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ON THE BEHAVIOUR OF THE GENERALIZED SOLUTION OF THE BIHARMONIC EQUATION IN TWO DIMENSIONS NEAR SINGULAR POINT OF THE BOUNDARY

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The results of this contribution were obtained commonly by O. A. Olejnik, V. A. Kondratjev from Moscow University and me and will be published in [1].

1. Notations and definitions

Let $\Omega \subset \mathbb{R}^2$ be a bounded domain, $\Gamma \subset \delta \Omega$. By $H^2(\Omega, \Gamma)$ we denote the closure in the norm $\|u\|^2 = \int_{\Sigma} \sum_{u \in \Omega} |D^{u}u|^2 dx$ of the set of functions from $\mathbb{C}^{\infty}(\overline{\Omega})$, vanishing in some neighbourhood of Γ . By general: neighbourhood of Γ . By generalized solution of the problem (1) $\Delta^2 u = f in \Omega$, $u = \delta u/\delta n = 0$ on Γ we understand the function $u \in H^2(\Omega, \Gamma)$, satisfying the identity (1') $\int_{\Omega} \sum_{|\alpha|=2} D^{\alpha} u \cdot D^{\alpha} v \, dx = \int f \cdot v \, dx \text{ for all } v \in H^2(\Omega, \delta\Omega).$

2. The main theorem

Theorem. For Ω , Γ defined above let 0 be a point from Γ . Suppose that for some T > 0 the assumptions

i) $\delta \Omega \cap \{x \in \mathbb{R}^2, \|x - 0\| = t\} \neq \emptyset$ for all $t \in (0, T)$, ii) $\delta \Omega \cap \delta \Omega_T \subset \Gamma$

hold, where $\Omega_t = \Omega \cap \{x, \|x - 0\| < t\}$. Then there exist constants C, depending on T (but not on Ω and u) such that for arbitrary generalized solution u of (1) with f = 0 in Ω_{T} following estimates hold

(3)
$$\int_{\Omega_{t}} E(u) dx \leq C_{1} \cdot t^{0,304} \int_{\Omega_{T}} E(u) dx \text{ for } t \in (0, T),$$
(4)
$$|u(x)|^{2} \leq C_{2} ||x - 0||^{2,304} \int_{\Omega_{T}} E(u) dx \text{ for } x \in \Omega, ||x - 0|| \leq T/2,$$
where $E(u) = u_{x_{1}x_{1}}^{2} + 2 u_{x_{1}x_{2}}^{2} + u_{x_{2}x_{2}}^{2}.$
If in addition to above assumptions Ω satisfies

iii) for each
$$x^{\circ} \in \Gamma$$
, $\|x^{\circ} - 0\| \leq T/2$ the circle $\|x - x^{\circ}\| = t$
intersects Γ for all $t \in (0, \|x^{\circ} - 0\|/2)$,

then the estimate

(4')
$$| \text{grad } u(x) |^2 \leq C_3 ||x - 0||^{0,304} \int_T E(u) dx$$

holds too.

3. The sketch of the proof of the theorem

We suppose 0 = (0, 0) and introduce polar coordinates r and \P . a) for $\phi \in C^1(0, T)$, $\phi' \leq 0$, ϕ'' piecewise continuous on (0, T), $\phi(t) = 1, \phi'(T) = 0$ let $\tilde{\phi}(r, t) = \phi(r)$ for $t \le r \le T, \tilde{\phi} = 1$ for $r \ge T$, $\tilde{\phi}$ linear on (0, t) and in C¹(0, T). Substituting in (1') v = u. $(\vec{\phi}(\mathbf{r}, t) - 1)$, we get after some integrations by parts the identity (5) $\int_{\Omega_{+}} \left\{ \left[\phi(t) + (r - t) \cdot \phi'(t) \right] \cdot E(u) + r^{-1} |\phi'(t)| \cdot K_{2}(u) \right\} dx$ + $\int_{\Omega_{m}} \left[\phi(\mathbf{r}) \cdot E(\mathbf{u}) - \phi''(\mathbf{r}) \cdot K_1(\mathbf{u}) - \mathbf{r}^{-1} \phi'(\mathbf{r}) \cdot K_2(\mathbf{u}) \right] d\mathbf{x} \leq \int_{\Omega_{m}} E(\mathbf{u}) d\mathbf{x}$ with $K_1(u) = u_r^2 + (r^{-1}u_r)^2 - u \cdot u_{rr}$, $K_2(u) = u_r^2 + 2(r^{-1}u_r)^2 - r^{-1}u \cdot u_r$. b) If ϕ moreover satisfies (6) $\int_{S} \left[E(u) \phi(r) - \phi''(r) \cdot K_{1}(u) - r^{-1} \phi'(r) \cdot K_{2}u \right] df \ge 0$ for all $r \in (0, T)$ and all $u \in C^2$ in a neighbourhood of \widetilde{S}_r , vanishing with first derivatives in points of $\delta\Omega \cap \delta\Omega_T$ (by S_r we denote the set of f such that the point $x = (r \cos f, r \sin f)$ belongs to Λ , by \tilde{S}_{\pm} the set of all such points x) we get from (5) using $\phi' \leq 0$, $\phi \geq 1$, $\int_{\infty} K_2(u) df \ge 0 \text{ the estimate}$ (7) $\int_{\Omega_{m}}^{\infty} E(u) \, dx \leq (\phi(t))^{-1} \int_{\Omega_{m}}^{\infty} E(u) \, dx.$ c) Now we shall construct the function ϕ . Denoting $z = r \cdot u_{rr}$, $w = r^{-1} \cdot u$, $v = u_r$, we see, that ϕ will satisfy (6) if it satisfies

(8)
$$\int_{0}^{2\pi} (\mathcal{E}(w, v, z) \cdot \phi(r) - r^{2} \phi'' \mathcal{K}_{1}(w, v, z) - r \cdot \phi' \mathcal{K}_{2}(w, v, z)) df \geq 0$$

for all re(0, T) and all we $\tilde{H}^{2}(0, 2\tilde{k}), v \in \tilde{H}^{1}(0, 2\tilde{k}), z \in L_{2}(0, 2\tilde{k})$ where $\tilde{E}(w, v, z) = z^{2} + (w'' + v)^{2} + 2(v' - w')^{2}, \mathcal{K}_{1} = v^{2} + (w')^{2} - w \cdot z,$ $\begin{aligned} &\mathcal{H}_{2} = v^{2} + 2(w')^{2} - v \cdot w. \text{ (We use that } S_{r} \subset \{0, 2\} \text{ and the equality} \\ & E(u) = u_{rr}^{2} + r^{-2}((r^{-1} \cdot u_{ff} + u_{r})^{2} + 2(u_{rf} - r^{-1} \cdot u_{f})^{2}). \end{aligned}$ For fixed w, v, z with $\hat{\mathcal{E}}$ (w, v, z) = 1 (8) will be satisfied by the function

 $(\beta(r/T)^{-\alpha} + \alpha(r/T)^{\beta})/(\alpha + \beta)$ (9)

where

(9')
$$\propto = \frac{-(k_1 - k_2) + \left[(k_1 - k_2)^2 + 4 k_1\right]}{\frac{2 k_1}{k_1 - k_2}}$$

 $\beta = \frac{(k_1 - k_2) + \sqrt{(k_1 - k_2)^2 + 4 k_1}}{\frac{2 k_1}{k_1}}$

which solves the equation

(10) $r^2 \phi''(r)k_1 + r \phi'(r)k_2 = \phi, \phi(T) = 1, \phi'(T) = 0$ with $k_i = \int_0^{2T} \mathcal{H}_i(w, v, z) df$. (We may assume $k_1 > 0$.) The function (9) satisfies $\phi' \leq 0, \phi'' \geq 0, \phi''(r) + r^{-1}\phi'(r) \geq 0$ for $0 < r \leq T$. Using this it is easily seen that if ϕ satisfies (10) then it satisfies

(10')
$$r^2 \phi''(r) \widetilde{k}_1 + r \phi'(r) \widetilde{k}_2 \leq \phi(r)$$

with \widetilde{k}_1 , \widetilde{k}_2 satisfying one of the following pairs of conditions
 $\widetilde{k}_1 = k_1$, $\widetilde{k}_1 \leq k_1$, $\widetilde{k}_1 \leq k_1$,
 $\widetilde{k}_2 \geq k_2$, $\widetilde{k}_2 = k_2$, $\widetilde{k}_1 - \widetilde{k}_2 = k_1 - k_2$.

Now we prove the estimates

max $k_1(w, v, z) = \int_1 \le 3,2679$, max $(k_1(w, v, z) - k_2(w, v)) \le 2,29395$ where the max are taken on the set $\{w,v,z, \int_0^T \pounds(w,v,z) df = 1\}$. So the function (9) corresponding to $k_1 = 3,2679$ and $k_2 = 3,2679 -$ 2,29395 satisfies (8) for all w, v, z. Calculating the corresponding α from (9'), we obtain \geq 0,304. So the estimate (3) is proved. (4) follows from (3) e.g. by imbedding theorem.

<u>Remark 1</u>. From results of [2], [3] it follows that for

holds for arbitrary generalized solution considered in theorem proved above. It is an open problem if such estimate is valid for general Ω as in our theorem.

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<u>Remark 2</u>. One may derive other sufficient conditions for \oint to satisfy (6) e.g. (I) \oint (T) = 1, \oint '(T) = 0, \oint ' \leq 0, \oint '' \geq 0, \oint ''(r) + $r^{-1} \oint$ '(r) \geq 0, $0 < r \leq T$, 2($\ell(r)/V$)²(ϕ ''(r) + $r^{-1} \oint$ '(r)) + ($\ell(r)/2V$)² ϕ ''(r) + 4 r($\ell(r)/2rV$)³ $| \phi'| \leq \phi$ or (II) instead of inequality from (I) one assume the inequality (2 + 2 $\sqrt{2}$) ($\ell(r)/2V$)² ϕ ''(r) $\leq \phi$ (r), $0 < r \leq T$,

where $\ell(\mathbf{r})$ is the length of largest arc in $\widetilde{S}_{\mathbf{r}}$.

These conditions yield in the case of general domain not so good estimates as in our theorem, but in special cases of \mathcal{A} may give e.g. exponential decay of $\int_{\Omega_t} E(u) dx$ as $t \rightarrow 0$, since they exploit the form of Ω .

4. Further results

a) From above results one may derive uniqueness theorems for Dirichlet problem for $\Delta^2 u = f$ in unbounded domains. b) Analogously one may study the decay of energy or of the solution of (1) in unbounded domains as $\|x\| \to \infty$. One can prove e.g. that for the domain $\Omega^{\omega} = \{x = (x_1, x_2) = (r \cos \beta, r \sin \beta), \beta \in (0, \omega)\}$ the solution with finite energy tends to zero as $\|x\| \to \infty$ for ω sufficiently small. Since for $\omega = 2T$ this is not true, it is an open problem to find the largest ω for which this is true.

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