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EQUIVALENCE PROBLEM FOR LAGRANGIANS

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The equivalence problem for Lagrangians was solved by E. Cartan by his own methods. Recently, a solution in the same spirit was presented by R. B. Gardner, see [1]. Here, I present another approach to the same problem which seems to me to be simpler and more effective. I restrict myself just to the cases of order one and two; in order two, I get as the special case Lagrangians of the form (2.17).

I am grateful to Prof. I. Kolář for several interesting discussions.

1. FIRST ORDER LAGRANGIANS

Let a fibred manifold $\tilde{\pi}: \tilde{M} \rightarrow \tilde{N}$ be given such that $\dim \tilde{N} = 1, \dim \tilde{M} = 2$; let $J^1(\tilde{M})$ be its first jet prolongation and $\tilde{m} \in J^1(\tilde{M})$. On $J^1(\tilde{M})$, let us choose local fibre coordinates (ξ, η, ζ) such that $\tilde{\pi}_{1,0}(\xi, \eta, \zeta) = (\xi, \eta), \tilde{\pi}(\xi, \eta) = \zeta$ and $\tilde{m} = (0, 0, 0)$.

Further, let a fibred manifold $\pi: M \rightarrow N$ be given such that $\dim N = 1, \dim M = 2$; let $J^1(M)$ be its first jet prolongation. On $J^1(M)$, let us introduce fibre coordinates $(x, y, z) \equiv (x, y, \dot{y})$ in a similar way. Let $f_0: \tilde{M} \rightarrow M$ be a (local) bundle isomorphism, $f \equiv j^1(f_0): J^1(\tilde{M}) \rightarrow J^1(M)$ its prolongation, and let $m \equiv f(\tilde{m}) = (x_0, y_0, z_0) \in J^1(M)$. The (local) isomorphism f is given by

$$(1.1) \quad x = x(\xi), \quad y = y(\xi, \eta), \quad z = \left(\frac{dx}{d\xi}\right)^{-1} \left(\frac{\partial y}{\partial \xi} + \frac{\partial y}{\partial \eta} \zeta\right),$$

and we have

$$(1.2) \quad x_0 = x(0), \quad y_0 = y(0, 0), \quad z_0 = \left(\frac{dx(0)}{d\xi}\right)^{-1} \frac{\partial y(0, 0)}{\partial \xi}.$$

It is easy to check that

$$(1.3) \quad df \left(\frac{\partial}{\partial \xi} \Big|_{\tilde{m}} \right) = \left(\frac{dx(0)}{d\xi} \right)^{-1} \frac{\partial y(0, 0)}{\partial \eta} \frac{\partial}{\partial z} \Big|_m,$$

$$df \left(\frac{\partial}{\partial \eta} \Big|_{\tilde{m}} \right) = \frac{\partial y(0, 0)}{\partial \eta} \frac{\partial}{\partial y} \Big|_m + \left(\frac{dx(0)}{d\xi} \right)^{-1} \frac{\partial^2 y(0, 0)}{\partial \xi \partial \eta} \frac{\partial}{\partial z} \Big|_m,$$

$$df\left(\frac{\partial}{\partial \xi}\Big|_{\tilde{m}}\right) = \frac{dx(0)}{d\xi} \frac{\partial}{\partial x}\Big|_m + \frac{\partial y(0,0)}{\partial \xi} \frac{\partial}{\partial y}\Big|_m + \left\{ -\left(\frac{dx(0)}{d\xi}\right)^{-2} \frac{d^2x(0)}{d\xi^2} \frac{\partial y(0,0)}{\partial \xi} + \left(\frac{dx(0)}{d\xi}\right)^{-1} \frac{\partial^2 y(0,0)}{\partial \xi^2} \right\} \frac{\partial}{\partial z}\Big|_m.$$

Let us write

$$(1.4) \quad D := \frac{dx(0)}{d\xi}, \quad C := \frac{\partial y(0,0)}{\partial \eta}, \quad A := D^{-1} \frac{\partial^2 y(0,0)}{\partial \xi \partial \eta}, \\ B := -D^{-2} \frac{d^2x(0)}{d\xi^2} \frac{\partial y(0,0)}{\partial \xi} + D^{-1} \frac{\partial^2 y(0,0)}{\partial \xi^2}.$$

Notice that

$$(1.5) \quad DC \neq 0,$$

f being a (local) isomorphism; from (1.2₃),

$$(1.6) \quad \frac{\partial y(0,0)}{\partial \xi} = z_0 D.$$

Thus the differential $df: T_{\tilde{m}}(J^1(\tilde{M})) \rightarrow T_m(J^1(M))$ maps the vectors $\partial/\partial \xi, \partial/\partial \eta, \partial/\partial \xi$ at \tilde{m} to the vectors

$$(1.7) \quad v_1 = CD^{-1} \frac{\partial}{\partial z}, \quad v_2 = C \frac{\partial}{\partial y} + A \frac{\partial}{\partial z}, \quad v_3 = D \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} \right) + B \frac{\partial}{\partial z}$$

at m . Each triple of vectors $\{v_1, v_2, v_3\}$, $v_i \in T_m(J^1(M))$, of the form (1.7) is called a *frame* of $J^1(M)$ at m .

On $J^1(M)$, let a Lagrangian

$$(1.8) \quad \lambda = f(x, y, z) dx, \quad ff_z \neq 0$$

be given; here, $f_z = \partial f/\partial z$ etc. A frame $\{v_1, v_2, v_3\}$ of $J^1(M)$ at m will be called a *λ -frame* if

$$(1.9) \quad \lambda(v_3) = 1$$

at m ; from (1.7₃), we get $f(x_0, y_0, z_0) D = 1$. Now, let us consider, on $J^1(M)$, a field of λ -frames

$$(1.10) \quad v_1 = Cf \frac{\partial}{\partial z}, \quad v_2 = C \frac{\partial}{\partial y} + A \frac{\partial}{\partial z}, \quad v_3 = f^{-1} \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} \right) + B \frac{\partial}{\partial z},$$

A, B and C being functions on $J^1(M)$. Let functions $c_{jk}^i: J^1(M) \rightarrow \mathbb{R}$ be defined by

$$(1.11) \quad [v_i, v_j] = \sum_{k=1}^3 c_{ij}^k v_k; \quad i, j = 1, 2, 3;$$

$[v_i, v_j]$ being the Lie bracket of the vector fields v_i, v_j . Let us try to exhibit special fields of λ -frames by imposing suitable conditions on the functions c_{ij}^k .

We get $[v_1, v_3] = c_{13}^1 v_1 + v_2 - f_z C v_3$, and the condition

$$(1.12) \quad c_{13}^3 = -1$$

implies

$$(1.13) \quad C = f_z^{-1}.$$

Further,

$$[v_2, v_3] = c_{23}^1 v_1 + \{f^{-1} f_z A + f^{-1} f_z^{-1} (f_{xz} + z f_{yz}) + f_z^{-1} f_{zz} B\} v_2 - f^{-1} (f_y f_z^{-1} + f_z A) v_3.$$

The condition

$$(1.14) \quad c_{23}^3 = 0$$

yields

$$(1.15) \quad A = -f_y f_z^{-2}$$

and

$$(1.16) \quad c_{23}^2 = f^{-1} f_z^{-1} (f_{xz} + z f_{yz} - f_y) + f_z^{-1} f_{zz} B.$$

Proposition 1.1. *Let the Lagrangian λ (1.8) satisfy $f_{zz} \neq 0$. Then there is, on $J^1(M)$, exactly one field of λ -frames $\{v_1, v_2, v_3\}$ such that we have, in (1.11), $c_{13}^3 = -1$, $c_{23}^3 = c_{23}^2 = 0$ (this implying $c_{13}^2 = 1$). This field is given by*

$$(1.17) \quad v_1 = f f_z^{-1} \frac{\partial}{\partial z}, \quad v_2 = f_z^{-1} \frac{\partial}{\partial y} - f_y f_z^{-2} \frac{\partial}{\partial z},$$

$$v_3 = f^{-1} \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} \right) - f^{-1} f_{zz}^{-1} (f_{xz} + z f_{yz} - f_y) \frac{\partial}{\partial z}.$$

Now, let $f_{zz} = 0$, i.e.,

$$(1.18) \quad f(x, y, z) = a(x, y) + b(x, y) z.$$

Let us suppose (1.12) and (1.14), i.e., (1.13) and (1.15). Then

$$(1.19) \quad [v_2, v_3] = c_{23}^1 v_1 + F v_2 \quad \text{with}$$

$$F = f^{-1} f_z^{-1} (f_{xz} + z f_{yz} - f_y) = f^{-1} b^{-1} (b_x - a_y).$$

Further, $[v_1, v_3] = c_{13}^1 v_1 + v_2 - v_3$, i.e.,

$$(1.20) \quad [v_3, [v_2, v_3]] = c_{323}^1 v_1 + (v_3 F - F^2 - c_{23}^1) v_2 + c_{23}^1 v_3,$$

the functions $c_{jkl}^i: J^1(M) \rightarrow \mathbb{R}$ being given by

$$(1.21) \quad [v_j, [v_k, v_l]] = \sum_{i=1}^3 c_{jkl}^i v_i.$$

The condition

$$(1.22) \quad c_{23}^1 + c_{323}^2 = 0$$

reads

$$v_3 F - F^2 = f^{-1} (F_x + z F_y) + B F_z - F^2 = 0,$$

i.e.,

$$(1.23) \quad (b_x - a_y) B = f \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} \right) (f^{-1} b^{-1} (b_x - a_y)) - f^{-1} b^{-2} (b_x - a_y)^2.$$

Proposition 1.2. *Let the Lagrangian λ be of the form (1.18) with $b_x \neq a_y$. Then*

there is, on $J^1(M)$, exactly one field of λ -frames $\{v_1, v_2, v_3\}$ such that we have $c_{13}^3 = -1$, $c_{23}^3 = c_{23}^1 + c_{323}^2 = 0$; these λ -frames are given by (1.10) with (1.13), (1.15) and (1.23).

Finally, let us suppose (1.18) with $b_x = a_y$. Then there is a function c such that we may write

$$(1.24) \quad f(x, y, z) = c_x + c_y z, \quad c = c(x, y).$$

On $J^1(M)$, introduce new fibre coordinates (x, Y, Z) with $Y = c(x, y)$. Then $Z = c_x + c_y z$, and we have the following

Proposition 1.3. *Let the Lagrangian λ be of the form (1.18) with $b_x = a_y$. Then there are, on $J^1(M)$, new bundle coordinates (x, Y, Z) such that $\lambda = Z dx$.*

Let us remark that the coframe dual to the frame

$$(1.25) \quad v_1 = f f_z^{-1} \frac{\partial}{\partial z}, \quad v_2 = f_z^{-1} \frac{\partial}{\partial y} - f_y f_z^{-2} \frac{\partial}{\partial z}, \quad v_3 = f^{-1} \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} \right) + B \frac{\partial}{\partial z}$$

is given by the 1-forms

$$(1.26) \quad \omega^1 = -f^{-1} f_z B dx + f^{-1} f_y (dy - z dx) + f^{-1} f_z dz, \\ \omega^2 = f_z (dy - z dx), \quad \omega^3 = \lambda = f dx.$$

On $J^1(M)$, let a Lagrangian λ (1.8) with $f_{zz} \neq 0$ be given. The corresponding Euler equation is

$$(1.27) \quad E(\lambda) \equiv f_y - \frac{d}{dx} f_z = f_y - f_{xz} - f_{yz} y' - f_{zz} y'' = 0.$$

Let γ be a critical section of λ given by $y = y(x)$. Then $j^1 \gamma$ is given by $y = y(x)$, $z = y'(x)$, and its tangent vector at each point is

$$(1.28) \quad \frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + f_{zz}^{-1} (f_y - f_{xz} - f_{yz} z) \frac{\partial}{\partial z} = f v_3,$$

v_3 being exactly the vector (1.17₃).

2. SECOND ORDER LAGRANGIANS

Let $\tilde{\pi}: \tilde{M} \rightarrow \tilde{N}$ be a fibred manifold with $\dim \tilde{N} = 1$, $\dim \tilde{M} = 2$. On $J^2(\tilde{M})$, let a point \tilde{m} and fibre coordinates (ξ, η, ζ, τ) be given such that $\tilde{m} = (0, 0, 0, 0)$. Analogously, let $\pi: M \rightarrow N$ be a fibred manifold with fibre coordinates $(x, y, z, t) \equiv (x, y, \dot{y}, \ddot{y})$ on $J^2(M)$. The prolongation $f := j^2(f_0)$ of a (local) bundle isomorphism $f_0: \tilde{M} \rightarrow M$ is given by (1.1) and

$$(2.1) \quad t = - \left(\frac{dx}{d\xi} \right)^{-3} \frac{d^2 x}{d\xi^2} \left(\frac{\partial y}{\partial \xi} + \frac{\partial y}{\partial \eta} \zeta \right) + \left(\frac{dx}{d\xi} \right)^{-2} \left(\frac{\partial^2 y}{\partial \xi^2} + 2 \frac{\partial^2 y}{\partial \xi \partial \eta} \zeta + \frac{\partial^2 y}{\partial \eta^2} \zeta^2 + \frac{\partial y}{\partial \eta} \tau \right).$$

A set of vectors $v_1, v_2, v_3, v_4 \in T_m(J^2(M))$ is called a frame of $J^2(M)$ at $m =$

$= (x, y, z, t)$ if they are images of the vectors $\partial/\partial\tau, \partial/\partial\zeta, \partial/\partial\eta, \partial/\partial\xi \in T_{\tilde{m}}(J^2(\tilde{M}))$ under the mapping $df: T_{\tilde{m}}(J^2(\tilde{M})) \rightarrow T_m(J^2(M))$, $f_0: \tilde{M} \rightarrow M$ being an arbitrary (local) bundle isomorphism such that $f(\tilde{m}) = m$. It is easy to check

$$(2.2) \quad \begin{aligned} v_1 &= \left(\frac{dx}{d\xi}\right)^{-2} \frac{\partial y}{\partial \eta} \frac{\partial}{\partial t}, \\ v_2 &= \left(\frac{dx}{d\xi}\right)^{-1} \frac{\partial y}{\partial \eta} \frac{\partial}{\partial z} + \left\{ -\left(\frac{dx}{d\xi}\right)^{-3} \frac{d^2x}{d\xi^2} \frac{\partial y}{\partial \eta} + 2\left(\frac{dx}{d\xi}\right)^{-2} \frac{\partial^2 y}{\partial \xi \partial \eta} \right\} \frac{\partial}{\partial t}, \\ v_3 &= \frac{\partial y}{\partial \eta} \frac{\partial}{\partial y} + \left(\frac{dx}{d\xi}\right)^{-1} \frac{\partial^2 y}{\partial \xi \partial \eta} \frac{\partial}{\partial z} + \left\{ -\left(\frac{dx}{d\xi}\right)^{-3} \frac{d^2x}{d\xi^2} \frac{\partial^2 y}{\partial \xi \partial \eta} + \left(\frac{dx}{d\xi}\right)^{-2} \frac{\partial^3 y}{\partial \xi^2 \partial \eta} \right\} \frac{\partial}{\partial t}, \\ v_4 &= \frac{dx}{d\xi} \frac{\partial}{\partial x} + \frac{\partial y}{\partial \xi} \frac{\partial}{\partial y} + \left\{ -\left(\frac{dx}{d\xi}\right)^{-2} \frac{d^2x}{d\xi^2} \frac{\partial y}{\partial \xi} + \left(\frac{dx}{d\xi}\right)^{-1} \frac{\partial^2 y}{\partial \xi^2} \right\} \frac{\partial}{\partial z} \\ &\quad + \left\{ 3\left(\frac{dx}{d\xi}\right)^{-4} \left(\frac{d^2x}{d\xi^2}\right)^2 \frac{\partial y}{\partial \xi} - \left(\frac{dx}{d\xi}\right)^{-3} \frac{d^3x}{d\xi^3} \frac{\partial y}{\partial \xi} - \left(\frac{dx}{d\xi}\right)^{-3} \frac{d^2x}{d\xi^2} \frac{\partial^2 y}{\partial \xi^2} - \right. \\ &\quad \left. - 2\left(\frac{dx}{d\xi}\right)^{-3} \frac{d^2x}{d\xi^2} \frac{\partial^2 y}{\partial \xi^2} + \left(\frac{dx}{d\xi}\right)^{-2} \frac{\partial^3 y}{\partial \xi^3} \right\} \frac{\partial}{\partial t}, \end{aligned}$$

all the derivatives being calculated at $\tilde{m} = (0, 0, 0, 0)$. From (1.1) and (2.1), we obtain

$$(2.3) \quad z = \left(\frac{dx}{d\xi}\right)^{-1} \frac{\partial y}{\partial \xi}, \quad t = -\left(\frac{dx}{d\xi}\right)^{-3} \frac{d^2x}{d\xi^2} \frac{\partial y}{\partial \xi} + \left(\frac{dx}{d\xi}\right)^{-2} \frac{\partial^2 y}{\partial \xi^2},$$

again at \tilde{m} . From (2.3), we calculate $\partial y(0, 0)/\partial \xi, \partial^2 y(0, 0)/\partial \xi^2$ and substitute them into (2.2). Thus we see that the most general frame at $m \in J^2(M)$ is given by

$$(2.4) \quad \begin{aligned} v_1 &= F^{-2}A \frac{\partial}{\partial t}, \quad v_2 = F^{-1}A \frac{\partial}{\partial z} + B \frac{\partial}{\partial t}, \quad v_3 = A \frac{\partial}{\partial y} + C \frac{\partial}{\partial z} + D \frac{\partial}{\partial t}, \\ v_4 &= F \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + t \frac{\partial}{\partial z} \right) + E \frac{\partial}{\partial t}, \end{aligned}$$

A, \dots, F being arbitrary real numbers.

On $J^2(M)$, let a Lagrangian

$$(2.5) \quad \lambda = f(x, y, z, t) dx, \quad ff_t \neq 0$$

be given. A frame $\{v_1, v_2, v_3, v_4\}$ is called a λ -frame if $\lambda(v_4) = 1$. The most general λ -frame is

$$(2.6) \quad \begin{aligned} v_1 &= f^2A \frac{\partial}{\partial t}, \quad v_2 = fA \frac{\partial}{\partial z} + B \frac{\partial}{\partial t}, \quad v_3 = A \frac{\partial}{\partial y} + C \frac{\partial}{\partial z} + D \frac{\partial}{\partial t}, \\ v_4 &= f^{-1} \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + t \frac{\partial}{\partial z} \right) + E \frac{\partial}{\partial t}; \quad A, \dots, E \in \mathbb{R}. \end{aligned}$$

Consider a field of λ -frames (2.6) on $J^2(M)$, A, \dots, E being real-valued functions

on $J^2(M)$. We have

$$(2.7) \quad d\lambda = f_y dy \wedge dx + f_z dz \wedge dx + f_t dt \wedge dx$$

and

$$(2.8) \quad d\lambda(v_1, v_4) = ff_t A, \quad d\lambda(v_2, v_4) = f_z A + f^{-1}f_t B, \\ d\lambda(v_3, v_4) = f^{-1}(f_y A + f_z C + f_t D).$$

Lemma. *There are fields of λ -frames satisfying*

$$(2.9) \quad d\lambda(v_1, v_4) = 1, \quad d\lambda(v_2, v_4) = d\lambda(v_3, v_4) = 0.$$

A general field of λ -frames satisfying (2.9) is given by

$$(2.10) \quad v_1 = ff_t^{-1} \frac{\partial}{\partial t}, \quad v_2 = f_t^{-1} \frac{\partial}{\partial z} - f_z f_t^{-2} \frac{\partial}{\partial t}, \\ v_3 = f^{-1} f_t^{-1} \frac{\partial}{\partial y} - f^{-1} f_y f_t^{-2} \frac{\partial}{\partial t} + C f_t \left(f_t^{-1} \frac{\partial}{\partial z} - f_z f_t^{-2} \frac{\partial}{\partial t} \right), \\ v_4 = f^{-1} \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + t \frac{\partial}{\partial z} \right) + E \frac{\partial}{\partial t}.$$

Consider the vector fields

$$(2.11) \quad w_1 = ff_t^{-1} \frac{\partial}{\partial t}, \quad w_2 = f_t^{-1} \frac{\partial}{\partial z} - f_z f_t^{-2} \frac{\partial}{\partial t}, \\ w_3 = f^{-1} \left(f_t^{-1} \frac{\partial}{\partial y} - f_y f_t^{-2} \frac{\partial}{\partial t} \right), \quad w_4 = f^{-1} \left(\frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + t \frac{\partial}{\partial z} \right);$$

thus the general λ -frames satisfying (2.9) are

$$(2.12) \quad v_1 = w_1, \quad v_2 = w_2, \quad v_3 = w_3 + R w_2, \quad v_4 = w_4 + S w_1,$$

$R, S: J^2(M) \rightarrow \mathbb{R}$ being arbitrary functions. By a direct calculation,

$$(2.13) \quad [w_1, w_2] = -ff_t^{-2} f_{tt} w_2, \quad [w_1, w_3] = -(1 + ff_t^{-2} f_{tt}) w_3, \\ [w_1, w_4] = \{f^{-1} f_t^{-1} (f_z + f_{xt} + z f_{yt} + t f_{zt}) - \\ - f^{-2} (f_x + z f_y + t f_z)\} w_1 + w_2 - w_4, \\ [w_2, w_3] = f^{-1} (f_t^{-2} f_{yt} - f_y f_t^{-3} f_{tt}) w_2 + (f_z f_t^{-3} f_{tt} - f_t^{-2} f_{zt}) w_3, \\ [w_2, w_4] = \alpha_{24}^1 w_1 + \alpha_{24}^2 w_2 + w_3, \\ \alpha_{24}^1 = f^{-2} f_t^{-1} (f_{xz} + z f_{yz} + t f_{zz} + f_y) - f^{-2} f_t^{-2} f_z (f_{xt} + z f_{yt} + t f_{zt} + f_z), \\ \alpha_{24}^2 = f^{-1} f_t^{-1} (f_{xt} + z f_{yt} + t f_{zt} - f_z), \\ [w_3, w_4] = \alpha_{34}^1 w_1 + \alpha_{34}^2 w_2 + \alpha_{34}^3 w_3, \\ \alpha_{34}^2 = -f^{-2} f_y f_t^{-1}, \quad \alpha_{34}^3 = f^{-2} (f_x + z f_y + t f_z) + f^{-1} f_t^{-1} (f_{xt} + z f_{yt} + t f_{zt}).$$

From this,

$$(2.14) \quad [v_2, v_4] = c_{24}^1 v_1 + c_{24}^2 v_2 + v_3, \quad [v_3, v_4] = c_{34}^1 v_1 + c_{34}^2 v_2 + c_{34}^3 v_3$$

with

$$(2.15) \quad \begin{aligned} c_{24}^2 &= f^{-1}f_t^{-1}(f_{xt} + zf_{yt} + tf_{zt} - f_z) + ff_t^{-2}f_{tt}S - R, \\ c_{34}^3 &= f^{-2}(f_x + zf_y + tf_z) + f^{-1}f_t^{-1}(f_{xt} + zf_{yt} + tf_{zt}) + \\ &\quad + (1 + ff_t^{-2}f_{tt})S + R. \end{aligned}$$

Proposition 2.1. *On $J^2(M)$, let the Lagrangian (2.5) be given which does not satisfy*

$$(2.16) \quad 1 + 2ff_t^{-2}f_{tt} = 0.$$

Then there is, on $J^2(M)$, a unique field of λ -frames satisfying (2.9) and $c_{24}^2 = c_{34}^3 = 0$ in (2.14).

Proof. From (2.15) with $c_{24}^2 = c_{34}^3 = 0$, we may calculate R and S , and our field of λ -frames is given by (2.12) and (2.11). QED.

The special case is thus formed by Lagrangians with (2.16), i.e., by Lagrangians

$$(2.17) \quad \lambda = \{a(x, y, z) + b(x, y, z)t\}^{2/3} dx.$$

The coframe dual to (2.12) is given by

$$(2.18) \quad \begin{aligned} \omega^1 &= -Sf dx + f^{-1}f_y(dy - z dx) + f^{-1}f_z(dz - t dx) + f^{-1}f_t dt, \\ \omega^2 &= -Rff_t(dy - z dx) + f_t(dz - t dx), \\ \omega^3 &= ff_t(dy - z dx), \quad \omega^4 = \lambda = f dx. \end{aligned}$$

Especially, the forms

$$(2.19) \quad \begin{aligned} \omega^3 &= ff_t(dy - z dx), \quad \omega^2 \wedge \omega^3 = ff_t^2(dz - t dx) \wedge (dy - z dx), \\ \omega^1 \wedge \omega^2 \wedge \omega^3 \wedge \omega^4 &= ff_t^3 dx \wedge dy \wedge dz \wedge dt \end{aligned}$$

are invariant in all cases.

Reference

- [1] *R. B. Gardner: Differential Geometric Methods Interfacing Control Theory. In "Differential Geometric Control Theory", Progress in Math. 27, Birkhauser, Boston, 1983, 117–180.*

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