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ON TIIE BEHAVIOUR OF THE GENERALIZED SOLUTION OF THE BIHARMONIC EQUATION IN TWO DIMENSIONS NEAR SINGULAR POINT OF THE BOUNDARY

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The results of this contribution were obtained commonly by 0 . A. Olejnik, V. A. Kondratjev from Moscow University and me and will be published in [1].

1. Notations and definitions

Let $\Omega \subset R^{2}$ be a bounded domain, $\Gamma \subset \delta \Omega$. By $H^{2}(\Omega, F)$ we denote the closure in the norm $\|u\|^{2}=\int_{a} \sum_{k \mid E 2}\left|D^{\alpha} u\right|^{2} d x$ of the set of functions from $C^{\infty}(\bar{\Omega})$, vanishing in $\frac{\gamma}{}$ some solution of the problem
(1) $\Delta_{u}^{2}=f$ in $\Omega, u=\delta_{u} / \delta_{n}=0$ on $\Gamma$
we understand the function $u \in H^{2}(\Omega, \Gamma)$, satisfying the identity
(1') $\int_{\Omega} \sum_{|\alpha|=2} D^{\alpha} u \cdot D^{\alpha} v d x=\int_{\Omega} f \cdot v d x$ for all $v \in H^{2}(\Omega, \delta \Omega)$.

## 2. The main theorem

Theorem. For $\Omega, \Gamma$ defined above let 0 be a point from $\Gamma$. Suppose that for some $T>0$ the assumptions
i) $\delta \Omega \cap\left\{x \in R^{2},\|x-0\|=t\right\} \neq \emptyset$ for all $t \in(0, T)$,
ii) $\delta \Omega \cap \delta \Omega_{T} \subset \Gamma$
hold, where $\Omega_{t}=\Omega \cap\{x,\|x-0\|<t\}$. Then there exist constants $C_{i}$ depending on $T$ (but not on $\Omega$ and $u$ ) such that for arbitrary generalized solution $u$ of (1) with $f=0$ in $\Omega_{T}$ following estimates hold
(3) $\int_{\Omega_{t}} E(u) d x \leq C_{1} \cdot t^{0,304} \int_{\Omega_{T}} E(u) d x$ for $t \in(0, T)$,
(4) $|u(x)|^{2} \leq c_{2}\|x-0\|^{2,304} \int_{\Omega} E(u) d x$ for $x \in \Omega,\|x-0\| \leq T / 2$, where $E(u)=u_{x_{1} x_{1}}^{2}+2 u_{x_{1} x_{2}}^{2}+{u_{x_{2} x_{2}}^{2}}_{2^{2}}^{2}$
If in addition to above assumptions $\Omega$ satisfies
iii) for each $x^{0} \in \Gamma, \| x^{0}-0 \mid \leq T / 2$ the circle $\left\|x-x^{0}\right\|=t$ intersects $\Gamma$ for all $t \in\left(0,\left\|x^{0}-0\right\| / 2\right\rangle$,
then the estimate $|\operatorname{grad} u(x)|^{2} \leq c_{3}\|x-0\|^{0,304} \int_{\Omega_{T}} E(u) d x$ holds too.
3. The sketch of the proof of the theorem

We suppose $0=(0,0)$ and introduce polar coordinates $r$ and $\varphi$. a) for $\phi \in C^{1}(0, T\rangle, \phi^{\prime} \leq 0, \phi^{\prime \prime}$ piecewise continuous on $(0, T)$, $\phi(t)=1, \phi^{\prime}(T)=0 \operatorname{let} \tilde{\phi}(r, t)=\phi(r)$ for $t \leq r \leq T, \tilde{\phi}=1$ for $r \geqslant T, \tilde{\phi}_{\text {linear }}$ on ( $0, t$ ) and in $C^{1}(0, T)$. Substituting in (1') $v=u$. ( $\bar{\phi}(r, t)-1)$, we get after some integrations by parts the identity

$$
\begin{align*}
& \int_{\Omega_{t}}\left\{\left[\phi(t)+(r-t) \cdot \phi^{\prime}(t)\right] \cdot E(u)+r^{-1}\left|\phi^{\prime}(t)\right| \cdot K_{2}(u)\right\} d x  \tag{5}\\
& +\int_{\Omega_{T^{\prime}} \Omega_{t}}\left[\phi(r) \cdot E(u)-\phi^{\prime \prime}(r) \cdot K_{1}(u)-r^{-1} \phi^{\prime}(r) \cdot K_{2}(u)\right] d x \leqslant_{\Omega_{T}} E(u) d x
\end{align*}
$$

with $K_{1}(u)=u_{r}^{2}+\left(r^{-1} u_{\rho}\right)^{2}-u \cdot u_{r r}, K_{2}(u)=u_{r}^{2}+2\left(r^{-1} u_{\rho}\right)^{2}-r^{-1} u \cdot u_{r}$.
b) If $\boldsymbol{\phi}$ moreover satisfies

$$
\begin{equation*}
\int_{S_{r}}\left[E(u) \phi(r)-\phi^{\prime \prime}(r) \cdot K_{1}(u)-r^{-1} \phi^{\prime}(r) \cdot K_{2} u\right] d \varphi \geq 0 \tag{6}
\end{equation*}
$$

for all $r \in(0, T)$ and all $u \in C^{2}$ in a neighbourhood of $\tilde{S}_{r}$, vanishing with first derivatives in points of $\delta \Omega \cap \delta \Omega_{T}$ by $S_{r}$ we denote the set of $\varphi$ such that the point $x=(r \cos \varphi, r \sin \varphi)$ belongs to $\Omega$, by $\tilde{S}_{\dot{r}}$ the set of all such points $x$ ) we get from (5) using $\phi, \leqslant 0, \phi \geq 1$, $\int_{r} K_{2}(u) d \mathcal{I} \geq 0$ the estimate
(7)

$$
\int_{\Omega_{t}} E(u) d x \leq(\phi(t))^{-1} \int_{\Omega_{T}} E(u) d x
$$

c) Now we shall construct the function $\phi$. Denoting $z=r \cdot u_{r r}$, $w=r^{-1} \cdot u, v=u_{r}$, we see, that $\phi$ will satisfy (6) if it satisfies
(8)

$$
\int_{0}^{2 \pi}\left(\varepsilon(w, v, z) \cdot \phi(r)-r^{2} \phi \cdot \mathcal{K}_{1}(w, v, z)-r\right.
$$

- $\left.\phi^{\prime} \varkappa_{2}(w, v, 2)\right) d \rho \geq 0$
for all $r \in(0, T)$ and all $w \in \dot{H}^{0}(0,2 K), v \in \AA^{1}(0,2 T), z \in L_{2}(0,2 T)$ where $E(w, v, z)=z^{2}+\left(w^{\prime \prime}+v\right)^{2}+2\left(v^{\prime}-w^{\prime}\right)^{2}, \mathcal{K}_{1}=v^{2}+\left(w^{\prime}\right)^{2}-w \cdot z$,
$\mathcal{K}_{2}=v^{2}+2\left(w^{\prime}\right)^{2}-v \cdot w$. (We use that $s_{r} \subset(0,2 \pi)$ and the equality $\left.E(u)=u_{r r}^{2}+r^{-2}\left(\left(r^{-1} \cdot u_{\rho \rho}+u_{r}\right)^{2}+2\left(u_{r \rho}-r^{-1} \cdot u_{f}\right)^{2}\right).\right)$
For fixed $w, v, z$ with $\mathcal{(} w, v, z)=1$ (8) will be satisfied by the function
(9) $\quad\left(\beta(r / T)^{-\alpha}+\alpha(r / T)^{\beta}\right) /(\alpha+\beta)$
where
(9')

$$
\begin{aligned}
& \alpha=\frac{-\left(k_{1}-k_{2}\right)+\sqrt{\left(k_{1}-k_{2}\right)^{2}+4 k_{1}}}{2 k_{1}}, \\
& \beta=\frac{\left(k_{1}-k_{2}\right)+\sqrt{\left(k_{1}-k_{2}\right)^{2}+4 k_{1}}}{2 k_{1}}
\end{aligned}
$$

which solves the equation
(10) $r^{2} \phi^{\prime \prime}(r) k_{1}+r \phi^{\prime}(r) k_{2}=\phi, \phi(T)=1, \phi^{\prime}(T)=0$
with $k_{i}=\int_{0}^{2 \pi} \mathcal{K}_{i}(w, v, z) d \rho$. (We may assume $k_{1}>0$.)
The function (9) satisfies $\phi^{\prime} \leq 0, \phi^{\prime \prime} \geq 0, \phi^{\prime \prime}(r)+r^{-1} \phi^{\prime}(r) \geq 0$ for $0<r \leq T$. Using this it is easily seen that if $\phi$ satisfies (10) then it satisfies
(10') $r^{2} \phi^{\prime \prime}(r) \tilde{k}_{1}+r \phi^{\prime}(r) \tilde{k}_{2} \leq \phi(r)$
with $\tilde{k}_{1}, \tilde{k}_{2}$ satisfying one of the following pairs of conditions

$$
\begin{array}{ll}
\tilde{\mathbf{k}}_{1}=\mathrm{k}_{1}, & \tilde{\mathbf{k}}_{1} \leq k_{1}, \\
\tilde{\mathbf{k}}_{2} \geq \mathrm{k}_{2}, & \tilde{\mathrm{k}}_{2} \leq \mathbf{k}_{1}, \\
\mathrm{k}_{2}, & \tilde{\mathrm{k}}_{1}-\tilde{k}_{2}=k_{1}-\mathbf{k}_{2}
\end{array}
$$

Now we prove the estimates
$\max k_{1}(w, v, z)=\gamma_{1} \leq 3,2679$,
$\max \left(k_{1}(w, v, 2)-k_{2}(w, v)\right) \leq 2,29395$
where the max are taken on the set $\left\{w, v, z, \int_{0}^{2 T} \varepsilon(w, v, z) d f=1\right\}$. So the function (9) corresponding to $k_{1}=3,2679$ and $k_{2}=3,2679-$ 2,29395 satisfies (8) for all $\mathrm{w}, \mathrm{v}, \mathrm{z}$. Calculating the corresponding $\alpha$ from ( $9^{\prime}$ ), we obtain $\geq 0,304$. So the estimate (3) is proved.
(4) follows from (3) e.g. by imbedding theorem.

Remark 1. From results of [2], [3] it follows that for
$\Omega=\left\{x \in R^{2}, x=\left(x_{1}, x_{2}\right),\|x\| \neq x_{1}\right\}$ the estimate $|u(x)|^{2} \leq c\|x\|^{3}$
holds for arbitrary generalized solution considered in theorem proved above. It is an open problem if such estimate is valid for general $\Omega$ as in our theorem.

Remark 2. One may derive other sufficient conditions for $\phi$ to satisfy (6) e.g.
(I) $\phi(T)=1, \phi^{\prime}(T)=0, \phi^{\prime} \leq 0, \phi^{\prime \prime} \geq 0, \phi^{\prime \prime}(r)+r^{-1} \phi^{\prime}(r) \geq 0$, $0<r \leq T, 2(\ell(r) / \pi)^{2}\left(\phi^{\prime \prime}(r)+r^{-1} \phi^{\prime}(r)\right)+(\ell(r) / 2 \pi)^{2} \phi^{\prime \prime}(r)$ $+4 r(\ell(r) / 2 r \pi)^{3}\left|\phi^{\prime}\right| \leqslant \phi$
or
(II) instead of inequality from (I) one assume the inequality $(2+2 \sqrt{2})(\ell(r) / 2 \pi)^{2} \phi^{\prime \prime}(r) \leqslant \phi \quad(r), \quad 0<r \leq T$,
where $\ell(r)$ is the length of largest arc in $\tilde{S}_{\mathbf{r}}$.
These conditions yield in the case of general domain not so good estimates as in our theorem, but in special cases of $\Omega$ may give e.g. exponential decay of $\int_{\Omega_{t}} E(u) d x$ as $t \rightarrow 0$, since they exploit the
form of $\Omega$.
4. Further results
a) From above results one may derive uniqueness theorems for Dirichlet problem for $\Delta^{2} u=f$ in unbounded domains.
b) Analogously one may study the decay of energy or of the solution of (1) in unbounded domains as $\|x\| \rightarrow \infty$. One can prove e.g. that for the domain $\left.\Omega^{\omega}=\left\{x=\left(x_{1}, x_{2}\right)=(r \cos \rho, r \sin \}\right), \rho \in(0, \omega)\right\}$ the solution with finite energy tends to zero as $\|x\| \rightarrow \infty$ for $\omega$ sufficiently small. Since for $\omega=2 \pi$ this is not true, it is an open problem to find the largest $w$ for which this is true.

## References

[1] Olejnik, O. A., Kondratjev, V. A., Kopaček, J.: Differencialnyje uravněnija 1981, to appear.
[2] Kondratjev, V. A.: Prikl. mat. mech. 31 (1967), 119 - 123.
[3] Mazja, V. G.: DAN SSSR 235 (1977), 1263-1266.

