively. By means of formulae $(2,1)$ and $(2,2)$ of [2] we shall easily calculate that

$$
D_{T}\left(v^{a} B_{a}^{\alpha}\right)=B_{a}^{\alpha}\left(D_{T} v^{a}+w_{b}^{a} v^{b}\right)+{ }_{p} n^{\alpha}\left(h_{a b}^{p} v^{b} \frac{\mathrm{~d} u^{a}}{\mathrm{~d} T}+{ }^{p} m_{a} v^{a}\right)
$$

or

$$
\begin{gathered}
D_{T}\left(w^{a} B_{a}^{\alpha}+B_{t}^{x}\right)=B_{a}^{\alpha}\left(D_{T} w^{a}+w_{b}^{a} w^{b}+{ }^{\prime} w_{b}^{a} \frac{\mathrm{~d} u^{b}}{\mathrm{~d} T}+W^{a}\right)+ \\
\quad+{ }_{p} n^{\alpha}\left(h_{a b}^{p} w^{b} \frac{\mathrm{~d} u^{a}}{\mathrm{~d} T}+{ }^{p} m_{a} w^{a}+{ }^{p^{\prime}} m_{a} \frac{\mathrm{~d} u^{a}}{\mathrm{~d} T}+{ }^{p} W\right) .
\end{gathered}
$$

Hence and from (4) and (5) it follows that the equations (6) or (7) are equivalent to the equations

$$
\begin{equation*}
D_{T} v^{a}+w_{b}^{a} v^{b}=0 \tag{8}
\end{equation*}
$$

or

$$
\begin{equation*}
D_{T} w^{a}+w_{b}^{a} w^{b}+{ }^{\prime} w_{b}^{a} \frac{\mathrm{~d} u^{b}}{\mathrm{~d} T}+W^{a}=0 \tag{9}
\end{equation*}
$$

respectively.
If for every $T \in J$

$$
\begin{equation*}
D_{r} v^{a}=0 \tag{10}
\end{equation*}
$$

then the vector field (4) will be called the pseudoparallel field.
We have defined three "parallel" displacements, let us first consider the T-parallel displacement. We shall prove the following theorem:

Let $\left[x^{\alpha}\left({ }_{0} T\right),{ }_{0} T\right]$ be a point of the trajectory (2) and $\left({ }_{0} w^{a}{ }_{\circ} B_{a}^{\alpha}+{ }_{o} B_{t}^{\alpha}\right)$ the tangential vector of the rheonomous manifold $r-L_{n}^{m}(t)$ defined at that point. Then there exists along the trajectory (2) exactly one T-parallel field (5) such that $w^{a}\left({ }_{o} T\right)={ }_{o} w^{a}$.

The proof is easy. The functions $w^{a}(T)$ are solutions of the system of $m$-differential equations of (9). Writing in full the system (9) we obtain the equivalent linear system of in-differential equations of $m$-unknown functions in Cauchy's canonical form. From this result it follows immediately that the system (9) possesses on the interval $J$ exactly one solution for given initial conditions.

If the tangential vectors of a given trajectory of the rheonomous manifold $r-L_{n}^{m}(t)$ form the $T$-parallel field then we call such trajectory a $T$-geodesic. The following theorem holds:

Let $\left({ }_{0} w^{\alpha}\right)=\left({ }_{0} w^{a}{ }_{0} B_{a}^{\alpha}+{ }_{\circ} B_{t}^{\alpha}\right)$ be the tangential vector of the rheonomous manifold defined at its point $\left[{ }_{0} x^{\alpha},{ }_{\circ} T\right]$. Then there exists (locally) exactly one T-geodesic with
parametric description (2) such that

$$
\begin{equation*}
x^{\alpha}\left({ }_{\circ} T\right)={ }_{\circ} x^{\alpha}, \frac{\mathrm{d} x^{\alpha}\left({ }_{0} T\right)}{\mathrm{d} T}={ }_{\circ} v^{\alpha} \equiv{ }_{\circ} v^{a}{ }_{\circ} B_{a}^{\alpha}+{ }_{\circ} B_{t}^{\alpha} . \tag{11}
\end{equation*}
$$

Proof. By (9) the $W$-geodesic is described by the system of differential equations

$$
\begin{equation*}
D_{T} \frac{\mathrm{~d} u^{a}}{\mathrm{~d} T}+\left(w_{b}^{a}+{ }^{\prime} w_{b}^{a}\right) \frac{\mathrm{d} u^{a}}{\mathrm{~d} T}+W^{a}=0, \quad \frac{\mathrm{~d} x^{\alpha}}{\mathrm{d} T}=B_{a}^{\alpha} \frac{\mathrm{d} u^{a}}{\mathrm{~d} T}+B_{t}^{x} \tag{12}
\end{equation*}
$$

If we introduce the notation $\xi^{a}=\mathrm{d} u^{a} / \mathrm{d} T$ then (12) may be easely written in Cauchy's canonical form. For initial conditions (11), now of the form $x^{\alpha}\left({ }_{0} T\right)={ }_{0} x^{\alpha}, \xi^{a}\left({ }_{0} T\right)=$ $={ }_{0} w^{a}$, the new equivalent system has (locally) exactly one solution.

Preceding consideration of the $T$-parallel displacement may be extended also on the cases of the $W$-parallel and pseudoparallel displacement. We may show that under "usual" conditions the virtual vector may undergo a $W$-parallel or pseudoparallel displacement along the given trajectory in exactly one way. We may also define the notion of the $W$-geodesic or the pseudogeodesic and show its unique (local) existence for usual initial conditions. In case when $r-L_{n}^{m}(t)$ is a stationary manifold then $' w_{b}^{a}=w_{b}^{a}=0, W^{a}=0, B_{t}^{\alpha}=0$ and all three "parallel" displacement are identical.

If a metric field $\left(g_{\alpha \beta}\right)$ is defined in $L_{n}$ (everywhere symmetric and positively definite) then $L_{n}$ is a Riemannian space. We shall denote it $R_{n}$. In this case we shall denote $r-L_{n}(t)$ by $r-R_{n}(t)$. In the space $r-R_{n}(t)$ let be given such rheonomous manifold $r-L_{n}^{m}(t)$ that for every $\left\{x^{\omega}, t\right\} g_{\alpha \beta} B_{a p}^{\alpha} n^{\beta}=0, g_{\alpha \beta} n^{\alpha}{ }_{p} n^{\beta}={ }_{p q} \delta$. Then we shall denote $r-L_{n}^{m}(t)$ by $r-R_{n}^{m}(t)$. By means of the equations

$$
\begin{equation*}
g_{a b}=g_{\alpha \beta} B_{a}^{\chi} B_{b}^{\beta}, \quad g_{a b} g^{b c}=\delta_{a}^{c} \tag{13}
\end{equation*}
$$

the coordinates of the metric tensor field $\left(g_{a b}\right)$ are defined on the rheonomous manifold $r-R_{n}^{m}(t)$.

We shall show that the following equations hold at every point of the rheonomous manifold $r-R_{n}^{m}(t)$ :

$$
\begin{gather*}
D_{c} g_{b a}=0,  \tag{14}\\
\Gamma_{c b}{ }^{a}=\left\{\begin{array}{c}
a \\
c b
\end{array}\right\}+\Omega_{b c}{ }^{a}+\Omega_{c b}^{a}+\Omega_{b}{ }^{a}{ }_{c} \tag{15}
\end{gather*}
$$

where

$$
\begin{gather*}
\left\{\begin{array}{c}
a \\
c b
\end{array}\right\}=\frac{1}{2} g^{a e}\left(\partial_{c} g_{b e}+\partial_{b} g_{c e}-\partial_{e} g_{c b}\right), \quad \Omega_{b c}{ }^{a}=\partial_{[b} B_{|a|}^{\alpha} B_{c]}^{\alpha}, \\
D_{t} g_{\beta \alpha}=0 . \tag{16}
\end{gather*}
$$

I. From the equation

$$
\begin{equation*}
g_{\alpha \beta}=g_{\gamma \delta} \delta_{\alpha}^{\gamma} \delta_{\beta}^{\delta}=g_{a b} B_{\alpha}^{a} B_{\beta}^{b}+{ }_{p q} \delta^{p} n_{\alpha}^{q} n_{\beta} \tag{17}
\end{equation*}
$$

and the relation $\nabla_{\gamma} g_{\beta \alpha}=0$ it follows that

$$
D_{c} g_{b a}=B_{c}^{\gamma} B_{b}^{\beta} B_{a}^{\alpha} \nabla_{\gamma}\left(B_{\beta}^{e} B_{\alpha}^{f} g_{e f}\right)=B_{c}^{\gamma} B_{b}^{\beta} B_{a}^{\alpha} \nabla_{\gamma}\left(g_{\beta \alpha}-{ }_{p q} \delta^{p} n_{\beta}{ }^{q} n_{\alpha}\right)=0 .
$$

II. Obviously we may write

$$
\begin{gathered}
\Gamma_{c b}{ }^{a}=\frac{1}{2} B_{c}^{\gamma} B_{b}^{\beta} B_{\alpha}^{a}\left[\partial_{\gamma}\left(B_{\beta}^{e} B_{\omega}^{g} g_{e g}+{ }_{p q} \delta^{p} n_{\beta}{ }^{q} n_{\omega}\right)+\right. \\
\left.+\partial_{\beta}\left(B_{\gamma}^{f} B_{\omega}^{q} g_{f g}+{ }_{p q} \delta^{p} n_{\gamma}{ }^{q} n_{\omega \omega}\right)+\partial_{\omega}\left(B_{\gamma}^{f} B_{\beta}^{e} g_{f e}+{ }_{p q} \delta^{p} n_{\gamma}{ }^{q} n_{\beta}\right)\right] g^{\alpha \omega}-B_{b}^{\beta} \partial_{c} B_{\beta}^{a} .
\end{gathered}
$$

Applying a simple modification by means of (17) we obtain (15).
III. Let us denote by the symbol $\nabla_{t}$ the absolute derivate in the space $R_{n}$ along the curve which is described by the parametric equations $x^{\alpha}=x^{\alpha}\left({ }_{0} u^{A}, t\right), t \in I$. Evidently $D_{t} g_{\alpha \beta}=\nabla_{t} g_{\alpha \beta}=0$.

We shall show that the following theorem holds: Let $\left(u^{a}\right),\left(v^{b}\right)$ be two fields of $W$-parallel vectors which are defined along the trajectory (2). Then

$$
\begin{equation*}
\left.D_{T}\left(g_{a b} u^{a} v^{b}\right)=0^{1}\right) \tag{18}
\end{equation*}
$$

at every point of the trajectory (2).
In the proof of the theorem we shall make use of the notation $G_{a b}=D_{t} g_{a b}$. From $(13)$ and $(2,4)$ in [2] it follows that

$$
G_{a b}=g_{\alpha \beta}\left(D_{t} B_{a}^{\alpha}\right) B_{b}^{\beta}+g_{\alpha \beta} B_{a}^{\alpha} D_{t} B_{b}^{\beta}=w_{a b}+w_{b a} .
$$

Hence and from the equations

$$
D_{T} u^{a}+w_{b}^{a} u^{b}=0, \quad D_{T} v^{b}+w_{b}^{a} v^{b}=0, \quad D_{T} g_{a b}=G_{a b}
$$

we easily verify that (18) holds.
If $\left(u^{a}\right),\left(v^{b}\right)$ are two pseudoparallel fields of virtual vectors than $D_{T}\left(g_{a b} u^{a} v^{b}\right)=$ $=G_{a b} u^{a} v^{b}$ and for this season the scalar function $g_{a b} u^{a} v^{b}$ is generally not a constant. Similarly, for two $T$-parallel fields of virtual vectors $\left(u^{\alpha}\right),\left(v^{\beta}\right)$ the scalar function $g_{\alpha \beta} u^{\alpha} v^{\beta}$ is not in a general case a constant.

Let $\left(v^{a}\right)$ be a field of virtual vectors defined along the trajectory (2). If there exists the function $k=k(T), T \in J$ such that

$$
\begin{equation*}
D_{T}\left(k v^{a}\right)=0, \quad k \neq 0, \quad g_{a b} v^{a} v^{b}=\text { konst. }>0 \tag{19}
\end{equation*}
$$

holds for every $T \in J$, then we call $\left(v^{a}\right)$ a $\delta$-parallel field. The vector field $\left(v^{a}\right)$ is the $\delta$-parallel field if and only if

$$
\begin{equation*}
\delta_{T} v^{a} \equiv D_{1} v^{a}+\frac{1}{2} \frac{G_{b c} v^{b} v^{c}}{g_{b c} v^{b} v^{c}} v^{a}=0 \tag{20}
\end{equation*}
$$

holds for every $T \in J$.
The proof is analogical to the case of the rheonomous anholonomic manifold in [1],

[^0](20) and (22). It is easy to verify the unique existence of the parallel displacement and $\delta$-geodesic for the usual initial conditions. Let us remark that in entirely similar way as in [1] we may define the notion of $H$-parallel displacement on the anholonomic rheonomous manifold $r-R_{n}^{m}(t)$.

If the rheonomous manifold $r-R_{n}^{m}(t)$ is stationary than at every point $G_{a b}=0$ and the $\delta$-parallel displacement of vectors along the given trajectory is mutually identical with the $T$-parallel, $W$-parallel and pseudoparallel displacement.

## References

[1] Bruno Budinský: Parallel displacement of vectors in rheonomous Riemannian space. Cas. pro pěst. mat. 94 (1969), 34-42.
[2] Bruno Budinský: The Gauss and Gauss-Codazzi-Ricci equations for rheonomous anholonomic manifold. Čas. pro pěst. mat. 94 (1969), 270-276.
[3] J. A. Schouten, D. J. Struik: Einführung in die neueren Methoden der Differentialgeometrie I, II. Groningen Batavia 1935, 1938.

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[^0]:    ${ }^{1}$ ) I.e. the scalar product of two $W$-parallel "displaced" vectors is invariant.

