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# ON THE NORMALITY OF AN ALMOST CONTACT 3-STRUCTURE ON $Q R$-SUBMANIFOLDS 

S. Funabashi, Saitama, J. S. Pak, Taegu, and Y. J. Shin, Masan

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Abstract. We study $n$-dimensional $Q R$-submanifolds of $Q R$-dimension ( $p-1$ ) immersed in a quaternionic space form $Q P^{(n+p) / 4}(c), c \geqslant 0$, and, in particular, determine such submanifolds with the induced normal almost contact 3 -structure.

Keywords: quaternionic projective space, quaternionic number space, $Q R$-submanifold, normal almost contact 3 -structure

MSC 2000: 53C40

## 1. Introduction

Let $M$ be an $n$-dimensional $Q R$-submanifold of $Q R$-dimension ( $p-1$ ) isometrically immersed in a quaternionic Kähler manifold $\bar{M}^{(n+p) / 4}$. Denoting by $\{F, G, H\}$ the quaternionic Kähler structure of $\bar{M}^{(n+p) / 4}$, it follows by definition (cf. [9]) that there exists a $(p-1)$-dimensional subbundle $\nu$ of the normal bundle $T M^{\perp}$ such that

$$
\left\{\begin{array}{l}
F \nu_{x} \subset \nu_{x}, \quad G \nu_{x} \subset \nu_{x}, \quad H \nu_{x} \subset \nu_{x},  \tag{1.1}\\
F \nu_{x}^{\perp} \subset T_{x} M, \quad G \nu_{x}^{\perp} \subset T_{x} M, \quad H \nu_{x}^{\perp} \subset T_{x} M
\end{array}\right.
$$

for each $x \in M$, where $\nu^{\perp}$ denotes the complementary orthogonal subbundle to $\nu$ in $T M^{\perp}$. Thus there is a naturally distinguished unit normal vector field $\xi$ to $M$ such that $\nu_{x}^{\perp}=\operatorname{Span}\{\xi\}$ for each $x \in M$, and the vector fields $U, V, W$ defined by

$$
\begin{equation*}
U=-F \xi, \quad V=-G \xi, \quad W=-H \xi \tag{1.2}
\end{equation*}
$$

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are tangent to $M$. On the other hand, each tangent space $T_{x} M$ is decomposed as

$$
T_{x} M=D_{x} \oplus D_{x}^{\perp},
$$

where $D_{x}$ is the maximal quaternionic invariant subspace of $T_{x} M$ defined by

$$
D_{x}=T_{x} M \cap F T_{x} M \cap G T_{x} M \cap H T_{x} M
$$

and $D_{x}^{\perp}$ its orthogonal complement in $T_{x} M$. In our case, as already shown in [2], [9], $D_{x}^{\perp}=\operatorname{Span}\{U, V, W\}$ and so $D: x \mapsto D_{x}$ defines an $(n-3)$-dimensional distribution on $M$. But $D$ cannot be a quaternionic $C R$-distribution in the sense of [1]. Further it is clear that

$$
F T_{x} M, G T_{x} M, H T_{x} M \subset T_{x} M \oplus \operatorname{Span}\{\xi\}
$$

and, consequently, for any tangent vector $X$ to $M$, we have the following decomposition in tangential and normal components

$$
\left\{\begin{array}{l}
F X=\varphi X+u(X) \xi, \quad G X=\psi X+v(X) \xi  \tag{1.3}\\
H X=\theta X+w(X) \xi
\end{array}\right.
$$

By means of the hermitian property of $\{F, G, H\}$ it can be easily shown that $\varphi, \psi$ and $\theta$ are skew-symmetric endomorphisms acting on $T_{x} M$. Moreover it is known ([9], [10], [11]) that the aggregate $\{\varphi, \psi, \theta, u, v, w\}$ gives an almost contact 3 -structure on the $Q R$-submanifold $M$ of $Q R$-dimension $(p-1)$ in $\bar{M}^{(n+p) / 4}$ (see also Proposition 2.1).

On the other hand the normality of an almost contact 3 -structure was defined by one of the present authors ([13]) and by Yano, Ishihara and Konishi ([14]) in a different point of view. But, in this paper, it will be shown that the normalities of the induced almost contact 3 -structure in the sense of [13] and [14] are equivalent to each other, and the submanifold with the induced normal almost contact 3 -structure will be determined when the ambient manifold $\bar{M}$ is a quaternionic space form of constant $Q$-sectional curvature $c \geqslant 0$.

## 2. Fundamental formulas for $Q R$-submanifolds

Let $\bar{M}^{(n+p) / 4}$ be a real $(n+p)$-dimensional quaternionic Kähler manifold. Then, by definition, there is a 3 -dimensional vector bundle $V$ consisting of tensor fields of type $(1,1)$ over $\bar{M}$ satisfying the following conditions (a), (b) and (c):
(a) In any coordinate neighborhood $\overline{\mathscr{U}}$, there is a local basis $\{F, G, H\}$ of $V$ such that

$$
\left\{\begin{array}{l}
F^{2}=-I, \quad G^{2}=-I, \quad H^{2}=-I,  \tag{2.1}\\
F G=-G F=H, \quad G H=-H G=F, \quad H F=-F H=G .
\end{array}\right.
$$

(b) There is a Riemannian metric $g$ which is hermitian with respect to all of $F, G$ and $H$.
(c) For the Riemannian connection $\bar{\nabla}$ with respect to $g$

$$
\left(\begin{array}{c}
\bar{\nabla} F  \tag{2.2}\\
\bar{\nabla} G \\
\bar{\nabla} H
\end{array}\right)=\left(\begin{array}{rrr}
0 & r & -q \\
-r & 0 & p \\
q & -p & 0
\end{array}\right)\left(\begin{array}{l}
F \\
G \\
H
\end{array}\right)
$$

where $p, q$ and $r$ are local 1-forms defined in $\overline{\mathscr{U}}$. Such a local basis $\{F, G, H\}$ is called a canonical local basis of the bundle $V$ in $\overline{\mathscr{U}}$.
For canonical local bases $\{F, G, H\}$ and $\left\{{ }^{\prime} F,{ }^{\prime} G,{ }^{\prime} H\right\}$ of $V$ in coordinate neighborhoods $\overline{\mathscr{U}}$ and ${ }^{\prime} \overline{\mathscr{U}}$, it follows that in $\overline{\mathscr{U}} \cap{ }^{\prime} \overline{\mathscr{U}}$

$$
\left(\begin{array}{c}
{ }^{\prime} F  \tag{2.3}\\
\prime G \\
\prime \\
\prime
\end{array}\right)=\left(s_{x y}\right)\left(\begin{array}{l}
F \\
G \\
H
\end{array}\right) \quad(x, y=1,2,3)
$$

with differentiable functions $s_{x y}$, where the matrix $S=\left(s_{x y}\right)$ is contained in $S O(3)$ as a consequence of (2.1). As is well known [5], [6], every quaternionic Kähler manifold is orientable.

From now on we consider a real $n$-dimensional $Q R$-submanifold $M$ of $Q R$ dimension $(p-1)$ immersed in $\bar{M}^{(n+p) / 4}$ and use the same notations as in Section 1. We now take a local orthonormal basis $\left\{\xi_{\alpha} ; \alpha=1, \ldots, p\right\}\left(\xi_{1}=\xi\right)$ of normal vectors to $M$ and consider the following decompositions in tangential and normal components:

$$
\left\{\begin{array}{l}
F \xi_{\alpha}=-U_{\alpha}+P_{1} \xi_{\alpha}, \quad G \xi_{\alpha}=-V_{\alpha}+P_{2} \xi_{\alpha}  \tag{2.4}\\
H \xi_{\alpha}=-W_{\alpha}+P_{3} \xi_{\alpha}
\end{array}\right.
$$

$(\alpha=1, \ldots, p)$. Then $P_{1}, P_{2}$ and $P_{3}$ are skew-symmetric endomorphisms acting on $T_{x} M^{\perp}$. Moreover, by means of (1.3), the hermitian property of $\{F, G, H\}$ and (2.4) imply

$$
\begin{align*}
& \left\{\begin{array}{l}
g\left(X, \varphi U_{\alpha}\right)=-u(X) g\left(\xi_{1}, P_{1} \xi_{\alpha}\right), \\
g\left(X, \psi V_{\alpha}\right)=-v(X) g\left(\xi_{1}, P_{2} \xi_{\alpha}\right), \\
g\left(X, \theta W_{\alpha}\right)=-w(X) g\left(\xi_{1}, P_{3} \xi_{\alpha}\right), \quad \alpha=1, \ldots, p
\end{array}\right.  \tag{2.5}\\
& \left\{\begin{array}{l}
g\left(U_{\alpha}, U_{\beta}\right)=\delta_{\alpha \beta}-g\left(P_{1} \xi_{\alpha}, P_{1} \xi_{\beta}\right), \\
g\left(V_{\alpha}, V_{\beta}\right)=\delta_{\alpha \beta}-g\left(P_{2} \xi_{\alpha}, P_{2} \xi_{\beta}\right), \\
g\left(W_{\alpha}, W_{\beta}\right)=\delta_{\alpha \beta}-g\left(P_{3} \xi_{\alpha}, P_{3} \xi_{\beta}\right), \quad \alpha, \beta=1, \ldots, p
\end{array}\right. \tag{2.6}
\end{align*}
$$

Also, from $g\left(F X, \xi_{\alpha}\right)=-g\left(X, F \xi_{\alpha}\right), g\left(G X, \xi_{\alpha}\right)=-g\left(X, G \xi_{\alpha}\right)$ and $g\left(H X, \xi_{\alpha}\right)=$ $-g\left(X, H \xi_{\alpha}\right)$, it follows that

$$
g\left(X, U_{\alpha}\right)=u(X) \delta_{1 \alpha}, \quad g\left(X, V_{\alpha}\right)=v(X) \delta_{1 \alpha}, \quad g\left(X, W_{\alpha}\right)=w(X) \delta_{1 \alpha}
$$

and hence

$$
\begin{gather*}
g\left(U_{1}, X\right)=u(X), \quad g\left(V_{1}, X\right)=v(X), \quad g\left(W_{1}, X\right)=w(X),  \tag{2.7}\\
U_{\alpha}=0, \quad V_{\alpha}=0, \quad W_{\alpha}=0, \quad \alpha=2, \ldots, p
\end{gather*}
$$

On the other hand, comparing (1.2) and (2.4) with $\alpha=1$, we have $U_{1}=U, V_{1}=V$, $W_{1}=W$, which together with (2.7) imply

$$
\begin{equation*}
g(U, X)=u(X), \quad g(V, X)=v(X), \quad g(W, X)=w(X) \tag{2.8}
\end{equation*}
$$

In the sequel we shall use the notations $U, V, W$ instead of $U_{1}, V_{1}, W_{1}$.
Next, applying $F$ to the first equation of (1.3) and using (2.4), (2.7) and (2.8), we have

$$
\varphi^{2} X=-X+u(X) U, \quad u(X) P_{1} \xi=-u(\varphi X) \xi
$$

Similarly we have

$$
\begin{align*}
& \left\{\begin{array}{l}
\varphi^{2} X=-X+u(X) U, \quad \psi^{2} X=-X+v(X) V \\
\theta^{2} X=-X+w(X) W,
\end{array}\right.  \tag{2.9}\\
& \left\{\begin{array}{l}
u(X) P_{1} \xi=-u(\varphi X) \xi, \quad v(X) P_{2} \xi=-v(\psi X) \xi \\
w(X) P_{3} \xi=-w(\theta X) \xi,
\end{array}\right. \tag{2.10}
\end{align*}
$$

from which, taking account of the skew-symmetry of $P_{1}, P_{2}$ and $P_{3}$ and using (2.5), we also have

$$
\left\{\begin{array}{l}
u(\varphi X)=0, \quad v(\psi X)=0, \quad w(\theta X)=0  \tag{2.11}\\
\varphi U=0, \quad \psi V=0, \quad \theta W=0 \\
P_{1} \xi=0, \quad P_{2} \xi=0, \quad P_{3} \xi=0
\end{array}\right.
$$

So (2.4) can be rewritten in the form

$$
\left\{\begin{array}{l}
F \xi=-U, \quad G \xi=-V, \quad H \xi=-W  \tag{2.12}\\
F \xi_{\alpha}=P_{1} \xi_{\alpha}, \quad G \xi_{\alpha}=P_{2} \xi_{\alpha}, \quad H \xi_{\alpha}=P_{3} \xi_{\alpha}
\end{array}\right.
$$

where $\alpha=2, \ldots, p$.

Applying $G$ and $H$ to the first equation of (1.3) and using (1.3), (2.1) and (2.12), we have

$$
\begin{aligned}
& \theta X+w(X) \xi=-\psi(\varphi X)-v(\varphi X) \xi+u(X) V \\
& \psi X+v(X) \xi=\theta(\varphi X)+w(\varphi X) \xi-u(X) W
\end{aligned}
$$

and consequently

$$
\left\{\begin{array}{l}
\psi(\varphi X)=-\theta X+u(X) V, \quad v(\varphi X)=-w(X)  \tag{2.13}\\
\theta(\varphi X)=\psi X+u(X) W, \quad w(\varphi X)=v(X)
\end{array}\right.
$$

From the other equations of (1.3) we have by a quite similar method

$$
\begin{align*}
& \left\{\begin{array}{l}
\varphi(\psi X)=\theta X+v(X) U, \quad u(\psi X)=w(X) \\
\theta(\psi X)=-\varphi X+v(X) W, \quad w(\psi X)=-u(X)
\end{array}\right.  \tag{2.14}\\
& \left\{\begin{array}{l}
\varphi(\theta X)=-\psi X+w(X) U, \quad u(\theta X)=-v(X) \\
\psi(\theta X)=\varphi X+w(X) V, \quad v(\theta X)=u(X)
\end{array}\right. \tag{2.15}
\end{align*}
$$

From the first three equations of (2.12), we also have

$$
\left\{\begin{array}{l}
\psi U=-W, \quad v(U)=0, \quad \theta U=V, \quad w(U)=0  \tag{2.16}\\
\varphi V=W, \quad u(V)=0, \quad \theta V=-U, \quad w(V)=0 \\
\varphi W=-V, \quad u(W)=0, \quad \psi W=U, \quad v(W)=0
\end{array}\right.
$$

On the other hand, we may put

$$
\begin{cases}P_{1} \xi_{\alpha}=\sum_{\beta=2}^{p} P_{1 \alpha \beta} \xi_{\beta}, & P_{2} \xi_{\alpha}=\sum_{\beta=2}^{p} P_{2 \alpha \beta} \xi_{\beta}  \tag{2.17}\\ P_{3} \xi_{\alpha}=\sum_{\beta=2}^{p} P_{3 \alpha \beta} \xi_{\beta}, & \alpha=2, \ldots, p\end{cases}
$$

from which, substituting into the last three equations of (2.12) and using the hermitian property of $\{F, G, H\}$, we have

$$
\left\{\begin{array}{l}
\sum_{\gamma} P_{1 \alpha \gamma} P_{1 \gamma \beta}=-\delta_{\alpha \beta}, \quad \sum_{\gamma} P_{2 \alpha \gamma} P_{2 \gamma \beta}=-\delta_{\alpha \beta},  \tag{2.18}\\
\sum_{\gamma} P_{3 \alpha \gamma} P_{3 \gamma \beta}=-\delta_{\alpha \beta} .
\end{array}\right.
$$

Also, from (2.1), (2.12) and (2.17), we have

$$
\begin{cases}\sum_{\beta} P_{1 \alpha \beta} P_{2 \beta \gamma}=-P_{3 \alpha \gamma}, & \sum_{\beta} P_{1 \alpha \beta} P_{3 \beta \gamma}=P_{2 \alpha \gamma}  \tag{2.19}\\ \sum_{\beta} P_{2 \alpha \beta} P_{3 \beta \gamma}=-P_{1 \alpha \gamma}, & \sum_{\beta} P_{2 \alpha \beta} P_{1 \beta \gamma}=P_{3 \alpha \gamma} \\ \sum_{\beta} P_{3 \alpha \beta} P_{1 \beta \gamma}=-P_{2 \alpha \gamma}, & \sum_{\beta} P_{3 \alpha \beta} P_{2 \beta \gamma}=P_{1 \alpha \gamma}\end{cases}
$$

The equations (2.6)-(2.11) and (2.13)-(2.16) tell us

Proposition 2.1 ([9], [10], [11]). An n-dimensional $Q R$-submanifold of $Q R$ dimension $(p-1)$ in a quaternionic Kähler manifold $\bar{M}^{(n+p) / 4}$ admits an almost contact 3 -structure.

In general if the condition

$$
\left[\varphi_{i}, \varphi_{i}\right]+d u_{i} \otimes U_{i}=0
$$

is satisfied for some $1 \leqslant i \leqslant 3$, then the almost contact structure $\left(\varphi_{i}, U_{i}, u_{i}\right)$ is said to be normal, where we put

$$
\begin{gathered}
\varphi_{1}=\varphi, \quad \varphi_{2}=\psi, \quad \varphi_{3}=\theta \\
U_{1}=U, \quad U_{2}=V, \quad U_{3}=W ; \quad u_{1}=u, \quad u_{2}=v, \quad u_{3}=w
\end{gathered}
$$

and $\left[\varphi_{i}, \varphi_{i}\right]$ denotes the Nijenhuis tensor of $\varphi_{i}$. In their papers [8] and [14], Ishihara, Konishi, Kuo and Yano have proved

Lemma 2.2. If, for an almost contact 3-structure $\left\{\left(\varphi_{i}, U_{i}, u_{i}\right) ; i=1,2,3\right\}$, any two of the almost contact structures $\left(\varphi_{i}, U_{i}, u_{i}\right)$ are normal, then so is the third.

Moreoreover, in [14] the following lemma was proved.

Lemma 2.3. For an almost contact 3 -structure $\left\{\left(\varphi_{i}, U_{i}, u_{i}\right) ; i=1,2,3\right\}$, a necessary and sufficient condition in order that the almost contact structures $\left(\varphi_{i}, U_{i}, u_{i}\right)$ are all normal is that the condition

$$
\left\{\begin{array}{l}
2\left[\varphi_{1}, \varphi_{2}\right]+d u_{1} \otimes U_{2}+d u_{2} \otimes U_{1}=0  \tag{2.20}\\
\mathscr{L}_{U_{1}} \varphi_{2}+\mathscr{L}_{U_{2}} \varphi_{1}=0, \quad d u_{1} \bar{\wedge} \varphi_{2}+d u_{2} \bar{\wedge} \varphi_{1}=0
\end{array}\right.
$$

be valid, where $\left[\varphi_{1}, \varphi_{2}\right]$ denotes the Nijenhuis tensor of $\varphi_{1}$ and $\varphi_{2}, d u_{i} \bar{\wedge} \varphi_{j}$ the 2 -form defined by

$$
\left(d u_{i} \bar{\wedge} \varphi_{j}\right)(X, Y)=d u_{i}\left(\varphi_{j} X, Y\right)+d u_{i}\left(X, \varphi_{j} Y\right)
$$

and $\mathscr{L}_{U_{i}}$ the Lie derivative with respect to $U_{i}$.

## 3. Further properties of the induced almost contact 3-structure

In this section we shall use the same notations and terminology as in the previous section.

Now let $\nabla$ be the Levi-Cività connection on $M$ and $\nabla^{\perp}$ the normal connection induced from $\bar{\nabla}$ in the normal bundle $T M^{\perp}$ of $M$. Then Gauss and Weingarten formulae are given by

$$
\begin{align*}
\bar{\nabla}_{X} Y & =\nabla_{X} Y+h(X, Y)  \tag{3.1}\\
\bar{\nabla}_{X} \xi_{\alpha} & =-A_{\alpha} X+\nabla_{X}^{\perp} \xi_{\alpha}, \quad \alpha=1, \ldots, p \tag{3.2}
\end{align*}
$$

for $X, Y$ tangent to $M$. Here $h$ denotes the second fundamental form and $A_{\alpha}$ the shape operator corresponding to $\xi_{\alpha}$. They are related by $h(X, Y)=$ $\sum_{\alpha=1}^{p} g\left(A_{\alpha} X, Y\right) \xi_{\alpha}$. Furthermore, put

$$
\begin{equation*}
\nabla \stackrel{\perp}{X} \xi_{\alpha}=\sum_{\beta=1}^{p} s_{\alpha \beta}(X) \xi_{\beta} \tag{3.3}
\end{equation*}
$$

where $\left(s_{\alpha \beta}\right)$ is the skew-symmetric matrix of connection forms of $\nabla^{\perp}$.
Differentiating the first equation of (1.3) covariantly and using (1.3), (2.2), (2.4), (2.7), (3.1) and (3.2), we have

$$
\begin{align*}
& \left(\nabla_{Y} \varphi\right) X=r(Y) \psi X-q(Y) \theta X+u(X) A_{1} Y-g\left(A_{1} Y, X\right) U  \tag{3.4}\\
& \left(\nabla_{Y} u\right) X=r(Y) v(X)-q(Y) w(X)+g\left(\varphi A_{1} Y, X\right)
\end{align*}
$$

From the other equations of (1.3) we also have

$$
\begin{align*}
\left(\nabla_{Y} \psi\right) X & =-r(Y) \varphi X+p(Y) \theta X+v(X) A_{1} Y-g\left(A_{1} Y, X\right) V  \tag{3.5}\\
\left(\nabla_{Y} v\right) X & =-r(Y) u(X)+p(Y) w(X)+g\left(\psi A_{1} Y, X\right) \\
\left(\nabla_{Y} \theta\right) X & =q(Y) \varphi X-p(Y) \psi X+w(X) A_{1} Y-g\left(A_{1} Y, X\right) W  \tag{3.6}\\
\left(\nabla_{Y} w\right) X & =q(Y) u(X)-p(Y) v(X)+g\left(\theta A_{1} Y, X\right)
\end{align*}
$$

Next, differentiating the first equation of (2.12) covariantly and comparing the tangential and normal parts, we have

$$
\left\{\begin{array}{l}
\nabla_{Y} U=r(Y) V-q(Y) W+\varphi A_{1} Y,  \tag{3.7}\\
g\left(A_{\alpha} U, Y\right)=-\sum_{\beta=2}^{p} s_{1 \beta}(Y) P_{1 \beta \alpha}, \quad \alpha=2, \ldots, p
\end{array}\right.
$$

From the other equations of (2.12), we have similarly

$$
\begin{align*}
& \left\{\begin{array}{l}
\nabla_{Y} V=-r(Y) U+p(Y) W+\psi A_{1} Y, \\
g\left(A_{\alpha} V, Y\right)=-\sum_{\beta=2}^{p} s_{1 \beta}(Y) P_{2 \beta \alpha}, \quad \alpha=2, \ldots, p
\end{array}\right.  \tag{3.8}\\
& \left\{\begin{array}{l}
\nabla_{Y} W=q(Y) U-p(Y) V+\theta A_{1} Y, \\
g\left(A_{\alpha} W, Y\right)=-\sum_{\beta=2}^{p} s_{1 \beta}(Y) P_{3 \beta \alpha}, \quad \alpha=2, \ldots, p
\end{array}\right. \tag{3.9}
\end{align*}
$$

In what follows we assume that the distinguished normal vector field $\xi$ is parallel with respect to the normal connection, that is, $\nabla \frac{\perp}{X} \xi=0$. Hence it follows from (3.3) that $s_{\beta 1}=0, \beta=2, \ldots, p$, and, consequently,

$$
A_{\alpha} U=0, \quad A_{\alpha} V=0, \quad A_{\alpha} W=0, \quad \alpha=2, \ldots, p
$$

because of (3.7)-(3.9).
In particular when the ambient manifold is a quaternionic space form $\bar{M}^{(n+p) / 4}(c)$, that is, a quaternionic Kähler manifold of constant $Q$-sectional curvature $c$, the curvature tensor $\bar{R}$ of $\bar{M}^{(n+p) / 4}(c)$ has the form

$$
\begin{aligned}
\bar{R}_{\overline{X Y}} \bar{Z}= & \frac{c}{4}\{g(\bar{Y}, \bar{Z}) \bar{X}-g(\bar{X}, \bar{Z}) \bar{Y} \\
& +g(F \bar{Y}, \bar{Z}) F \bar{X}-g(F \bar{X}, \bar{Z}) F \bar{Y}-2 g(F \bar{X}, \bar{Y}) F \bar{Z} \\
& +g(G \bar{Y}, \bar{Z}) G \bar{X}-g(G \bar{X}, \bar{Z}) G \bar{Y}-2 g(G \bar{X}, \bar{Y}) G \bar{Z} \\
& +g(H \bar{Y}, \bar{Z}) H \bar{X}-g(H \bar{X}, \bar{Z}) H \bar{Y}-2 g(H \bar{X}, \bar{Y}) H \bar{Z}\}
\end{aligned}
$$

for $\bar{X}, \bar{Y}, \bar{Z}$ tangent to $\bar{M}^{(n+p) / 4}(c)$ (cf. [5], [6]). So the above assumption implies that the equation of Codazzi and Ricci is of the form

$$
\begin{align*}
g\left(\left(\nabla_{X} A_{1}\right) Y\right. & \left.-\left(\nabla_{Y} A_{1}\right) X, Z\right)  \tag{3.10}\\
= & \frac{c}{4}\{g(\varphi Y, Z) u(X)-g(\varphi X, Z) u(Y)-2 g(\varphi X, Y) u(Z) \\
& +g(\psi Y, Z) v(X)-g(\psi X, Z) v(Y)-2 g(\psi X, Y) v(Z) \\
& +g(\theta Y, Z) w(X)-g(\theta X, Z) w(Y)-2 g(\theta X, Y) w(Z)\},
\end{align*}
$$

$$
\begin{equation*}
g\left(\bar{R}(X, Y) \xi_{\alpha}, \xi_{\beta}\right)=g\left(R^{\perp}(X, Y) \xi_{\alpha}, \xi_{\beta}\right)+g\left(\left[A_{\beta}, A_{\alpha}\right] X, Y\right) \tag{3.11}
\end{equation*}
$$

for any $X, Y, Z$ tangent to $M$, where $R$ and $R^{\perp}$ denote the curvature tensor of $\nabla$ and $\nabla^{\perp}$, respectively (cf. [3], [9], [10], [11]).

Finally we introduce a theorem due to Kwon and one of the present authors ([9]) for later use.

Theorem K-P. Let $M$ be an $n$-dimensional $Q R$-submanifold of $Q R$-dimension $(p-1)$ in a quaternionic projective space $Q P^{(n+p) / 4}(4)$ and let the normal vector field $\xi$ be parallel with respect to the normal connection. If

$$
A_{1} \varphi=\varphi A_{1}, \quad A_{1} \psi=\psi A_{1}, \quad A_{1} \theta=\theta A_{1}
$$

on $M$, then $\pi^{-1}(M)$ is locally a product of $M_{1} \times M_{2}$ where $M_{1}$ and $M_{2}$ lie on some $\left(4 n_{1}+3\right)$ - and $\left(4 n_{2}+3\right)$-spheres, respectively, and $A_{1}$ denotes the shape operator corresponding to $\xi$ ( $\pi$ is the Hopf fibration $S^{n+p+3}(1) \rightarrow Q P^{(n+p) / 4}(4)$ ).

## 4. The submanifolds with the induced normal ALMOST CONTACT 3-STRUCTURE

In this section we introduce the notion of the normality of almost contact 3-structure in the sense of [13].

From now on we put in each coordinate neighborhood $\mathscr{U}$ of $M$

$$
\begin{align*}
& \left(\begin{array}{c}
\stackrel{\circ}{\nabla} \varphi \\
\stackrel{\circ}{\nabla} \psi \\
\stackrel{\circ}{\nabla} \theta
\end{array}\right)=\left(\begin{array}{c}
\nabla \varphi \\
\nabla \psi \\
\nabla \theta
\end{array}\right)+\left(\begin{array}{rrr}
0 & r & -q \\
-r & 0 & p \\
q & -p & 0
\end{array}\right)\left(\begin{array}{l}
\varphi \\
\psi \\
\theta
\end{array}\right),  \tag{4.1}\\
& \left(\begin{array}{c}
\stackrel{\circ}{\nabla} U \\
\stackrel{\circ}{\nabla} V \\
\stackrel{\circ}{\nabla} W
\end{array}\right)=\left(\begin{array}{c}
\nabla U \\
\nabla V \\
\nabla W
\end{array}\right)+\left(\begin{array}{rrr}
0 & r & -q \\
-r & 0 & p \\
q & -p & 0
\end{array}\right)\left(\begin{array}{c}
U \\
V \\
W
\end{array}\right) . \tag{4.2}
\end{align*}
$$

Then it follows from (2.3) that

$$
\left(\begin{array}{c}
\stackrel{\circ}{\nabla^{\prime}} \varphi  \tag{4.3}\\
\stackrel{\circ}{\nabla^{\prime}} \psi \\
\stackrel{\circ}{\nabla^{\prime}} \theta
\end{array}\right)=\left(s_{x y}\right)\left(\begin{array}{c}
\stackrel{\circ}{\nabla} \varphi \\
\stackrel{\circ}{\nabla} \psi \\
\stackrel{\circ}{\nabla} \theta
\end{array}\right), \quad\left(\begin{array}{c}
\stackrel{\circ}{\nabla^{\prime}} U \\
\stackrel{\circ}{\nabla^{\prime}} V \\
\stackrel{\circ}{\nabla}^{\prime} W
\end{array}\right)=\left(s_{x y}\right)\left(\begin{array}{c}
\stackrel{\circ}{\nabla} U \\
\stackrel{\circ}{\nabla} V \\
\stackrel{\circ}{\nabla} W
\end{array}\right)
$$

in $\mathscr{U} \cap^{\prime} \mathscr{U}$. Now, in each coordinate neighborhood $\mathscr{U}$, we consider local tensor fields $S\left(\varphi_{i}, \varphi_{j}\right)(i, j=1,2,3)$ of type $(1,2)$ such that

$$
\begin{align*}
S\left(\varphi_{i}, \varphi_{j}\right) & (X, Y)  \tag{4.4}\\
= & \left(\stackrel{\circ}{\nabla}_{\varphi_{i} X} \varphi_{j}\right) Y-\left(\stackrel{\circ}{\nabla}_{\varphi_{i} Y} \varphi_{j}\right) X+\left(\stackrel{\circ}{\nabla}_{\varphi_{j} X} \varphi_{i}\right) Y-\left(\stackrel{\circ}{\nabla}_{\varphi_{j} Y} \varphi_{i}\right) X \\
& +\varphi_{i}\left\{\left(\stackrel{\circ}{\nabla}_{Y} \varphi_{j}\right) X-\left(\stackrel{\circ}{\nabla}_{X} \varphi_{j}\right) Y\right\}+\varphi_{j}\left\{\left(\stackrel{\circ}{\nabla}_{Y}\right) X-\left(\stackrel{\circ}{\nabla}_{X} \varphi_{i}\right) Y\right\} \\
& +\left\{\left(\stackrel{\circ}{\nabla}_{X} u_{i}\right) Y-\left(\stackrel{\circ}{\nabla}_{Y} u_{i}\right) X\right\} U_{j}+\left\{\left(\stackrel{\circ}{\nabla}_{X} u_{j}\right) Y-\left(\stackrel{\circ}{\nabla}_{Y} u_{j}\right) X\right\} U_{i}
\end{align*}
$$

where we again put

$$
\varphi_{1}=\varphi, \quad \varphi_{2}=\psi, \quad \varphi_{3}=\theta, \quad U_{1}=U, \quad U_{2}=V, \quad U_{3}=W
$$

and

$$
\begin{equation*}
\left(\stackrel{\circ}{\nabla}_{X} u_{i}\right) Y=g\left(\stackrel{\circ}{\nabla}_{X} U_{i}, Y\right), \quad i=1,2,3 . \tag{4.5}
\end{equation*}
$$

Then a simple computation using (4.3) implies that

$$
S\left({ }^{\prime} \varphi_{i},{ }^{\prime} \varphi_{j}\right)=\left(s_{x y}\right)\left(S\left(\varphi_{i}, \varphi_{j}\right)\right)\left(s_{x y}\right)^{-1}
$$

in $\mathscr{U} \cap^{\prime} \mathscr{U}$. Hence we have the global tensor fields $\Sigma_{1}$ and $\Sigma_{2}$ on $M$ defined by

$$
\begin{align*}
\Sigma_{1}= & S\left(\varphi_{1}, \varphi_{1}\right)+S\left(\varphi_{2}, \varphi_{2}\right)+S\left(\varphi_{3}, \varphi_{3}\right)  \tag{4.6}\\
\Sigma_{2}= & S\left(\varphi_{1}, \varphi_{1}\right) \otimes S\left(\varphi_{2}, \varphi_{2}\right)+S\left(\varphi_{2}, \varphi_{2}\right) \otimes S\left(\varphi_{3}, \varphi_{3}\right)  \tag{4.7}\\
& +S\left(\varphi_{3}, \varphi_{3}\right) \otimes S\left(\varphi_{1}, \varphi_{1}\right)-S\left(\varphi_{1}, \varphi_{2}\right) \otimes S\left(\varphi_{2}, \varphi_{1}\right) \\
& -S\left(\varphi_{2}, \varphi_{3}\right) \otimes S\left(\varphi_{3}, \varphi_{2}\right)-S\left(\varphi_{3}, \varphi_{1}\right) \otimes S\left(\varphi_{1}, \varphi_{3}\right)
\end{align*}
$$

up to a sign. It is said that the induced almost contact 3-structure is normal if $\Sigma_{1}=0$ and $\Sigma_{2}=0$ (for details see [13]).

Remark 4.1 ([13]). A necessary and sufficient condition in order for the almost contact 3 -structure to be normal is

$$
S\left(\varphi_{i}, \varphi_{j}\right)=0, \quad i, j=1,2,3
$$

We next consider the traceless part of $\delta$-decomposition of the global tensor field $\Sigma_{1}$ in the sense of Krupka ([7]). Since $\Sigma_{1}$ is of type (1,2) and $n \geqslant 2$, using (3.4)-(3.6) and (4.4)-(4.6) we can easily verify that the traceless part $\Sigma_{1}^{\circ}$ of $\Sigma_{1}$ is given by

$$
\begin{align*}
\stackrel{\circ}{\Sigma}_{1}(X, Y)= & \Sigma_{1}(X, Y)-\frac{1}{2(n-1)}\left\{u\left(A_{1} \varphi Y\right) X-u\left(A_{1} \varphi X\right) Y\right.  \tag{4.8}\\
& \left.+v\left(A_{1} \psi Y\right) X-v\left(A_{1} \psi X\right) Y+w\left(A_{1} \theta Y\right) X-w\left(A_{1} \theta X\right) Y\right\}
\end{align*}
$$

or equivalently

$$
\begin{align*}
2 \stackrel{\circ}{\Sigma}_{1}(X, Y)= & u(Y)\left(A_{1} \varphi-\varphi A_{1}\right) X-u(X)\left(A_{1} \varphi-\varphi A_{1}\right) Y  \tag{4.8}\\
& +v(Y)\left(A_{1} \psi-\psi A_{1}\right) X-v(X)\left(A_{1} \psi-\psi A_{1}\right) Y \\
& +w(Y)\left(A_{1} \theta-\theta A_{1}\right) X-w(X)\left(A_{1} \theta-\theta A_{1}\right) Y \\
& -\frac{1}{n-1}\left\{u\left(A_{1} \varphi Y\right) X-u\left(A_{1} \varphi X\right) Y+v\left(A_{1} \psi Y\right) X\right. \\
& \left.-v\left(A_{1} \psi X\right) Y+w\left(A_{1} \theta Y\right) X-w\left(A_{1} \theta X\right) Y\right\} .
\end{align*}
$$

From now on we assume that $\Sigma_{1}=0$ identically on $M$. Putting $Y=U$ in (4.8) ${ }^{\prime}$ with $\Sigma_{1}^{\circ}=0$ and using (2.13)-(2.16), we obtain

$$
\begin{align*}
0= & \left(A_{1} \varphi-\varphi A_{1}\right) X+u(X) \varphi A_{1} U+v(X)\left\{A_{1} W+\psi A_{1} U\right\}  \tag{4.9}\\
& -w(X)\left\{A_{1} V-\theta A_{1} U\right\} \\
& +\frac{1}{n-1}\left\{u\left(A_{1} \varphi X\right)+v\left(A_{1} \psi X\right)+w\left(A_{1} \theta X\right)\right\} U,
\end{align*}
$$

from which, taking the inner product with $U$, it follows that

$$
\begin{equation*}
\frac{1}{n-1}\left(n \varphi A_{1} U+\psi A_{1} V+\theta A_{1} W\right)=2\left\{u\left(A_{1} W\right) V-u\left(A_{1} V\right) W\right\} \tag{4.10}
\end{equation*}
$$

Taking the inner product of (5.3) with $V$ and $W$, respectively, and using (2.13)(2.16), we have

$$
u\left(A_{1} W\right)=u\left(A_{1} V\right)=0
$$

which together with (4.10) yields

$$
n \varphi A_{1} U+\psi A_{1} V+\theta A_{1} W=0
$$

Similarly we have

$$
\begin{aligned}
& n \varphi A_{1} U+\psi A_{1} V+\theta A_{1} W=0 \\
& \varphi A_{1} U+n \psi A_{1} V+\theta A_{1} W=0 \\
& \varphi A_{1} U+\psi A_{1} V+n \theta A_{1} W=0
\end{aligned}
$$

and, consequently,

$$
\varphi A_{1} U=\psi A_{1} V=\theta A_{1} W=0
$$

Moreover, the last equations imply

$$
A_{1} U=u\left(A_{1} U\right) U, \quad A_{1} V=v\left(A_{1} V\right) V, \quad A_{1} W=w\left(A_{1} W\right) W
$$

which together with (4.8) gives the following implication:

$$
\stackrel{\circ}{\Sigma}_{1}=0 \Longrightarrow \Sigma_{1}=0
$$

Since the converse is trivial, we have

Lemma 4.1. Let $M$ be an $n$-dimensional $Q R$-submanifold of $Q R$-dimension $(p-1)$ in a quaternionic Kähler manifold $\bar{M}^{(n+p) / 4}$ and let the normal vector field $\xi$ be parallel with respect to the normal connection. Then we have

$$
\stackrel{\circ}{\Sigma}_{1}=0 \Longleftrightarrow \Sigma_{1}=0
$$

By means of Lemma 4.1 we have
Theorem 1. Let $M$ be as in Lemma 4.1. Then the following are equivalent to each other:
(a) The almost contact 3-structure is normal.
(b) The global tensor field $\Sigma_{1}$ defined by (4.6) vanishes.
(c) The traceless part $\Sigma_{1}$ of $\Sigma_{1}$ vanishes.
(d) The relation given by $(2.20)$ is valid.
(e) $A_{1} \varphi=\varphi A_{1}, \quad A_{1} \psi=\psi A_{1}, \quad A_{1} \theta=\theta A_{1}$.

Proof. Substituting (3.4)-(3.9) into (4.4), we can easily obtain that

$$
\begin{align*}
S(\varphi, \varphi)(X, Y) & =2\left\{u(Y)\left(A_{1} \varphi-\varphi A_{1}\right) X-u(X)\left(A_{1} \varphi-\varphi A_{1}\right) Y\right\}  \tag{4.11}\\
S(\psi, \psi)(X, Y) & =2\left\{v(Y)\left(A_{1} \psi-\psi A_{1}\right) X-v(X)\left(A_{1} \psi-\psi A_{1}\right) Y\right\} \\
S(\theta, \theta)(X, Y) & =2\left\{w(Y)\left(A_{1} \theta-\theta A_{1}\right) X-w(X)\left(A_{1} \theta-\theta A_{1}\right) Y\right\}
\end{align*}
$$

and

$$
\begin{align*}
S(\varphi, \psi)(X, Y)= & v(Y)\left(A_{1} \varphi-\varphi A_{1}\right) X-v(X)\left(A_{1} \varphi-\varphi A_{1}\right) Y  \tag{4.12}\\
& +u(Y)\left(A_{1} \psi-\psi A_{1}\right) X-u(X)\left(A_{1} \psi-\psi A_{1}\right) Y \\
S(\psi, \theta)(X, Y)= & w(Y)\left(A_{1} \psi-\psi A_{1}\right) X-w(X)\left(A_{1} \psi-\psi A_{1}\right) Y \\
& +v(Y)\left(A_{1} \theta-\theta A_{1}\right) X-v(X)\left(A_{1} \theta-\theta A_{1}\right) Y \\
S(\theta, \varphi)(X, Y)= & u(Y)\left(A_{1} \theta-\theta A_{1}\right) X-u(X)\left(A_{1} \theta-\theta A_{1}\right) Y \\
& +w(Y)\left(A_{1} \varphi-\varphi A_{1}\right) X-w(X)\left(A_{1} \varphi-\varphi A_{1}\right) Y
\end{align*}
$$

which together with Lemmas 2.2, 2.3 and Remark 4.1 yields the implications

$$
(\mathrm{e}) \Longrightarrow(\mathrm{a}), \quad(\mathrm{e}) \Longrightarrow(\mathrm{b}), \quad(\mathrm{e}) \Longrightarrow(\mathrm{d})
$$

In order to prove that the other implications are valid, it suffices to show the implication $(\mathrm{b}) \Longrightarrow(\mathrm{e})$. Now we assume that $(\mathrm{b})$ is valid. Then (4.11) implies

$$
\begin{align*}
u(Y)\left(A_{1} \varphi\right. & \left.-\varphi A_{1}\right) X-u(X)\left(A_{1} \varphi-\varphi A_{1}\right) Y  \tag{4.13}\\
& +v(Y)\left(A_{1} \psi-\psi A_{1}\right) X-v(X)\left(A_{1} \psi-\psi A_{1}\right) Y \\
& +w(Y)\left(A_{1} \theta-\theta A_{1}\right) X-w(X)\left(A_{1} \theta-\theta A_{1}\right) Y=0
\end{align*}
$$

Putting $Y=U$ in (4.13) and using (2.11) and (2.16), we have

$$
\begin{align*}
\left(A_{1} \varphi-\varphi A_{1}\right) X & -u(X) \varphi A_{1} U+v(X)\left(A_{1} W+\psi A_{1} U\right)  \tag{4.14}\\
& -w(X)\left(A_{1} V-\theta A_{1} U\right)=0
\end{align*}
$$

from which, taking the inner product with $U$, it follows that

$$
g\left(\varphi A_{1} U, X\right)=2 u\left(A_{1} W\right) v(X)-2 u\left(A_{1} V\right) w(X)
$$

and, consequently,

$$
\varphi A_{1} U=0, \quad u\left(A_{1} W\right)=0, \quad u\left(A_{1} V\right)=0
$$

Similarly we have

$$
\begin{gather*}
A_{1} U=u\left(A_{1} U\right) U, \quad A_{1} V=v\left(A_{1} V\right) V, \quad A_{1} W=w\left(A_{1} W\right) W  \tag{4.15}\\
u\left(A_{1} V\right)=v\left(A_{1} U\right)=u\left(A_{1} W\right)=w\left(A_{1} U\right)  \tag{4.16}\\
=v\left(A_{1} W\right)=w\left(A_{1} V\right)=0
\end{gather*}
$$

Substituting (4.15) into (4.14) and using (2.16), we have

$$
\begin{align*}
\left(A_{1} \varphi-\varphi A_{1}\right) X & +v(X)\left\{w\left(A_{1} W\right)-u\left(A_{1} U\right)\right\} W  \tag{4.17}\\
& -w(X)\left\{v\left(A_{1} V\right)-u\left(A_{1} U\right)\right\} V=0
\end{align*}
$$

from which, taking the symmetric part,

$$
\begin{aligned}
2 g\left(\left(A_{1} \varphi-\varphi A_{1}\right) X, Y\right) & +\left\{w\left(A_{1} W\right)-v\left(A_{1} V\right)\right\} \\
& \times\{v(X) w(Y)+v(Y) w(X)\}=0
\end{aligned}
$$

Putting $X=V$ and $Y=W$ in the last equation and using (2.16) and (4.15), we obtain

$$
v\left(A_{1} V\right)=w\left(A_{1} W\right)
$$

Similarly we have

$$
u\left(A_{1} U\right)=v\left(A_{1} V\right)=w\left(A_{1} W\right)
$$

which together with (4.17) gives

$$
A_{1} \varphi=\varphi A_{1}
$$

By the quite similar method we have

$$
A_{1} \varphi=\varphi A_{1}, \quad A_{1} \psi=\psi A_{1}, \quad A_{1} \theta=\theta A_{1}
$$

which yields the implication $(\mathrm{b}) \Longrightarrow(\mathrm{e})$.
Combining Theorem 1 with Theorem K-P, we have

Theorem 2. Let $M$ be an $n$-dimensional $Q R$-submanifold of $Q R$-dimension ( $p-1$ ) in $Q P^{(n+p) / 4}(4)$ and let the normal vector field $\xi$ be parallel with respect to the normal connection. If one of the conditions (a)-(e) stated in Theorem 1 is valid on $M$, then $\pi^{-1}(M)$ is locally a product $M_{1} \times M_{2}$ where $M_{1}$ and $M_{2}$ lie on some $\left(4 n_{1}+3\right)$ - and $\left(4 n_{2}+3\right)$-dimensional spheres, respectively ( $\pi$ is the Hopf fibration $\left.S^{n+p+3}(1) \rightarrow Q P^{(n+p) / 4}(4)\right)$.

## 5. The special case of an ambient quaternionic Kähler manifold

In this section we specify the ambient manifold $\bar{M}$ as a quaternionic space form $\bar{M}^{(n+p) / 4}(c)$ with $c=0$ and assume that one of the conditions (a)-(e) stated in Theorem 1 is valid on $M$. Then Theorem 1 implies

$$
\begin{equation*}
A_{1} \varphi=\varphi A_{1}, \quad A_{1} \psi=\psi A_{1}, \quad A_{1} \theta=\theta A_{1} \tag{5.1}
\end{equation*}
$$

from which, taking account of (2.9) and (2.11), we have

$$
A_{1} U=\lambda U, \quad A_{1} V=\mu V, \quad A_{1} W=\nu W
$$

where $\lambda=u\left(A_{1} U\right), \mu=v\left(A_{1} V\right), \nu=w\left(A_{1} W\right)$. But, applying $\psi$ to the first equation of (5.1) and using (2.13) and (5.1) itself, we have

$$
u(X) A_{1} V=u\left(A_{1} X\right) V
$$

from which, putting $X=U$, it follows that

$$
A_{1} V=\lambda V
$$

and, consequently, $\lambda=\mu$. Similarly we $\lambda=\mu=\nu$ which yields

$$
\begin{equation*}
A_{1} U=\lambda U, \quad A_{1} V=\lambda V, \quad A_{1} W=\lambda W \tag{5.2}
\end{equation*}
$$

Differentiating the first equation of (5.2) covariantly and using (3.7), (5.1) and (5.2) itself, we have

$$
g\left(\left(\nabla_{X} A_{1}\right) Y, U\right)+g\left(\varphi A_{1}^{2} X, Y\right)=(X \lambda) u(Y)+\lambda g\left(\varphi A_{1} X, Y\right)
$$

from which, taking the skew-symmetric part and making use of (3.10) with $c=0$ and (5.1), it follows that

$$
\begin{equation*}
2 g\left(\varphi A_{1}^{2} X, Y\right)=(X \lambda) u(Y)-(Y \lambda) u(X)+2 \lambda g\left(\varphi A_{1} X, Y\right) \tag{5.3}
\end{equation*}
$$

Now we put $Y=U$ in (5.3). Then the skew-symmetry of $\varphi$ and (2.11) imply $X \lambda=(U \lambda) u(X)$. Similary we have

$$
X \lambda=(U \lambda) u(X)=(V \lambda) v(X)=(W \lambda) w(X)
$$

and consequently $U \lambda=V \lambda=W \lambda=0$ which yield that $\lambda$ is constant. Combining this fact with (5.3) gives $\varphi\left(A_{1}^{2} X-\lambda A_{1} X\right)=0$, from which, applying $\varphi$ and using (2.9) and (5.2), we obtain $A_{1}^{2}=\lambda A_{1}$. Thus we have

Lemma 5.1. Let $M$ be an $n$-dimensional $Q R$-submanifold of $Q R$-dimension $(p-1)$ in a quaternionic space form $\bar{M}^{(n+p) / 4}(c)$ with $c=0$ such that the distinguished normal vector field $\xi$ is parallel with respect to the normal connection. If one of the conditions (a)-(e) stated in Theorem 1 is valid on $M$, then

$$
\begin{equation*}
A_{1}^{2}=\lambda A_{1} \tag{5.4}
\end{equation*}
$$

and $\lambda$ is constant.
In particular, we can prove

Lemma 5.2. Let $M$ be as in Lemma 5.1. Then

$$
\begin{equation*}
\nabla A_{1}=0 \tag{5.5}
\end{equation*}
$$

provided $\lambda \neq 0$.
Proof. Differentiating (5.4) covariantly and using the fact that $\lambda$ is constant, we have

$$
\begin{equation*}
\left(\nabla_{Y} A_{1}\right) A_{1} X+A_{1}\left(\nabla_{Y} A_{1}\right) X=\lambda\left(\nabla_{Y} A_{1}\right) X \tag{5.6}
\end{equation*}
$$

from which, taking the skew-symmetric part and using (3.10) with $c=0$, we find

$$
\left(\nabla_{Y} A_{1}\right) A_{1} X=\left(\nabla_{X} A_{1}\right) A_{1} Y
$$

and, consequently,

$$
g\left(\left(\nabla_{Y} A_{1}\right) A_{1} X, Z\right)=g\left(\left(\nabla_{X} A_{1}\right) A_{1} Y, Z\right)=g\left(A_{1}\left(\nabla_{X} A_{1}\right) Z, Y\right)
$$

On the other hand

$$
g\left(\left(\nabla_{Y} A_{1}\right) A_{1} X, Z\right)=g\left(\left(\nabla_{Z} A_{1}\right) A_{1} X, Y\right)
$$

which together with the last equation gives

$$
g\left(\left(\nabla_{Y} A_{1}\right) A_{1} X, Z\right)=g\left(A_{1}\left(\nabla_{X} A_{1}\right) Y, Z\right)
$$

that is, $\left(\nabla_{Y} A_{1}\right) A_{1} X=A_{1}\left(\nabla_{Y} A_{1}\right) X$. Hence (5.6) reduces to

$$
2 A_{1}\left(\nabla_{Y} A_{1}\right) X=\lambda\left(\nabla_{Y} A_{1}\right) X
$$

from which, applying $A_{1}$ and using (5.4), it is clear that

$$
\lambda A_{1}\left(\nabla_{Y} A_{1}\right) X=0
$$

and therefore $\lambda\left(\nabla_{Y} A_{1}\right) X=0$. Thus we complete the proof.
Remark 5.1. When the ambient space is a quaternionic projective space $Q P^{(n+p) / 4}$, the assumptions stated in Lemma 5.1 yield that the shape operator $A_{1}$ is cyclic-parallel, that is,

$$
\left.\left.\left.g\left(\nabla_{X} A_{1}\right) Y, Z\right)+g\left(\nabla_{Y} A_{1}\right) Z, X\right)+g\left(\nabla_{Z} A_{1}\right) X, Y\right)=0
$$

But, in this case we don't need the hypothesis $\lambda \neq 0$. (For details, see [9].)

## 6. The main results when $\bar{M}=Q^{(n+p) / 4}$

In this section we specialize to the case of an ambient quaternionic number space $Q^{(n+p) / 4}$. In this case, as already shown in Lemma 5.1, the eigenvalues $\kappa$ of the shape operator $A_{1}$ satisfy

$$
\kappa(\kappa-\lambda)=0 .
$$

Moreover it is clear from (5.1) and (5.2) that the multiplicity of $\lambda$ must be $4 m+3$ for some integer $m$ at each point in $M$. Since $\lambda$ is constant and trace $A_{1}$ is continuous, the multiplicity $r$ of $\lambda$ is constant. Hence it suffices to consider the following three cases

$$
\text { (i) } r=0, \quad \text { (ii) } r=n, \quad \text { (iii) } 3 \leqslant r<n \text {. }
$$

We will start with the first case (i). In this case $A_{1}=0$. Since, by assumption, the normal vector field $\xi$ is parallel with respect to the normal connection, Erbacher's reduction theorem ([4]) yields that there exists a totally geodesic hypersurface $R^{n+p-1}$ in $Q^{(n+p) / 4}$ which contains $M$.

Next, we consider the case (ii). In this case $A_{1}=\lambda I$. Let $\bar{x}$ be the position vector of $M$ and put $\bar{p}:=\bar{x}+\lambda^{-1} \xi$. Then

$$
\bar{\nabla}_{X} \bar{p}=\bar{\nabla}_{X}\left(\bar{x}+\lambda^{-1} \xi\right)=X-\lambda^{-1}\left(A_{1} X-\nabla_{X}^{\perp} \xi\right)=0,
$$

which means that $\bar{p}$ is a fixed point in $Q^{(n+p) / 4}$. Moreover, it is clear that $\| \bar{x}-$ $\bar{p} \|=|\lambda|^{-1}$ and consequently $M$ is contained in the hypersphere $S^{n+p-1}\left(|\lambda|^{-1}\right)$ of radius $|\lambda|^{-1}$ centered at $\bar{p}$.

Finally we consider the case (iii). Since the multiplicity $r$ of $\lambda$ is constant, the eigenspaces corresponding to $\lambda$ and 0 determine distributions of dimension $r$ and $n-r$, which will be denoted by $D_{\lambda}$ and $D_{0}$, respectively. Furthermore, by means of Lemma 5.2, $\nabla A_{1}=0$ and consequently it is easily verified that $D_{\lambda}$ and $D_{0}$ are both involutive and that $D_{\lambda}$ is parallel along $D_{0}$ and vice versa. Denoting by $M_{\lambda}$ and $M_{0}$ the integral submanifolds of $D_{\lambda}$ and $D_{0}$, respectively, we can see that $M$ is locally the Riemannian product $M_{\lambda} \times M_{0}$.

From now on we shall study $M_{\lambda}$ and $M_{0}$ in more detail and start with $M_{\lambda}$. Let $Z_{1}, \ldots, Z_{n-r}$ be orthonormal vector fields belonging to $D_{0}$. Since $M_{\lambda}$ is totally geodesic in $M$, the shape operators $A_{1}^{\prime}, \ldots, A_{n-r}^{\prime}$ corresponding to those normal vectors vanish. On the other hand we may consider $M_{\lambda}$ as a submanifold of $Q^{(n+p) / 4}$. Then the vector fields $Z_{1}, \ldots, Z_{n-r}, \xi_{1}, \ldots, \xi_{p}$ form an orthonormal set of local vector fields normal to $M_{\lambda}$. In this case the shape operators corresponding to $Z_{1}, \ldots, Z_{n-r}$ also vanish. Hence it is clear from (3.11) that

$$
\begin{equation*}
{ }^{\prime} R_{X, Y}^{\perp} Z_{i}=0, \quad i=1, \ldots, n-r \tag{6.1}
\end{equation*}
$$

and moreover $\left[A_{1}, A_{\alpha}\right]=0$, where ' $R^{\perp}$ denotes the curvature tensor of the normal connection ${ }^{\prime} \nabla^{\perp}$ of $M_{\lambda}$ in $Q^{(n+p) / 4}$. On the other hand, we can easily see that for any $X \in D_{\lambda}$

$$
g\left(\nabla_{X}^{\perp} Z_{i}, \xi_{\beta}\right)=g\left(Z_{i}, A_{\beta} X\right), \quad \beta=1, \ldots, p .
$$

But, since $\left[A_{1}, A_{\beta}\right]=0, \beta=1, \ldots, p$, which is a direct consequence of (3.11) and $\nabla^{\perp} \xi_{1}=0$, we have $A_{\beta} X \in D_{\lambda}$ and, consequently,

$$
g\left({ }^{\prime} \nabla_{X}^{\perp} Z_{i}, \xi_{\beta}\right)=0, \quad \beta=1, \ldots, p
$$

that is, ${ }^{\prime} \nabla \frac{1}{X} Z_{i} \in D_{0}$. Thus, by the same method as in the proof of Proposition 1.1 in [3, p. 99], we may prove that (6.1) yields the existence of the normal vector fields $Z_{1}, \ldots, Z_{n-r}$ such that

$$
\begin{equation*}
{ }^{\prime} \nabla_{X}^{\perp} Z_{i}=0, \quad i=1, \ldots, n-r \tag{6.2}
\end{equation*}
$$

for any tangent vector field $X$ to $M_{\lambda}$.

Now let $\bar{x}$ be the position vector of $M_{\lambda}$ in $Q^{(n+p) / 4}$ and $X \in D_{\lambda}$. Then, by using (6.2) and $A_{i}^{\prime}=0, i=1, \ldots, n-r$, we have

$$
X g\left(\bar{x}, Z_{i}\right)=g\left(X, Z_{i}\right)=0, \quad i=1, \ldots, n-r
$$

that is,

$$
\begin{equation*}
g\left(\bar{x}, Z_{i}\right)=c_{i}, \quad i=1, \ldots, n-r \tag{6.3}
\end{equation*}
$$

where $c_{i}$ is constant. Moreover, putting $\bar{p}:=\bar{x}+\lambda^{-1} \xi$, we can see that

$$
\bar{\nabla}_{X} \bar{p}=X-\lambda^{-1} A_{1} X=0
$$

and $\|\bar{x}-\bar{p}\|=|\lambda|^{-1}$. Therefore $M_{\lambda}$ belongs to the intersection of the hypersphere of radius $|\lambda|^{-1}$ centered at $\bar{p}$ and the $n-r$ hyperplanes defined by (6.3). We notice that $\bar{p}$ is contained in the $n-r$ hyperplanes.

In a similar way it can be shown that $M_{0}$ belongs to the intersection of the $r+1$ hyperplanes given by

$$
g(\bar{x}, \xi)=c, \quad g\left(\bar{x}, Z_{s}\right)=c_{s}, \quad s=n-r+1, \ldots, n
$$

Summing up, we may conclude

Theorem 2. Let $M$ be an $n$-dimensional $Q R$-submanifold of $Q R$-dimension $(p-1)$ in $Q^{(n+p) / 4}$ which satisfies one of the conditions stated in Theorem 1. If the distinguished normal vector field $\xi$ is parallel with respect to the normal connection, then we have one of the following cases:
(a) $M$ is contained in a hyperplane orthogonal to $\xi$.
(b) $M$ is contained in a hypersphere orthogonal to $\xi$.
(c) $M$ is locally a Riemannian product $M_{\lambda} \times M_{0}$, where $M_{\lambda}$ is contained in a $(p+r-1)$-dimensional sphere $S^{(p+r-1)}$ and $M_{0}$ is contained in an $(n+p-r-1)$ dimensional subspace $R^{(n+p-r-1)}$.

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