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## THE LAPLACE TRANSFORM OF ANALYTIC VECTOR-VALUED FUNCTIONS (REAL CONDITIONS)

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This paper is concerned with the problem of characteristic properties of the Laplace transform of vector-valued exponentially bounded functions on positive half-axis which are analytic in the sense that they are, roughly speaking, developable in power series at the points of positive half-axis with a linearly increasing radius of convergence. The properties can be described in a simple way: the functions in question are infinitely differentiable on the positive half-axis  $R^+$  and their derivatives satisfy the inequalities (B) (I) in Theorem 5 with certain constants  $M \ge 0$ ,  $\omega \ge 0$  and  $g \ge 0$ .

We give necessary and sufficient representability conditions of purely real (Widder's) type in terms of the behavior of derivatives of the Laplace image on the real half-axis as shown in Theorem 5.

The proof of Theorem 5 is mainly based on the representability theorems for Lipschitzian functions [1], [2]. In a reflexive space, the representability theorems from [3], [4] may be used, in the numerical case, the original representability theorem of Widder [5] is sufficient (see Remark 7).

The analyticity of semigroups is examined in Theorem 9.

It seems that the results presented are new even in the simplest, i.e. the numerical case.

1. In the sequel, R will denote the real number field and  $R^+$  the set of all positive numbers. If  $M_1, M_2$  are arbitrary sets, then  $M_1 \rightarrow M_2$  will denote the set of all mappings of the whole set  $M_1$  into the set  $M_2$ .

2. By E we denote a Banach space over R with the norm  $\|\cdot\|$ .

3. The notions of differentiability, measurability and integrability of functions with values in E are used in the strong (norm) sense.

4. Lemma.  $(p + 1)!/\lambda^{p+2} \leq p!/(\lambda - 1)^{p+1}$  for any  $\lambda > 1$  and  $p \in \{0, 1, ...\}$ .

5. Fundamental theorem (real form). Let  $M, \omega, \varrho$  be nonnegative constants and let  $F \in (\omega, \infty) \rightarrow E$ . Then the following two conditions (A) and (B) are equivalent:

(A) (I) the function F is infinitely differentiable in  $(0, \infty)$ ,

(II) 
$$\left\| \frac{\mathrm{d}^{p+q}}{\mathrm{d}\lambda^{p+q}} \left( \lambda^q F(\lambda) \right) \right\| \leq \frac{Mp! q! \varrho^q}{(\lambda - \omega)^{p+1}} \text{ for any } \lambda > \omega \text{ and } p, q \in \{0, 1, \ldots\},$$

(B) there exists an infinitely differentiable function  $f \in \mathbb{R}^+ \to E$  such that

(I) 
$$||f^{(q)}(t)|| \leq \frac{Me^{\omega t}q! \varrho^{q}}{t^{q}}$$
 for any  $t \in \mathbb{R}^{+}$  and  $q \in \{0, 1, ...,\}$   
(II)  $F(\lambda) = \int_{0}^{\infty} e^{-\lambda \tau} f(\tau) d\tau$  for any  $\lambda > \omega$ .

Proof. (A)  $\Rightarrow$  (B): Let us first denote

$$G_q(\lambda) = -\frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}} \left(\lambda^q F(\lambda)\right)$$

for any  $\lambda > \omega$  and  $q \in \{0, 1, \ldots\}$ .

It follows easily from (A) by means of Lemma 4 that

(1) 
$$\left\|\frac{\mathrm{d}^{p}}{\mathrm{d}\lambda^{p}}G_{q}(\lambda)\right\| = \left\|\frac{\mathrm{d}^{p+q+1}}{\mathrm{d}\lambda^{p+q+1}}\left(\lambda^{q}F(\lambda)\right)\right\| \leq \frac{M(p+1)! q! \varrho^{q}}{(\lambda-\omega)^{p+2}} \leq \frac{Mp! q! \varrho^{q}}{(\lambda-\omega-1)^{p+1}}$$

for any  $\lambda > \omega + 1$  and  $p, q \in \{0, 1, \ldots\}$ ,

(2) 
$$\left\|\frac{\mathrm{d}^{p}}{\mathrm{d}\lambda^{p}}\left(\lambda \ G_{q}(\lambda)\right)\right\| = \left\|\frac{\mathrm{d}^{p}}{\mathrm{d}\lambda^{p}}\left[\lambda \frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}}\left(\lambda^{q} F(\lambda)\right)\right]\right\| =$$
$$= \left\|\frac{\mathrm{d}^{p}}{\mathrm{d}\lambda^{p}}\left[\frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}}\left(\lambda^{q+1} F(\lambda)\right) - \left(q+1\right)\frac{\mathrm{d}^{q}}{\mathrm{d}\lambda^{q}}\left(\lambda^{q} F(\lambda)\right)\right]\right\| \leq$$
$$\leq \left\|\frac{\mathrm{d}^{p+q+1}}{\mathrm{d}\lambda^{p+q+1}}\left(\lambda^{q+1} F(\lambda)\right)\right\| + \left\|\left(q+1\right)\frac{\mathrm{d}^{p+q}}{\mathrm{d}\lambda^{p+q}}\left(\lambda^{q} F(\lambda)\right)\right\| \leq$$
$$\leq \frac{Mp! \left(q+1\right)! \varrho^{q+1}}{\left(\lambda-\omega\right)^{p+1}} + \left(q+1\right)\frac{Mp! \ q! \ \varrho^{q}}{\left(\lambda-\omega\right)^{p+1}} \leq$$
$$\leq \frac{Mp! \left(q+1\right)! \varrho^{q} (1+\varrho)}{\left(\lambda-\omega\right)^{p+1}} \leq \frac{Mp! \left(q+1\right)! \varrho^{q} (1+\varrho)}{\left(\lambda-\omega-1\right)^{p+1}}$$

for any  $\lambda > \omega + 1$  and  $p, q \in \{0, 1, ...\}$ . With regard to (1) and (2), we obtain from Theorem 4 in [2] (\*) that there exists

.

<sup>(\*)</sup> Let us remark that the property (4) serves only to deduce the properties (6) and (7). Hence only the properties (3), (5), (6), (7) are needed and used in the sequel. But these properties can be also obtained from Proposition 4.17 in [1] since it is easy to see from its proof that the function f has also the property that  $f(0_+)$  exists, though this fact is not explicitly stated.

a sequence  $\psi_q \in \mathbb{R}^+ \to E$ ,  $q \in \{0, 1, ...\}$  such that

(3) 
$$\|\psi_q(t)\| \leq Mq! \, \varrho^q e^{(\omega+1)t}$$
 for any  $t \in \mathbb{R}^+$  and  $q \in \{0, 1, \ldots\}$ ,

(4) 
$$\|\psi_q(t_1) - \psi_q(t_2)\| \leq M(q+1)! \, \varrho^q(1+\varrho) \int_{t_1}^{t_2} e^{(\omega+1)\tau} \, d\tau$$
  
for any  $t_1, t_2 \in \mathbb{R}^+$ ,  $t_1 < t_2$  and  $q \in \{0, 1, ...\}$ ,

(5) 
$$\int_0^\infty e^{-\lambda \tau} \psi_q(\tau) \, \mathrm{d}\tau = -\frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}} \left(\lambda^q F(\lambda)\right) \text{ for any } \lambda > \omega + 1 \text{ and } q \in \{0, 1, \ldots\}.$$

It follows from (4) that

(6) the function 
$$\psi_q$$
 is continuous on  $R^+$  for any  $q \in \{0, 1, ...\}$ ,

(7) 
$$\psi_q(0_+)$$
 exists for any  $q \in \{0, 1, \ldots\}$ .

Now we shall prove that

(8) 
$$\lambda \int_0^\infty e^{-\lambda \tau} \psi_q(\tau) \, \mathrm{d}\tau - \psi_q(0_+) \xrightarrow[\lambda \to \infty]{} 0 \quad \text{for any} \quad q \in \{0, 1, \ldots\}.$$

Indeed, let  $q \in \{0, 1, ...\}$  be fixed. By (3), we have

$$\begin{aligned} \left\|\lambda\int_{0}^{\infty} e^{-\lambda\tau}\psi_{q}(\tau) \,\mathrm{d}\tau - \psi_{q}(0_{+})\right\| &= \left\|\lambda\int_{0}^{\infty} e^{-\lambda\tau}(\psi_{q}(\tau) - \psi_{q}(0_{+})) \,\mathrm{d}\tau\right\| &= \\ &= \left\|\lambda\int_{0}^{T} e^{-\lambda\tau}(\psi_{q}(\tau) - \psi_{q}(0_{+})) \,\mathrm{d}\tau + \lambda\int_{T}^{\infty} e^{-\lambda\tau}(\psi_{q}(\tau) - \psi_{q}(0_{+})) \,\mathrm{d}\tau\right\| &\leq \\ &\leq \lambda\int_{0}^{T} e^{-\lambda\tau} \,\mathrm{d}\tau \sup_{0 < \tau < T} \left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\| + \lambda\int_{T}^{\infty} e^{-\lambda\tau}\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\| \,\mathrm{d}\tau \leq \\ &\leq \lambda\int_{0}^{\infty} e^{-\lambda\tau} \,\mathrm{d}\tau \sup_{0 < \tau < T} \left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\| + \lambda\int_{T}^{\infty} e^{-\lambda\tau}(2Mq! \,\varrho^{q}e^{(\omega+1)\tau}) \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda\int_{T}^{\infty} e^{-(\lambda-\omega-1)\tau} \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda\int_{T}^{\infty} e^{-(\lambda-\omega-1)\tau} \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda\int_{T}^{\infty} e^{-(\lambda-\omega-1)\tau} \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda\int_{T}^{\infty} e^{-(\lambda-\omega-1)\tau} \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda\int_{T}^{\infty} e^{-(\lambda-\omega-1)\tau} \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda\int_{T}^{\infty} e^{-(\lambda-\omega-1)\tau} \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda\int_{T}^{\infty} e^{-(\lambda-\omega-1)\tau} \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda\int_{T}^{\infty} e^{-(\lambda-\omega-1)\tau} \,\mathrm{d}\tau = \\ &= \sup_{0 < \tau < T} \left(\left\|\psi_{q}(\tau) - \psi_{q}(0_{+})\right\|\right) + 2Mq! \,\varrho^{q}\lambda = 0 \end{aligned}$$

Let now  $\varepsilon > 0$ . By (7), we choose T > 0 so small that  $\sup_{0 < \tau < T} \|\psi_q(\tau) - \psi_q(0_+)\| \leq \frac{1}{2}\varepsilon$ 

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and for this T, we choose  $\lambda > \omega + 1$  so large that

$$2Mq! \, \varrho^q \lambda \frac{e^{-(\lambda-\omega-1)T}}{\lambda-\omega-1} \leq \frac{\varepsilon}{2}.$$

Then we obtain immediately from the preceding formula that  $\|\lambda \int_0^\infty e^{-\lambda \tau} \psi_q(\tau) d\tau - -\psi_q(0_+)\| \leq \varepsilon$  for sufficiently large  $\lambda$  which proves (8).

On the other hand, by (A) we have

$$\left\|\frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}}\left(\lambda^{q} F(\lambda)\right)\right\| \leq \frac{Mq! \, \varrho^{q}}{(\lambda - \omega)^{2}}$$

and consequently

(9) 
$$\lambda \frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}} \left(\lambda^q F(\lambda)\right) \xrightarrow{\lambda \to \infty} 0 \text{ for any } q \in \{0, 1, \ldots\}.$$

It follows from (5), (8) and (9) that

(10) 
$$\psi_q(0_+) = 0 \text{ for any } q \in \{0, 1, ...\}.$$

Further, by (4) and (10) we find that

(11) 
$$\|\psi_q(t)\| \leq M(q+1)! \varrho^q(1+\varrho) t e^{(\omega+1)t}$$
 for any  $t \in \mathbb{R}^+$  and  $q \in \{0, 1, ...\}$ .  
Let us now take  $\varphi_q = (1/t) \psi_q(t)$  for any  $t \in \mathbb{R}^+$  and  $q \in \{0, 1, ...\}$ .

It follows from (6) and (11) that

(12) the function 
$$\varphi_q$$
 is continuous for any  $q \in \{0, 1, ...\}$ ,

(13) 
$$\|\varphi_q(t)\| \leq M(q+1)! \, \varrho^q(1+\varrho) \, e^{(\omega+1)t}$$
 for any  $t \in \mathbb{R}^+$  and  $q \in \{0, 1, \ldots\}$ .

Now (12) and (13) give easily

(14) 
$$\int_0^\infty e^{-\lambda \tau} \varphi_q(\tau) \, \mathrm{d}\tau \xrightarrow[\lambda \to \infty]{} 0 \quad \text{for any} \quad q \in \{0, 1, \ldots\}.$$

On the other hand, we get easily from (3) and (13) by means of Proposition 4.4 in [1] that

(15) 
$$\int_{0}^{\infty} e^{-\lambda \tau} \psi_{q}(\tau) d\tau = \int_{0}^{\infty} e^{-\lambda \tau} \tau \varphi_{q}(\tau) d\tau = -\frac{d}{d\lambda} \int_{0}^{\infty} e^{-\lambda \tau} \varphi_{q}(\tau) d\tau$$
for any  $\lambda > \omega + 1$  and  $q \in \{0, 1, ...\}$ .

By (5) and (15) we have

(16) 
$$\frac{\mathrm{d}}{\mathrm{d}\lambda}\int_0^\infty e^{-\lambda\tau}\,\varphi_q(\tau)\,\mathrm{d}\tau=\frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}}\,(\lambda^q\,F(\lambda))\,\,\text{for any }\,\lambda>\omega+1\,\,\text{and }\,q\in\{0,\,1,\,\ldots\}\,.$$

Using (9), (14) and (16) we conclude finally that

(17) 
$$\int_0^\infty e^{-\lambda \tau} \varphi_q(\tau) \, \mathrm{d}\tau = \frac{\mathrm{d}^q}{\mathrm{d}\lambda^q} \left(\lambda^q F(\lambda)\right) \quad \text{for any} \quad \lambda > \omega + 1 \quad \text{and} \quad q \in \{0, 1, \ldots\} \; .$$

Using Propositions 4.4, 4.6 and 4.12 in [1] we get from (A) and from (17) that

(18) 
$$\|\varphi_q(t)\| \leq Me^{\omega t}q! \varrho^q$$
 for any  $t \in \mathbb{R}^+$  and  $q \in \{0, 1, \ldots\}$ .

Now let us denote

 $(19) f = \varphi_0 .$ 

We shall prove that

(20) . the function f is q-times differentiable for any  $q \in \{0, 1, ...\}$ ,

(21) 
$$t^q f^{(q)}(t) = (-1)^q \varphi_q(t)$$
 for any  $t \in \mathbb{R}^+$  and  $q \in \{0, 1, ...\}$ .

To this goal we proceed by induction on q.

Clearly, for q = 0, the statements (20) and (21) are true by (19).

Now we shall suppose the validity of (20), (21) for a fixed  $q \in \{0, 1, ...\}$  and prove it for q + 1.

According to the induction hypothesis just made, and according to (12), (13) and (17) we can establish the following identities for any  $\lambda > \omega + 1$ :

$$(22) \quad \int_{0}^{\infty} e^{-\lambda \tau} \tau^{q+1} f^{(q)}(\tau) \, \mathrm{d}\tau = \int_{0}^{\infty} e^{-\lambda \tau} \tau(\tau^{q} f^{(q)}(\tau)) \, \mathrm{d}\tau = -\frac{\mathrm{d}}{\mathrm{d}\lambda} \int_{0}^{\infty} e^{-\lambda \tau} \tau^{q} f^{(q)}(\tau) \, \mathrm{d}\tau = \\ = (-1)^{q+1} \frac{\mathrm{d}}{\mathrm{d}\lambda} \int_{0}^{\infty} e^{-\lambda \tau} \varphi_{q}(\tau) \, \mathrm{d}\tau = (-1)^{q+1} \frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}} \left(\lambda^{q} F(\lambda)\right), \\ (23) \qquad \int_{0}^{\infty} e^{-\lambda \tau} \left(\int_{0}^{\tau} \sigma^{q} f^{(q)}(\sigma) \, \mathrm{d}\sigma\right) \mathrm{d}\tau = (-1)^{q} \int_{0}^{\infty} e^{-\lambda \tau} \left(\int_{0}^{\tau} \varphi_{q}(\sigma) \, \mathrm{d}\sigma\right) \mathrm{d}\tau = \\ = \frac{(-1)^{q}}{\lambda} \int_{0}^{\infty} e^{-\lambda \tau} \varphi_{q}(\tau) \, \mathrm{d}\tau = \frac{(-1)^{q}}{\lambda} \frac{\mathrm{d}^{q}}{\mathrm{d}\lambda^{q}} \left(\lambda^{q} F(\lambda)\right), \\ (24) \qquad \int_{0}^{\infty} e^{-\lambda \tau} \int_{0}^{\tau} \varphi_{q+1}(\sigma) \, \mathrm{d}\sigma \, \mathrm{d}\tau = \frac{1}{\lambda} \int_{0}^{\infty} e^{-\lambda \tau} \varphi_{q+1}(\tau) \, \mathrm{d}\tau = \frac{1}{\lambda} \frac{\mathrm{d}^{q+1}}{\mathrm{d}\lambda^{q+1}} \left(\lambda^{q+1} F(\lambda)\right). \end{aligned}$$

Further, a simple calculation shows that

(25) 
$$\frac{1}{\lambda} \frac{d^{q+1}}{d\lambda^{q+1}} (\lambda^{q+1} F(\lambda)) = \frac{1}{\lambda} \frac{d^{q+1}}{d\lambda^{q+1}} (\lambda\lambda^{q} F(\lambda)) =$$
$$= \frac{1}{\lambda} \left[ \lambda \frac{d^{q+1}}{d\lambda^{q+1}} (\lambda^{q} F(\lambda)) + (q+1) \frac{d^{q}}{d\lambda^{q}} (\lambda^{q} F(\lambda)) \right] =$$
$$= \frac{d^{q+1}}{d\lambda^{q+1}} (\lambda^{q} F(\lambda)) + \frac{q+1}{\lambda} \frac{d^{q}}{d\lambda^{q}} (\lambda^{q} F(\lambda)) \text{ for any } \lambda > \omega + 1.$$

Using now the uniqueness theorem for the Laplace transform, see 4.11 in [1], we get from (22)-(25) that

(26) 
$$t^{q+1} f^{(q)}(t) = (q+1) \int_0^t \tau^q f^{(q)}(\tau) d\tau + (-1)^{q+1} \int_0^t \varphi_{q+1}(\tau) d\tau$$
 for any  $t \in \mathbb{R}^+$ .

Now (26) implies that

(27) the function 
$$t^{q+1}f^{q}(t)$$
,  $t \in \mathbb{R}^{+}$ , is differentiable,

(28) 
$$(t^{q+1}f^{(q)}(t))' = (q+1)t^q f^{(q)}(t) + (-1)^{q+1}\varphi_{q+1}(t)$$
 for all  $t \in \mathbb{R}^+$ .

But we see from (27) that the function  $f^{(q)}$  itself is differentiable and consequently

(29) 
$$f$$
 is  $(q + 1)$ -times differentiable

On the other hand, (28) and (29) imply

(30) 
$$t^{q+1} f^{(q+1)}(t) = (-1)^{q+1} \varphi_{q+1}(t)$$
 for every  $t \in \mathbb{R}^+$ .

By (29) and (30), the induction step for the properties (20) and (21) is verified and consequently (20) and (21) hold for  $q \in \{0, 1, 2, ...\}$ .

The desired property (B) follows from (18), (20) and (21).

(B)  $\Rightarrow$  (A): Let f be a fixed function satisfying the condition (B). It follows at once from this property that

(1) the function F is infinitely differentiable on  $(\omega, \infty)$ .

Now we will prove by induction on q that

(2) 
$$\int_0^\infty e^{-\lambda \tau} \tau^q f^{(q)}(\tau) d\tau = (-1)^q \frac{d^q}{d\lambda^q} (\lambda^q F(\lambda)) \text{ for any } \lambda > \omega \text{ and } q \in \{0, 1, \ldots\}.$$

The identity (2) is evidently true by the assumption (B) for q = 0.

Now let us suppose its validity for a fixed  $q \in \{0, 1, ...\}$  and proceed to prove it for q + 1. We have clearly  $(t^{q+1} f^{(q)}(t))' = t^{q+1} f^{(q+1)}(t) + (q+1) t^q f^{(q)}(t)$  for any  $t \in \mathbb{R}^+$ .

Using this identity, we can write by (1), (2) and (B) that

$$\int_{0}^{\infty} e^{-\lambda \tau} \tau^{q+1} f^{(q+1)}(\tau) d\tau = -(q+1) \int_{0}^{\infty} e^{-\lambda \tau} \tau^{q} f^{(q)}(\tau) d\tau + \int_{0}^{\infty} e^{-\lambda \tau} \frac{d}{d\tau} (\tau^{q+1} f^{(q)}(\tau)) d\tau =$$
$$= -(q+1) \int_{0}^{\infty} e^{-\lambda \tau} \tau^{q} f^{(q)}(\tau) d\tau + \lambda \int e^{-\lambda \tau} \tau^{q+1} f^{(q)}(\tau) d\tau =$$
$$= (-1)^{q+1} (q+1) \frac{d^{q}}{d\lambda^{q}} (\lambda^{q} F(\lambda)) - \lambda \frac{d}{d\lambda} \int_{0}^{\infty} e^{-\lambda \tau} \tau^{q} f^{(q)}(\tau) d\tau =$$

$$= (-1)^{q+1} (q+1) \frac{d^{q}}{d\lambda^{q}} (\lambda^{q} F(\lambda)) - (-1)^{q} \lambda \frac{d^{q+1}}{d\lambda^{q+1}} (\lambda^{q} F(\lambda)) =$$
  
=  $(-1)^{q+1} \left[ (q+1) \frac{d^{q}}{d\lambda^{q}} (\lambda^{q} F(\lambda)) + \lambda \frac{d^{q+1}}{d\lambda^{q+1}} (\lambda^{q} F(\lambda)) \right] = (-1)^{q+1} \frac{d^{q+1}}{d\lambda^{q+1}} (\lambda^{q+1} F(\lambda))$ 

which verifies the induction step.

Now, using (2), we obtain from (B) that

(3) 
$$\left\|\frac{\mathrm{d}^{p+q}}{\mathrm{d}\lambda^{p+q}}\left(\lambda^{q} F(\lambda)\right)\right\| = \left\|\frac{\mathrm{d}^{p}}{\mathrm{d}\lambda^{p}} \int_{0}^{\infty} e^{-\lambda\tau} (-\tau)^{q} f^{(q)} \,\mathrm{d}\tau\right\| = \left\|\int_{0}^{\infty} e^{-\lambda\tau} (-\tau)^{p+q} f^{(q)}(\tau) \,\mathrm{d}\tau\right\| \leq \int_{0}^{\infty} e^{-\lambda\tau} \tau^{p} M e^{\omega\tau} q! \,\varrho^{q} = \frac{Mp! \,q! \,\varrho^{q}}{(\lambda - \omega)^{p+1}}$$
for any  $t \in \mathbb{R}^{+}$  and  $p, q \in \{0, 1, \ldots\}$ .

But (1) and (3) give (A).

6. Remark. The preceding theorem has the advantage that the chatacteristic growth constants  $(M, \omega, \varrho)$  are preserved in the course of the transformation in both directions. Nevertheless, technically it is difficult to verify the determining property (A).

7. Remark. Some words about the proof of the implication  $(A) \Rightarrow (B)$  in the preceding theorem. The first part of this proof deals with the existence of a sequence of functions  $\varphi_q \in R^+ \to E$  with the properties (12), (17) and (18). In the case of a reflexive space *E*, these properties can be obtained directly from Corollary 9 in [3] or from Theorem 4 in [4], in the numerical case, E = R, from Theorems 16a and 16b in [5], if we replace in (12) "continuous" by "measurable". In the remaining part of the proof we must replace "differentiable" by "feebly differentiable" in the following sense: a function  $f \in R^+ \to E$  is feebly differentiable if there exists a function  $g \in$  $\in R^+ \to E$  such that for any  $0 < \alpha < \beta$ , g is integrable over  $(\alpha, \beta)$  and  $f(\beta) - f(\alpha) =$  $= \int_{\alpha}^{\beta} g(\sigma) d\sigma$ . Moreover, some identities are valid not for all  $t \in R^+$  but only for almost all  $t \in R^+$ . All these technicalities do not influence the final result.

8. Remark. The infinitely differentiable functions satisfying the inequality

(1) 
$$||f^{(q)}(t)|| \leq M e^{\omega t} \frac{q! \varrho^q}{t^q}$$

for every  $t \in \mathbb{R}^+$  and  $q \in \{0, 1, ...\}$  with some nonnegative constants M,  $\omega$ ,  $\varrho$  represent a special class of functions analytic on the positive half-axis. It is clear from Taylor's theorem that (1) implies

(2) 
$$f(\tau) = \sum_{q=0}^{\infty} \frac{f^{(q)}(t)}{q!} (\tau - t)^{q}$$

for every  $t, \tau \in R^+$  such that  $|\tau - t| < t/\rho$  (supposing  $\rho > 0$  because the case  $\rho = 0$  is uninteresting). We see from (2) that the radius of convergence of the power series in (2) depends linearly on the value  $t \in R^+$  (more precisely, it is not less than a positive linear function of t). This is a remarkable property of the class of functions considered but not characteristic because it does not imply the exponential growth of derivatives.

9. A function  $\mathscr{T}$  defined on  $\mathbb{R}^+$  with values in the Banach space of bounded linear operators on E into E will be called a semigroup of linear operators in E if

- $(T_1) \mathcal{T}(t_1 + t_2) = \mathcal{T}(t_1) \mathcal{T}(t_2)$  for any  $t_1, t_2 \in \mathbb{R}^+$ ,
- $(T_2)$   $\mathcal{T}(t) x \to x (t \to 0_+)$  for any x from a dense subset of E.

10. Theorem. Let  $\mathcal{F}$  be a semigroup of linear operators in E. Then the following two statements (A), (B) are equivalent:

(A) there exist  $M \ge 0$ ,  $\omega \ge 0$  and  $\varrho \ge 0$  so that for every  $x \in E$  for which the function  $\mathcal{T}(\cdot) x$  is infinitely differentiable on  $\mathbb{R}^+$ , we have

$$\left\|\frac{\mathrm{d}^{q}}{\mathrm{d}t^{q}}\,\mathscr{T}(t)\,x\right\| \leq Me^{\omega t}\,\frac{q\,!\,\varrho^{q}}{t^{q}}\,\|x\|$$

for every  $t \in R^+$  and  $q \in \{0, 1, ...\};$ 

(B) there exist  $\delta > 0$  and  $K \ge 0$  so that  $\|\mathscr{T}(t)\| \le K$  for every  $0 < t \le \delta$  and  $\|\mathscr{T}(t_1) - \mathscr{T}(t_2)\| \le K |\log t_1 - \log t_2|$  for every  $0 < t_1, t_2 \le \delta$ .

Proof. It follows easily from  $(T_1)$ ,  $(T_2)$  (see the proof of Theorem 10.3.4 in [8]) that there is a subset  $D \subseteq E$  such that

$$(1) D is dense in E,$$

(2) the function  $\mathscr{T}(\cdot) x$  is infinitely differentiable for every  $x \in D$ .

For  $x \in E$  and  $q \in \{0, 1, ...\}$ , we shall write simply  $\mathscr{T}^{(q)}(t) x$  instead of  $(d^q/dt^q) \mathscr{T}(t) x$ , if this derivative exists. Instead of  $\mathscr{T}^{(1)}(t) x$  we write  $\mathscr{T}'(t) x$ .

(A)  $\Rightarrow$  (B): It follows from (A) and (2) that

(3) 
$$\|\mathscr{T}(t)x\| \leq Me^{\omega t} \|x\|$$
 for every  $x \in D$ ,

$$(4) \qquad \left\|\mathscr{T}(t_{1}) x - \mathscr{T}(t_{2}) x\right\| = \left\|\int_{t_{1}}^{t_{2}} \mathscr{T}'(\tau) x \, \mathrm{d}\tau\right\| \leq \int_{t_{1}}^{t_{2}} \left\|\mathscr{T}'(\tau) x\right\| \, \mathrm{d}\tau \leq \\ \leq \int_{t_{1}}^{t_{2}} M e^{\omega \tau} \frac{\varrho}{\tau} \, \mathrm{d}\tau \|x\| \leq M \varrho e^{\omega t_{2}} \int_{t_{1}}^{t_{2}} \frac{1}{\tau} \, \mathrm{d}\tau \|x\| \leq \\ \leq M \varrho e^{\omega t_{2}} (\log t_{1} - \log t_{2}) \|x\| \quad \text{for every} \quad x \in D \quad \text{and} \quad 0 < t_{1} \leq t_{2} \end{cases}$$

But (1), (3) and (4) imply immediately (B) if we take, for example,  $\delta = 1$ ,  $K = M + M \rho e^{\omega}$ .

(B)  $\Rightarrow$  (A): Let us first fix a  $\delta > 0$  and a  $K \ge 0$  so that (B) holds.

It is easy to see that for any  $t \in \mathbb{R}^+$ , there is a positive integer  $n_t$  such that

(5) 
$$(n_t - 1) \delta < t \leq n_t \delta$$
 for any  $t \in \mathbb{R}^+$ .

Since (5) yields  $0 < t - (n_t - 1)\delta \leq \delta$  and  $n_t < t/\delta + 1$ , we get from (T<sub>1</sub>) and (B)

(6) 
$$\|\mathscr{T}(t)\| = \|\mathscr{T}((n_t - 1)\delta + t - (n_t - 1)\delta)\| =$$
$$= \|\mathscr{T}(\delta)^{n_t - 1} \mathscr{T}(t - (n_t - 1)\delta)\| \le K^{n_t - 1}K =$$
$$= K^{n_t} \le K^{t/\delta + 1} = K(K^{1/\delta})^t = Ke^{(\log K^{1/\delta})t} \text{ for every } t \in R^+$$

Let us now take  $\varkappa = \log K^{1/\delta}$ . Then (6) implies

(7) 
$$\|\mathscr{T}(t)\| \leq Ke^{\kappa t}$$
 for every  $t \in R^+$ .

On the other hand, it is easy to see from  $(T_1)$  and (B) that

(8) 
$$\mathcal{T}'(s + t) x = \mathcal{T}(s) \mathcal{T}'(t) x = \mathcal{T}(t) \mathcal{T}'(s) x$$
  
for every  $t, s \in \mathbb{R}^+$  and  $x \in D$ .

Further, by (5) we can write

(9) 
$$n_t < t/\delta + 1$$
 and  $t/n_t \leq \delta$  for every  $t \in R^+$ .

It follows from (B) and (2) that

(10) 
$$\|\mathscr{T}'(t)x\| = \left\|\lim_{h\to 0+} \frac{1}{-h} \left(\mathscr{T}(t-h)x - \mathscr{T}(t)x\right\| \leq$$

$$\leq \lim_{h \to 0_+} \frac{1}{h} \left\| \mathscr{T}(t-h) x - \mathscr{T}(t) x \right\| \leq K \lim_{h \to 0_+} \frac{1}{-h} \left( \log \left( t-h \right) - \log t \right) \left\| x \right\| = \frac{K}{t} \left\| x \right\|$$

for every  $x \in D$  and  $0 < t \leq \delta$ .

Using  $(T_1)$ , (2) and (7)-(10) we get

(11) 
$$\|\mathscr{T}'(t)x\| = \left\|\mathscr{T}\left(t - \frac{t}{n_t}\right)\mathscr{T}'\left(\frac{t}{n_t}\right)x\right\| \leq Ke^{x(t-t/n_t)}K\frac{n_t}{t}\|x\| \leq K^2e^{xt}\frac{1}{t}\left(\frac{t}{\delta} + 1\right)\|x\| \leq K^2e^{xt}e^{t/\delta}\frac{1}{t}\|x\| = K^2e^{(x+\delta^{-1})t}\frac{1}{t}\|x\|$$
for every  $x \in D$  and  $t \in R^+$ .

It follows from (1) and (11) by virtue of the theorem on differentiation of limit that

(12) the function  $\mathscr{T}(\cdot) x$  is continuously differentiable for every  $x \in E$ ,

(13) 
$$\|\mathscr{T}'(t)x\| \leq K^2 e^{(x+\delta^{-1})t} \frac{1}{t} \|x\|$$
 for every  $x \in E$  and  $t \in R^+$ .

We see from  $(T_1)$  and (12) that

(14) 
$$\mathscr{T}'(t+s) x = \mathscr{T}(s) \mathscr{T}'(t) x = \mathscr{T}(t) \mathscr{T}'(s) x$$
  
for every  $x \in E$  and  $t, s \in R^+$ .

We shall now prove that

(15) the function 
$$\mathscr{T}(\cdot) x$$
 is *n*-times differentiable for every

$$x \in E$$
 and  $n \in \{1, 2, ...\}$ ,

(16) 
$$\mathscr{T}^{(n)}(t) x = (\mathscr{T}'(t/n))^n x$$
 for every  $x \in E$ ,  $t \in \mathbb{R}^+$  and  $n \in \{1, 2, ...\}$ .

We proceed by induction. The case n = 1 is true by (12). Let now (15) and (16) hold for a fixed  $n \in \{1, 2, ...\}$ . Then we see from (14) that, for every  $x \in E$ , the function  $\mathcal{T}(\cdot) x$  is (n + 1)-times differentiable and

$$\mathscr{T}^{(n+1)}(\tau + \sigma) x = \frac{\mathrm{d}^n}{\mathrm{d}\tau^n} \left[ \mathscr{T}(\tau) \, \mathscr{T}'(\sigma) \, x \right] = \mathscr{T}^{(n)}(\tau) \, \mathscr{T}'(\sigma) \, x \quad \text{for every} \quad \tau, \, \sigma \in \mathbb{R}^+ \; .$$

Hence (15) holds with n + 1 instead of *n* and moreover, taking  $\tau = tn/(n + 1)$  and  $\sigma = t/(n + 1)$  in the preceding formula, we get immediately (16) with n + 1 instead of *n* again.

Making use of (13) and (16) we get

$$(17) \quad \left\|\mathscr{T}^{(n)}(t) x\right\| = \left\| \left(\mathscr{T}'\left(\frac{t}{n}\right)\right)^n x\right\| \leq \left[ K^2 e^{(x+\delta^{-1})t/n} \frac{n}{t} \right]^n \|x\| = K^{2n} e^{(x+\delta^{-1})t} \frac{n^n}{t^n} \|x\| = K^{2n} e^{(x+\delta^{-1})t} \frac{1}{t^n} \frac{n^n}{n!} n! \|x\| \leq K^{2n} e^{(x+\delta^{-1})t} \frac{1}{t^n} e^n n! \|x\| = e^{(x+\delta^{-1})t} \frac{n! (K^2 e)^n}{t^n} \|x\|$$
for every  $x \in E$ ,  $t \in R^+$  and  $n \in \{1, 2, ...\}$ .

Let us now take  $M = \max(1, K)$ ,  $\omega = \varkappa + \delta^{-1}$  and  $\varrho = K^2 e$ . Then we get from (7) and (17) that

(18) 
$$\|\mathscr{F}^{(q)}(t) x\| \leq M e^{\omega t} \frac{q! \, q^q}{t^q} \|x\|$$
 for every  $x \in E$ ,  $t \in R^+$  and  $q \in \{0, 1, ...\}$ .

But (15) and (18) prove (A). The proof is complete. 11. Remark. The proof of the implication  $(B) \Rightarrow (A)$  in the preceding theorem uses an idea of Yosida [6] – see the beginning of the proof of his Theorem 1.

12. Proposition. Let  $\mathcal{F}$  be a semigroup of linear operators in E. Further, let  $M, \omega, \varrho$  be nonnegative constants. Then the following two statements (A) and (B) are equivalent:

(A) for every  $x \in E$  for which the function  $\mathcal{T}(\cdot) x$  is infinitely differentiable on  $R^+$  we have

$$\left\|\frac{\mathrm{d}^{q}}{\mathrm{d}t^{q}}\,\mathcal{T}(t)\,\mathbf{x}\right\| \leq Me^{\omega t}\,\frac{q\,!\,\varrho^{q}}{t^{q}}\,\|\mathbf{x}\|$$

for every  $t \in R^+$  and  $q \in \{0, 1, ...\}$ ,

(B) the function  $\mathcal{T}$  is infinitely differentiable on  $R^+$  (as a function on  $R^+$  into the Banach space of bounded linear operators on E into E) and

$$\left\|\mathscr{T}^{(q)}(t)\right\| = M e^{\omega t} \frac{q! \varrho^{q}}{t^{q}}$$

for every  $t \in R^+$  and  $q \in \{0, 1, ...\}$ .

**Proof.** The implication  $(B) \Rightarrow (A)$  is trivial. In the proof of  $(A) \Rightarrow (B)$  we use the properties (1), (2) from the proof of Theorem 9.

13. Remark. The semigroups satisfying the condition 12 (B) are usually called holomorphic or analytic. Mizohata [9] calls them parabolic which name seems to be the most specific and adequate as an abstract extension of the meaning of this term in the theory of partial differential equations. See also Remark 8.

The usual essentially equivalent approach to the parabolic semigroups consists in their characterization by means of an analytic continuation into a wedge-shaped domain around the positive half-axis (cf., e.g., [7] and [8]).

**14.** Theorem. Let A be a linear operator from E into E. Then the following two statements (A), (B) are equivalent:

(A) the operator A is the generator of a semigroup  $\mathcal{T}$  of linear operators in E (i.e., x belongs to the domain of A if and only if there is  $y \in E$  such that

$$\frac{1}{t}\left(\mathscr{F}(t) x - x\right) \xrightarrow[t \to 0_+]{} y;$$

in that case Ax = y, satisfying the condition (A) from Theorem 9,

(B) there exist nonnegative constants  $M, \omega, \varrho$  such that

(I)  $(\omega, \infty) \subseteq \varrho(A)$  (the resolvent set of A),

(II) 
$$\left\| \frac{\mathrm{d}^{p+q}}{\mathrm{d}^{p+q}} \left( \lambda^q (\lambda I - A)^{-1} \right) \right\| \leq \frac{Mp!}{(\lambda - \omega)^{p+1}} q! \varrho^q \text{ for every } \lambda > \omega \text{ and } p, q \in \{0, 1, \ldots\}.$$

Proof. A relatively easy consequence of the Hille-Yosida theorem, of the uniqueness theorem for the Laplace transform and of Theorem 5. A direct proof based only on Theorem 5 is also possible.

15. Remark. The preceding theorem gives a "real" characterization of the generators of parabolic semigroups. However, this problem is usually studied in a somewhat different setting which is based on "complex" characterizations of both the parabolic semigroups (cf. Remark 13) and their generators (see, e.g.,  $\lceil 7 \rceil$ ,  $\lceil 8 \rceil$ ).

16. Remark. In a subsequent paper (The Laplace transform of analytic vector-valued functions (complex conditions), to appear in this Journal), we shall study complex equivalents of both properties (A) and (B) from Theorem 5. These equivalents will be found in the form of analytic continuations of the functions F and f to certain angular domains symmetric around the real half-axis. These analytic continuations must satisfy certain growth conditions, but the relations between the growth characterizations of F and f are less simple and direct.

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