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ON A PARTIAL COMPLEX STRUCTURE

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Let $M^3 \subset \mathcal{C}^2$ be a real hypersurface. \mathcal{C}^2 being identified with \mathcal{R}^4 endowed with an endomorphism $J: \mathcal{R}^4 \rightarrow \mathcal{R}^4$ satisfying $J^2 = -\text{id}$., we get, on M^3 , a structure consisting of a field of tangent planes $\tau_m = T_m(M^3) \cap JT(M^3)$ and the restrictions $J_m: \tau_m \rightarrow \tau_m$ of J to τ_m ; see [3] and [4] respectively. Such a structure is called the partial complex structure. An attempt to solve the equivalence problem for the structures of this type has been made by E. CARTAN [1]; unfortunately, his treatment is not a very clear one. In what follows, I am going to present a more simple method for solving the mentioned problem.

Let M be a 3-dimensional manifold; in what follows, all the considered manifolds and maps are supposed to be of class C^∞ .

Definition. Let G be the group of matrices of the type

$$(1) \quad \begin{pmatrix} \alpha & -\beta & 0 \\ \beta & \alpha & 0 \\ \gamma & \delta & \varphi \end{pmatrix}; \quad \alpha, \beta, \gamma, \delta, \varphi \in \mathcal{R}; \quad (\alpha^2 + \beta^2)\varphi \neq 0.$$

A G -structure B_G on M is called a *partial complex structure*.

Let $(v_1, v_2, v_3), (w_1, w_2, w_3)$ be two sections of B_G over a domain $U \subset M$; we have

$$(2) \quad v_1 = \alpha w_1 - \beta w_2, \quad v_2 = \beta w_1 + \alpha w_2, \quad v_3 = \gamma w_1 + \delta w_2 + \varphi w_3.$$

At each point $m \in M$ there are induced a plane $\tau_m = \{v_1, v_2\}$ and an endomorphism $J_m: \tau_m \rightarrow \tau_m$ given by $J_m(v_1) = v_2, J_m(v_2) = -v_1$; obviously, $J_m^2 = -\text{id}$. We are going to suppose that *the field of planes τ is non-integrable*.

Definition. A vector field v on M is called a τ -field if $v_m \in \tau_m$ for each $m \in M$. A vector field u is called an *infinitesimal motion* of B_G if: (i) $\mathcal{L}_u v$ is a τ -field for each τ -field v , (ii) $J(\mathcal{L}_u v') = \mathcal{L}_u(Jv')$ for each τ -field v' ; here, $\mathcal{L}_u v = [u, v]$ is the Lie derivative of v . The structure B_G is said to be *locally transitive* if, for each $m \in M$ and $t \in T_m(M)$, there is a neighbourhood $U \subset M$ of m and an infinitesimal motion u over U such that $u_m = t$.

Our main problem is to determine all locally transitive G -structures. Let (v_1, v_2, v_3) , (w_1, w_2, w_3) be two sections of B_G over $U \subset M$. Then

$$(3) \quad \begin{aligned} [v_1, v_2] &= a_1 v_1 + a_2 v_2 + a_3 v_3, & [w_1, w_2] &= A_1 w_1 + A_2 w_2 + A_3 w_3, \\ [v_1, v_3] &= b_1 v_1 + b_2 v_2 + b_3 v_3, & [w_1, w_3] &= B_1 w_1 + B_2 w_2 + B_3 w_3, \\ [v_2, v_3] &= c_1 v_1 + c_2 v_2 + c_3 v_3, & [w_2, w_3] &= C_1 w_1 + C_2 w_2 + C_3 w_3, \end{aligned}$$

the functions $a_1, \dots, c_3, A_1, \dots, C_3$ satisfying the Jacobi identities

$$(4) \quad \begin{aligned} [v_1, [v_2, v_3]] + [v_2, [v_3, v_1]] + [v_3, [v_1, v_2]] &= 0, \\ [w_1, [w_2, w_3]] + [w_2, [w_3, w_1]] + [w_3, [w_1, w_2]] &= 0. \end{aligned}$$

We get

$$\begin{aligned} [v_1, v_2] &= \{v_1 \beta - v_2 \alpha + (\alpha^2 + \beta^2) A_1\} w_1 + \{v_1 \alpha + v_2 \beta + (\alpha^2 + \beta^2) A_2\} w_2 + \\ &\quad + (\alpha^2 + \beta^2) A_3 w_3 = \\ &= (\alpha a_1 + \beta a_2 + \gamma a_3) w_1 + (\alpha a_2 - \beta a_1 + \delta a_3) w_2 + \varphi a_3 w_3, \\ [v_1, v_3] &= (\cdot) w_1 + (\cdot) w_2 + \{v_1 \varphi + (\alpha \delta + \beta \gamma) A_3 + \alpha \varphi B_3 - \beta \varphi C_3\} w_3 = \\ &= (\cdot) w_1 + (\cdot) w_2 + \varphi b_3 w_3, \\ [v_2, v_3] &= (\cdot) w_1 + (\cdot) w_2 + \{v_2 \varphi + (\beta \delta - \alpha \gamma) A_3 + \beta \varphi B_3 + \alpha \varphi C_3\} w_3 = \\ &= (\cdot) w_1 + (\cdot) w_2 + \varphi c_3 w_3. \end{aligned}$$

From this, we get the existence of sections satisfying $a_3 = 1$. Suppose $a_3 = A_3 = 1$, i.e., $\varphi = \alpha^2 + \beta^2$. Let us look for the existence of a section (w_1, w_2, w_3) satisfying $A_1 = A_2 = B_3 = C_3 = 0$. This amounts to the existence of $\alpha, \beta, \gamma, \delta$ such that

$$\begin{aligned} v_1 \beta - v_2 \alpha &= a_1 \alpha + a_2 \beta + \gamma, & 2\alpha v_1 \alpha + 2\beta v_1 \beta + \alpha \delta + \beta \gamma &= (\alpha^2 + \beta^2) b_3, \\ v_1 \alpha + v_2 \beta &= a_2 \alpha - a_1 \beta + \delta, & 2\alpha v_2 \alpha + 2\beta v_2 \beta + \beta \delta - \alpha \gamma &= (\alpha^2 + \beta^2) c_3. \end{aligned}$$

It is easy to see that this system has (at least locally) solutions such that $\alpha^2 + \beta^2 \neq 0$. Let $a_1 = a_2 = b_3 = c_3 = 0, a_3 = 1$ for (v_1, v_2, v_3) . From the Jacobi identity (4), we get $v_1 c_1 - v_2 b_1 = v_1 c_2 - v_2 b_2 = c_2 + b_1 = 0$. A τ -field v is called special if the section $(v_1 = v, v_2 = Jv, v_3 = [v, Jv])$ has the just described property.

Theorem. Let B_G be a partial complex structure over M . Let v, w be its special τ -fields, and let

$$(5) \quad \begin{aligned} [v, [v, Jv]] &= av + bJv, & [Jv, [v, Jv]] &= cv - aJv, \\ (Jv)b + va &= 0, & (Jv)a - vc &= 0; \end{aligned}$$

$$(6) \quad [w, [w, Jw]] = Aw + BJw, \quad [Jw, [w, Jw]] = Cw - AJw, \\ (Jw)B + wA = 0, \quad (Jw)A - wC = 0;$$

$$(7) \quad v = \alpha w - \beta Jw.$$

Consider the functions

$$(8) \quad j_1 = (v - Jv \cdot Jv)(c - b) + 8[v, Jv]a - 3(c^2 - b^2), \\ j_2 = (v \cdot Jv + Jv \cdot v)(c - b) + 4[v, Jv](b + c) + 6a(c - b);$$

J_1 and J_2 be defined similarly. Then:

(i) We have

$$(9) \quad j_1^2 + j_2^2 = (\alpha^2 + \beta^2)^4 (J_1^2 + J_2^2).$$

(ii) If $j_1 = j_2 = 0$, B_G is locally transitive. For each point $m \in M$, there is its neighbourhood $U \subset M$ and special τ -fields v over U satisfying

$$(10) \quad [v, [v, Jv]] = 0, \quad [Jv, [v, Jv]] = 0.$$

Choose such a τ -field v . Further, choose arbitrary real numbers R_1, \dots, R_8 . Then there is exactly one field $u \in \mathcal{L}(B_G)$ over a suitable neighbourhood $m \in U_1 \subset U$ such that

$$(11) \quad u_m = R_1 v_m + R_2 (Jv)_m + R_3 [v, Jv]_m, \\ [v, u]_m = R_4 v_m + R_5 (Jv)_m, \\ [v, [v, u]]_m = R_6 v_m + R_7 (Jv)_m + R_8 [v, Jv]_m, \\ [v, [v, [v, u]]]_m = R_8 (Jv)_m + 2R_7 [v, Jv]_m.$$

(iii) Let $j_1^2 + j_2^2 \neq 0$. To each point $m \in M$, there is its neighbourhood U and exactly two special τ -fields $v, v' = -v$ over U satisfying (5) and

$$(12) \quad j_1 = 1, \quad j_2 = 0.$$

B_G being transitive, v satisfies

$$(13) \quad [v, [v, Jv]] = bJv, \quad [Jv, [v, Jv]] = cv; \\ b, c \in \mathcal{R}, \quad 3(c^2 - b^2) + 1 = 0.$$

For each vector $t \in T_m(M)$, there is exactly one field $u \in \mathcal{L}(B_G)$ — defined over a suitable neighbourhood U_1 of m — such that $u_m = t$.

Proof. Let $T^c(M) = T(M) \oplus i T(M)$ be the complexification of the tangent bundle $T(M)$ of M . The bracket operation in $T^c(M)$ be introduced, quite naturally, by

$$(14) \quad [v + iv', w + iw'] = [v, w] - [v', w'] + i([v', w] + [v, w']).$$

In B_G , consider two special sections $(v_1, v_2, v_3), (w_1, w_2, w_3)$ satisfying

$$(15) \quad [v_1, v_2] = v_3, \quad [v_1, v_3] = av_1 + bv_2, \quad [v_2, v_3] = cv_1 - av_2;$$

$$[w_1, w_2] = w_3, \quad [w_1, w_3] = Aw_1 + Bw_2, \quad [w_2, w_3] = Cw_1 - Aw_2.$$

In $T^c(M)$, consider the vector fields

$$(16) \quad V_1 = v_1 + iv_2, \quad V_2 = v_1 - iv_2, \quad V_3 = -2iv_3;$$

$$W_1 = w_1 + iw_2, \quad W_2 = w_1 - iw_2, \quad W_3 = -2iw_3.$$

We get

$$(17) \quad [V_1, V_2] = V_3, \quad [V_1, V_3] = pV_1 + qV_2, \quad [V_2, V_3] = rV_1 - pV_2,$$

$$p = c - b, \quad q = b + c - 2ia, \quad r = -(b + c + 2ia),$$

$$V_2q = -V_1p, \quad V_1r = V_2p;$$

$$[W_1, W_2] = W_3, \quad [W_1, W_3] = PW_1 + QW_2, \quad [W_2, W_3] = RW_1 - PW_2,$$

$$P = C - B, \quad Q = B + C - 2iA, \quad R = -(B + C + 2iA),$$

$$W_2Q = -W_1P, \quad W_1R = W_2P.$$

From (2),

$$(18) \quad V_1 = \varrho W_1, \quad V_2 = \sigma W_2, \quad V_3 = \mu W_1 + \nu W_2 + \varphi W_3;$$

$$\varrho = \alpha + i\beta, \quad \sigma = \alpha - i\beta, \quad \mu = -(\delta + i\gamma), \quad \nu = \delta - i\gamma; \quad \varrho\sigma\varphi \neq 0.$$

Now,

$$(19) \quad \varphi = \varrho\sigma,$$

$$(20) \quad V_1\varrho = -2\varrho\sigma^{-1}\nu, \quad V_2\varrho = -\mu; \quad V_1\sigma = \nu, \quad V_2\sigma = 2\varrho^{-1}\sigma\mu;$$

$$V_2\mu = \varrho r - \varrho\sigma^2 R, \quad V_1\nu = \sigma q - \varrho^2\sigma Q,$$

$$V_1\mu - V_3\varrho = \varrho p - \varrho^2\sigma P, \quad V_2\nu - V_3\sigma = -\sigma p + \varrho\sigma^2 P.$$

It is known [2] that the integrability conditions of (20) imply

$$(21) \quad k_1 = \varrho^3\sigma K_1, \quad k_2 = \varrho\sigma^3 K_2$$

with

$$(22) \quad k_1 = V_1 V_1 p - 2V_3 q - 3pq, \quad k_2 = V_2 V_2 p - 2V_3 r + 3pr$$

and similar definition of K_1 and K_2 respectively. We get

$$(23) \quad k_1 = j_1 + ij_2, \quad k_2 = j_1 - ij_2,$$

j_1 and j_2 being given by (8). Thus (21) reduce to $k_1 = \varrho^3 \bar{\varrho} K_1$; the equation $k_1 \bar{k}_1 = \varrho^4 \bar{\varrho}^4 K_1 \bar{K}_1$ is exactly (9). The equation $k_1 = \varrho^3 \bar{\varrho}$ has solutions $\varrho \neq 0$ for each $k_1 \neq 0$; the equation $1 = \varrho^3 \bar{\varrho}$ has exactly two solutions $\varrho = \pm 1$. Thus (i) and the first part of (iii) have been proved.

Suppose $j_1 = j_2 = 0$, i.e., $k_1 = k_2 = 0$, and consider the system (19) + (20), $P = Q = R = 0$. According to [2], this system is completely integrable, and we have proved the first part of (ii).

Now, consider the structure B_G , given by a section (v_1, v_2, v_3) satisfying

$$(24) \quad [v_1, v_2] = v_3, \quad [v_1, v_3] = 0, \quad [v_2, v_3] = 0.$$

Let $u \in \mathcal{L}(B_G)$,

$$(25) \quad u = Av_1 + Bv_2 + Cv_3.$$

From

$$\begin{aligned} [v_1, u] &= v_1 A \cdot v_1 + v_1 B \cdot v_2 + (v_1 C + B) v_3, \\ [v_2, u] &= v_2 A \cdot v_1 + v_2 B \cdot v_2 + (v_2 C - A) v_3, \end{aligned}$$

we get

$$(26) \quad v_1 A - v_2 B = v_2 A + v_1 B = v_1 C + B = v_2 C - A = 0.$$

For $D := v_1 A$, $E := v_2 A$, we have

$$(27) \quad \begin{aligned} v_1 A &= D, \quad v_2 A = E; \\ v_1 B &= -E, \quad v_2 B = D; \quad v_1 C = -B, \quad v_2 C = A. \end{aligned}$$

From (27_{5,6}),

$$(28) \quad v_3 C = 2D;$$

the integrability conditions of (27₁) + (27₂), (27₃) + (27₄), (27₅) + (28) and (27₆) + (28) are

$$\begin{aligned} v_3 A &= v_1 E - v_2 D, \quad v_3 B = v_1 D + v_2 E, \\ 0 &= 2v_1 D + v_3 B, \quad 0 = 2v_2 D - v_3 A \end{aligned}$$

respectively; for $3F := v_1E$, $3G := v_2E$, they become

$$(29) \quad \begin{aligned} v_3A &= 2F; \quad v_3B = 2G; \\ v_1D &= -G, \quad v_2D = F; \quad v_1E = 3F, \quad v_2E = 3G. \end{aligned}$$

The integrability conditions of $(27_1) + (29_1)$, $(27_2) + (29_1)$, $(27_3) + (29_2)$, $(27_4) + (29_2)$, $(29_3) + (29_4)$ and $(29_5) + (29_6)$ being

$$\begin{aligned} 2v_1F &= v_3D, \quad 2v_2F = v_3E, \quad 2v_1G = -v_3E, \quad 2v_2G = v_3D, \\ v_3D &= v_1F + v_2G, \quad v_3E = 3v_1G - 3v_2F, \end{aligned}$$

we have

$$(30) \quad v_3D = 2H; \quad v_3E = 0; \quad v_1F = H, \quad v_2F = 0; \quad v_1G = 0, \quad v_2G = H$$

for $H := v_1F$. The integrability conditions of $(29_3) + (30_1)$, $(29_4) + (30_1)$, $(29_5) + (30_2)$, $(29_6) + (30_2)$, $(30_3) + (30_4)$ and $(30_5) + (30_6)$ are

$$\begin{aligned} 2v_1H &= -v_3G, \quad 2v_2H = v_3F, \quad v_3F = 0, \quad v_3G = 0, \\ v_3F &= -v_2H, \quad v_3G = v_1H; \end{aligned}$$

from these, we get

$$(31) \quad v_3F = 0; \quad v_3G = 0; \quad v_1H = 0, \quad v_2H = 0.$$

The integrability conditions of $(30_3) + (31_1)$, $(30_4) + (31_1)$, $(30_5) + (31_2)$, $(30_6) + (31_2)$ and $(31_3) + (31_4)$ reduce to

$$(32) \quad v_3H = 0.$$

Thus the system (27)–(32) is completely integrable. For $u \in \mathcal{L}(B_G)$ given by (25), we get

$$(33) \quad \begin{aligned} [v_1, u] &= Dv_1 - Ev_2, \quad [v_1, [v_1, u]] = -Gv_1 - 3Fv_2 - Ev_3, \\ [v_1, [v_1, [v_1, u]]] &= -3Hv_2 - 6Fv_3, \end{aligned}$$

this completing the proof of (ii).

Finally, let B_G be transitive with $j_1^2 + j_2^2 \neq 0$. Then they are exactly two special sections of B_G satisfying $j_1 = 1, j_2 = 0$; let (v_1, v_2, v_3) be one of them. The functions a, b, c being now invariants of B_G , they have to be constants, and we get $-3(c^2 - b^2) = 1, a(c - b) = 0$. Because of $c - b \neq 0, a = 0$, i.e.,

$$(34) \quad \begin{aligned} [v_1, v_2] &= v_3, \quad [v_1, v_3] = bv_2, \quad [v_2, v_3] = cv_1; \\ b, c &\in \mathcal{R}, \quad 3(c^2 - b^2) + 1 = 0. \end{aligned}$$

Let $u \in \mathcal{L}(B_G)$, u being given by (25). From

$$\begin{aligned} [v_1, u] &= v_1A \cdot v_1 + (v_1B + bC)v_2 + (v_1C + B)v_3, \\ [v_2, u] &= (v_2A + cC)v_1 + v_2B \cdot v_2 + (v_2C - A)v_3, \end{aligned}$$

we get

$$v_1A - v_2B = v_2A + v_1B + (b + c)C = v_1C + B = v_2C - A = 0.$$

For $D := v_1A$, $E := v_1B + bC$,

$$(35) \quad \begin{aligned} v_1A = D, \quad v_2A = -E - cC; \quad v_1B = E - bC, \quad v_2B = D; \\ v_1C = -B, \quad v_2C = A. \end{aligned}$$

The integrability condition of (35₅) and (35₆) is

$$(36) \quad v_3C = 2D.$$

The integrability conditions of (35₁) + (35₂), (35₃) + (35₄), (35₅) + (36) and (35₆) + (36) are

$$\begin{aligned} v_3A + v_1E + v_2D = cB, \quad v_3B - v_1D + v_2E = bA, \\ 2v_1D + v_3B = bA, \quad v_3A - 2v_2D = cB. \end{aligned}$$

For $3F := v_1E$, $3G := v_2A$, we get

$$(37) \quad \begin{aligned} v_3A = -2F + cB; \quad v_3B = -2G + bA; \\ v_1D = G, \quad v_2D = -F; \quad v_1E = 3F, \quad v_2E = 3G. \end{aligned}$$

The integrability conditions of (35₁) + (37₁), (35₂) + (37₁), (35₃) + (37₂), (35₄) + (37₂), (37₃) + (37₄) and (37₅) + (37₆) are

$$\begin{aligned} 2v_1F + v_3D = (b + c)E, \quad 2v_2F - v_3E = 2cD, \\ 2v_1G + v_3E = 2bD, \quad 2v_2G + v_3D = -(b + c)E, \\ v_3D + v_1F + v_2G = 0, \quad v_3E - 3v_1G + 3v_2F = 0. \end{aligned}$$

For $H := v_1F - bE$, they are

$$(38) \quad \begin{aligned} v_3D = -2H + (c - b)E; \quad v_3E = 2(b - c)D; \\ v_1F = H + bE, \quad v_2F = \frac{1}{3}(3b + 5c)D; \\ v_1G = \frac{1}{3}(5b + 3c)D, \quad v_2G = H - cE. \end{aligned}$$

The integrability conditions of $(37_3) + (38_1)$, $(37_4) + (38_1)$, $(37_5) + (38_2)$, $(37_6) + (38_2)$, $(38_3) + (38_4)$ and $(38_5) + (38_6)$ are

$$(39) \quad 2v_1H + v_3G = (3c - 2b)F, \quad 2v_2H - v_3F = (2c - 3b)G;$$

$$(40) \quad v_3F = -\frac{1}{3}(b + 2c)G; \quad v_3G = -\frac{1}{3}(2b + c)F;$$

$$(41) \quad v_2H = \frac{1}{3}(7c - 5b)G, \quad v_1H = \frac{1}{3}(5c - 7b)F.$$

Substituting from (40) and (41) into (39), we get

$$(42) \quad bF = 0, \quad cG = 0.$$

Let $b \neq 0 \neq c$, i.e., $F = G = 0$. From $(38_{4,5})$, $D = 0$. The equations $(38_{3,6})$ imply $(b + c)E = 0$; $b + c = 0$ being impossible because of (34), we have $E = 0$. We get $H = 0$ from (38_3) . Thus we obtain the completely integrable system

$$(43) \quad \begin{aligned} v_1A &= 0, & v_2A &= -cC, & v_3A &= cB, \\ v_1B &= -bC, & v_2B &= 0, & v_3B &= bA, \\ v_1C &= -B, & v_2C &= A, & v_3C &= 0. \end{aligned}$$

Next, suppose $b \neq 0, c = 0$ (the case $b = 0, c \neq 0$ being analogous). Then $F = 0$ because of (42_1) . We get $D = 0$ and $G = 0$ from (38_4) and (40_1) respectively. From $(38_{1,3})$, $v_3D + v_1F = -H$, $v_3D + 2v_1F = bE$, i.e., $E = H = 0$. Thus we obtain the system (43) with $c = 0$. This proves the second part of (iii) and the Theorem.

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