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# Lower bounds for simultaneous Diophantine approximation constants ${ }^{1}$ 

Werner Georg Nowak


#### Abstract

After a brief exposition of the state-of-art of research on the (Euclidean) simultaneous Diophantine approximation constants $\theta_{s}$, new lower bounds are deduced for $\theta_{6}$ and $\theta_{7}$.


## 1 Introduction

For a fixed positive integer $s$, the (Euclidean) simultaneous Diophantine approximation constant $\theta_{s}$ is defined as the supremum of all constants $c$ such that, for every point $\mathbf{a}$ in $\mathbb{R}^{s} \backslash \mathbb{Q}^{s}$, there exist infinitely many $(s+1)$-tuples $(\mathbf{p}, q) \in \mathbb{Z}^{s} \times \mathbb{N}^{*}$ with

$$
\begin{equation*}
\left|\mathbf{a}-\frac{1}{q} \mathbf{p}\right| \leq \frac{1}{q \sqrt[s]{c q}}, \tag{1}
\end{equation*}
$$

where $|\cdot|$ denotes the Euclidean norm in $\mathbb{R}^{s}$.
This notion generalizes a question whose answer is known as Hurwitz' classic theorem. This involves the special case $s=1$ and tells us that $\theta_{1}=\sqrt{5}$; see, e.g., Niven and Zuckerman 11, p. 189 and p. 221].

By some very deep analysis, Davenport and Mahler 6] were able to prove that $\theta_{2}=\frac{1}{2} \sqrt{23}$.

For $s \geq 3$, the exact values of $\theta_{s}$ are unknown, and only more or less precise bounds have been established.

We remark parenthetically that the problem becomes even considerably more difficult if one replaces in (1) the Euclidean norm by the maximum norm: The constants arising, say $\theta_{s}^{(\infty)}$, are unknown for all $s \geq 2$, the only general successful approach being due to Spohn 16 who combined the calculus of variation with a classic method of Blichfeldt 2 to estimate $\theta_{s}^{(\infty)}$ from below.

[^0]
## 2 Survey of methods and known results

The usual approach to estimate $\theta_{s}$ is based on tools from the geometry of numbers; cf. throughout the monograph by Gruber and Lekkerkerker 7 .

We briefly recall a few basic concepts of this theory: For a star-body $\mathcal{B}$ in $\mathbb{R}^{s}$, a lattice $\Lambda=A \mathbb{Z}^{s}$ ( $A$ a nonsingular real $(s \times s)$-matrix) is called admissible if its only point in the interior of $\mathcal{B}$ is the origin. The critical determinant $\Delta(\mathcal{B})$ of $\mathcal{B}$ is then defined as the infimum of the lattice constants $d(\Lambda)=|\operatorname{det} A|$, taken over all lattices $\Lambda$ admissible for $\mathcal{B}$.

By a celebrated result of Davenport 5 (see also 7, p. 480, Theorem 4]), $\theta_{s}$ is equal to the critical determinant $\Delta\left(\mathcal{B}_{s+1}\right)$ of the $(s+1)$-dimensional star body of points $\left(x_{0}, x_{1}, \ldots, x_{s}\right) \in \mathbb{R}^{s+1}$

$$
\begin{equation*}
\mathcal{B}_{s+1}: \quad\left|x_{0}\right|\left(\sum_{j=1}^{s} x_{j}^{2}\right)^{s / 2} \leq 1 \tag{2}
\end{equation*}
$$

For $s=3,4,5$, the sharpest known lower ${ }^{2}$ estimates for $\theta_{s}$ have been obtained by a method which is based on inequalities relating the critical determinants of star bodies in different dimensions. In its essence it goes back to Mordell 8, 9, 10, and Armitage 1. Combining these tools with the known critical determinant $\overline{\Delta(\mathcal{P})}=\frac{1}{2}$ of the three-dimensional double paraboloid

$$
\begin{equation*}
\mathcal{P}: \quad x^{2}+y^{2}+|z| \leq 1, \tag{3}
\end{equation*}
$$

it has been proved 13, 14 that

$$
\begin{equation*}
\theta_{3} \geq 1.879 \ldots, \quad \theta_{4} \geq 1.3225 \ldots, \quad \theta_{5} \geq 0.876 \tag{4}
\end{equation*}
$$

For $s \leq 5$, it appears that this is the very limit of present methods. For $s \geq 6$, the only successful approach is due to Prasad 15. This is based on the simple idea to apply the arithmetic-geometric mean inequality to the left-hand side of (2). In terms of geometry, this amounts to inscribing an $(s+1)$-dimensional ellipsoid

$$
\begin{equation*}
\mathcal{E}_{s+1}: \quad \frac{1}{s+1} x_{0}^{2}+\frac{s}{s+1} \sum_{j=1}^{s} x_{j}^{2} \leq 1 \tag{5}
\end{equation*}
$$

into $\mathcal{B}_{s+1}$. It follows that

$$
\begin{equation*}
\theta_{s}=\Delta\left(\mathcal{B}_{s+1}\right) \geq \Delta\left(\mathcal{E}_{s+1}\right)=\frac{(s+1)^{(s+1) / 2}}{s^{s / 2}} \Delta\left(\mathcal{S}_{s+1}\right) \tag{6}
\end{equation*}
$$

where $\mathcal{S}_{s+1}$ is the unit sphere in $\mathbb{R}^{s+1}$. Now the critical determinants of the unit spheres are known up to dimension 8: See 7. p. 410]; in particular, $\Delta\left(\mathcal{S}_{7}\right)=\frac{1}{8}$, $\Delta\left(\mathcal{S}_{8}\right)=\frac{1}{16}$. Hence, using (6), it readily follows that

$$
\begin{equation*}
\theta_{6} \geq \frac{343}{1728} \sqrt{7}, \quad \theta_{7} \geq \frac{256}{343} \frac{1}{\sqrt{7}} . \tag{7}
\end{equation*}
$$

[^1]We conclude this section by the remark that the question of upper bounds for $\theta_{s}$ has been dealt with in 14 , section 4]. It amounts to finding certain number fields of degree $s+1$ with small absolute discriminant.

## 3 Improvement of the estimate (7)

The critical determinants $\Delta\left(\mathcal{S}_{7}\right), \Delta\left(\mathcal{S}_{8}\right)$ once known, the deduction of the lower bounds (7) seems so natural that one might believe that this could be the end-of--the-art for this problem in the cases $s=6,7$. In this little note, however, we will establish a slight refinement.

Theorem 1. The inequalities

$$
\theta_{6} \geq \frac{343}{1728} \sqrt{7}\left(1+\omega_{6}\right), \quad \theta_{7} \geq \frac{256}{343} \frac{1}{\sqrt{7}}\left(1+\omega_{7}\right)
$$

hold true, with certain small constants $\omega_{6}>9 \times 10^{-4}$, $\omega_{7}>3 \times 10^{-4}$. I.e., numerically, $\theta_{6} \geq 0.52564, \theta_{7} \geq 0.28218$.

Of course, this improvement is fairly small, the main interest lying in the method applied. This in turn is inspired by classic work due to Davenport 3, 4, and Žilinskas 17, as well as by an earlier article by the author 12$]^{3}$.

## 4 Proof of the theorem

For better readability, we give the details only for $s=6$, the case $s=7$ being completely analogous. In principle, the argument can be extended to $s>7$ as well, but this is of less importance, since $\Delta\left(\mathcal{S}_{s+1}\right)$ is known for $s \leq 7$ only.

The star body $\mathcal{B}_{7}$ defined in (2) is automorphic, hence there exists a critical lattic $\S^{4} \Lambda$ with a point on the boundary of $\mathcal{B}_{7} ;$ cf. 7 7, p. 305, Theorem 4]. Applying to $\Lambda$ a suitable automorphism of $\mathcal{B}_{7}$, if necessary, we can assume this point to be $\mathbf{e}=(1,1,0,0,0,0,0)$. With $\mathbf{x} \in \mathbb{R}^{7}$, the function

$$
G(\mathbf{x}):=\left(\frac{1}{7} x_{0}^{2}+\frac{6}{7} \sum_{j=1}^{6} x_{j}^{2}\right)^{1 / 2}
$$

which is the square-root of the left-hand side of (5) for $s=6$, is called the distance function of the ellipsoid $\mathcal{E}_{7}$; it is homogeneous of order 1 . Since, according to 7 p. 195, Theorem 3], any o-symmetric ellipsoid has anomaly 1 , there exist seven linearly independent lattice points $\mathbf{u}^{(k)}$ of $\Lambda$, with $\left(G\left(\mathbf{u}^{(k)}\right)\right)_{k=1}^{7}$ nondecreasing, and

$$
\begin{equation*}
\Delta\left(\mathcal{E}_{7}\right) \prod_{k=1}^{7} G\left(\mathbf{u}^{(k)}\right) \leq d(\Lambda)=\Delta\left(\mathcal{B}_{7}\right)=\theta_{6} \tag{8}
\end{equation*}
$$

[^2]We pick $\mathbf{u}=\left(u_{0}, u_{1}, \ldots, u_{6}\right) \in\left\{\mathbf{u}^{(1)}, \mathbf{u}^{(2)}\right\}$ in such a way that $\mathbf{u} \neq \pm \mathbf{e}$. Then $\mathbf{u} \pm \mathbf{e}$ are nontrivial lattice points of $\Lambda$. Since $\Lambda$ is admissible for $\mathcal{B}_{7}$, a look back to (2) shows that

$$
\begin{equation*}
\left|u_{0}\right|\left(\sum_{j=1}^{6} u_{j}^{2}\right)^{3} \geq 1, \quad\left|u_{0} \pm 1\right|\left(\left(u_{1} \pm 1\right)^{2}+\sum_{j=2}^{6} u_{j}^{2}\right)^{3} \geq 1 \tag{9}
\end{equation*}
$$

Since $\mathcal{E}_{7} \subset \mathcal{B}_{7}$, it follows that $G\left(\mathbf{u}^{(1)}\right) \geq 1$, hence (8) implies that

$$
\begin{equation*}
\frac{\theta_{6}}{\Delta\left(\mathcal{E}_{7}\right)} \geq(G(\mathbf{u}))^{6} \tag{10}
\end{equation*}
$$

To prove the Theorem, it remains to minimize $G(\mathbf{u})$ under the constraints (9). We put $S=\sum_{j=2}^{6} u_{j}^{2}$ for short, and may assume, w.l.o.g., that $u_{0}>0$. Hence we have to deal with a minimization problem in three variables only, namely $u_{0}, u_{1}$ and $S$. In fact,

$$
\begin{equation*}
M:=\min _{[9]} G^{2}(\mathbf{u})=\min _{[12]}\left(\frac{1}{7} u_{0}^{2}+\frac{6}{7}\left(u_{1}^{2}+S\right)\right), \tag{11}
\end{equation*}
$$

with

$$
\begin{equation*}
u_{0}\left(u_{1}^{2}+S\right)^{3} \geq 1, \quad\left|u_{0} \pm 1\right|\left(\left(u_{1} \pm 1\right)^{2}+S\right)^{3} \geq 1 \tag{12}
\end{equation*}
$$

Solving (12), we infer that

$$
u_{1}^{2}+S \geq \max \left(u_{0}^{-1 / 3},\left(u_{0}+1\right)^{-1 / 3}-1-2 u_{1},\left|u_{0}-1\right|^{-1 / 3}-1+2 u_{1}\right)
$$

Hence, in view of (11),

$$
\begin{array}{r}
7 M=\min _{u_{0}>0, u_{1}}\left(\operatorname { m a x } \left(u_{0}^{2}+6 u_{0}^{-1 / 3}, u_{0}^{2}+6\left(u_{0}+1\right)^{-1 / 3}-6-12 u_{1}\right.\right. \\
\\
\left.\left.u_{0}^{2}+6\left|u_{0}-1\right|^{-1 / 3}-6+12 u_{1}\right)\right) .
\end{array}
$$

Keeping $u_{0}$ fixed for the moment and seeking the minimum with respect to $u_{1}$, we observe that the maximum of the last two expressions becomes minimal when they are equal. This obviously happens for

$$
u_{1}=\frac{1}{4}\left(\left(u_{0}+1\right)^{-1 / 3}-\left|u_{0}-1\right|^{-1 / 3}\right) .
$$

Consequently,

$$
\begin{equation*}
7 M=\min _{u_{0}>0}\left(\max \left(u_{0}^{2}+6 u_{0}^{-1 / 3}, u_{0}^{2}+3\left(u_{0}+1\right)^{-1 / 3}+3\left|u_{0}-1\right|^{-1 / 3}-6\right)\right) \tag{13}
\end{equation*}
$$

In order to solve this ultimate minimization problem, the figure below is very helpful.


Figure 1: The graphs of $f_{1}: u_{0} \mapsto u_{0}^{2}+6 u_{0}^{-1 / 3}$ (dashed)

$$
\text { and } f_{2}: u_{0} \mapsto u_{0}^{2}+3\left(u_{0}+1\right)^{-1 / 3}+3\left|u_{0}-1\right|^{-1 / 3}-6
$$

In fact, let $f_{1}, f_{2}$ be defined as in the graphics, then an easy calculus exercise shows that $f_{1}$ decreases on $] 0,1[$ and increases on $] 1, \infty\left[\right.$, and that $f_{2}$ increases on $] 0,1\left[\right.$. Furthermore, $f_{1}-f_{2}$ increases on $] 1, \infty[$, since there

$$
\frac{\mathrm{d}}{\mathrm{~d} u_{0}}\left(f_{1}\left(u_{0}\right)-f_{2}\left(u_{0}\right)\right)=\frac{1}{\left(u_{0}+1\right)^{4 / 3}}+\frac{1}{\left(u_{0}-1\right)^{4 / 3}}-\frac{2}{u_{0}^{4 / 3}}>0
$$

on applying the mean inequality to the first two fractions. It readily follows that the equation $f_{1}\left(u_{0}\right)=f_{2}\left(u_{0}\right)$ has exactly two solutions, $u_{0}^{(1)}<1$ and $u_{0}^{(2)}>1$, say, and that, recalling (13),

$$
7 M=\min \left(f_{1}\left(u_{0}^{(1)}\right), f_{1}\left(u_{0}^{(2)}\right)\right)
$$

Carrying out the numerics, we get $u_{0}^{(1)}=0.97012 \ldots, u_{0}^{(2)}=1.030799 \ldots$, hence,

$$
7 M=\min (7.002111 \ldots, 7.00218 \ldots) \geq 7.002111
$$

Going back to (10) and 11), we finally infer that

$$
\frac{\theta_{6}}{\Delta\left(\mathcal{E}_{7}\right)} \geq\left(\frac{7.002111}{7}\right)^{3} \geq 1.0009
$$

which completes the proof of the Theorem.

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[^1]:    ${ }^{2}$ Obviously, lower bounds are the more interesting ones, since they guarantee, for every $c$ less than the bound, the existence of infinitely many solutions of the inequality (1).

[^2]:    ${ }^{3}$ Carrying out the numerical details on the basis of the argument developed in that latter paper, one would get only $\omega_{6}>6 \times 10^{-5}, \omega_{7}>1.5 \times 10^{-5}$.
    ${ }^{4}$ I.e., $\Lambda$ is admissible for $\mathcal{B}_{7}$, and $d(\Lambda)=\Delta\left(\mathcal{B}_{7}\right)$.

