Aplikace matematiky

Václav Vítek

Periodic solutions of a weakly nonlinear hyperbolic equation in E_2 and E_3

Aplikace matematiky, Vol. 19 (1974), No. 4, 232-245

Persistent URL: http://dml.cz/dmlcz/103537

Terms of use:

© Institute of Mathematics AS CR, 1974

Institute of Mathematics of the Czech Academy of Sciences provides access to digitized documents strictly for personal use. Each copy of any part of this document must contain these *Terms of use*.



This document has been digitized, optimized for electronic delivery and stamped with digital signature within the project *DML-CZ: The Czech Digital Mathematics Library* http://dml.cz

PERIODIC SOLUTIONS OF A WEAKLY NONLINEAR HYPERBOLIC EQUATION IN E_2 AND E_3

VÁCLAV VÍTEK

(Received March 30, 1973)

1. Introduction. For n = 2 and 3 I investigate the classical solutions of the equation with constant coefficients

$$(1.1^{(n)}) \qquad \Box_n u + 2au_t + 2(B, \nabla_n u) + cu = h(t, x) + \varepsilon f(t, x, u, \varepsilon).$$

The initial conditions are

(1.2⁽ⁿ⁾)
$$u(0, x) = p(x), \quad u_t(0, x) = q(x), \quad x \in E_n.$$
In (1.1⁽ⁿ⁾)
$$\square_n \equiv \partial^2/\partial t^2 - \partial^2/\partial x_1^2 - \partial^2/\partial x_2^2 - \dots - \partial^2/\partial x_n^2,$$

$$\nabla_n = (\partial/\partial x_1, \partial/\partial x_2, \dots, \partial/\partial x_n),$$

$$B = (b_1, b_2, \dots, b_n),$$

$$x = (x_1, x_2, \dots, x_n).$$

 ε is a small parameter, $t \in E_1^+ = \langle 0, +\infty \rangle$.

The study of $(1.1^{(n)})$ -type equations was initiated in [2] by F. A. Ficken and B. A. Fleishman. In [3] J. Havlová studied $(1.1^{(1)})$ with a more general right-hand side, viz. that of $f(t, x, u, u_x, u_t, \varepsilon)$.

2. Preliminaries. Denote for $n \ge 2$ $E_{n+1}^+ = E_1^+ \times E_n$, $\omega^2 = ||B||^2 + c - a^2 = b_1^2 + b_2^2 + \ldots + b_n^2 + c - a^2$,

(2.1)
$$\omega = \frac{\sqrt{(\|B\|^2 + c - a^2)}}{i\sqrt{(a^2 - \|B\|^2 - c)}} \quad \text{if} \quad \omega^2 \ge 0,$$

 $\Omega = |\omega|$, i.e.,

(2.2)
$$\omega = \frac{\Omega}{i\Omega} \qquad \omega^2 \ge 0,$$

$$\omega^2 \ge 0,$$

$$\omega^2 < 0$$

(2.3)
$$\lambda = \begin{cases} a & \omega^2 \ge 0 \\ a - \Omega & \text{if } \omega^2 < 0 \end{cases}$$

Throughout the paper we assume

$$(2.4) \lambda > 0.$$

Let us note that this condition is equivalent to a > 0, $||B||^2 + c > 0$ in the case of $\omega^2 < 0$.

Further we denote

(2.5)
$$\Gamma^{2}(\tau, t; \xi, x) \equiv (t - \tau)^{2} - ||x - \xi||^{2},$$
$$\Gamma_{0}^{2} = \Gamma^{2}(0, t; \xi, x).$$

The facility of multiindex notation is used:

$$D_{i} = D_{i}^{1} = \partial/\partial x_{i}, \quad D_{x}^{\alpha} = D_{1}^{\alpha_{1}}D_{2}^{\alpha_{2}}...D_{n}^{\alpha_{n}}, \quad |\alpha| = \alpha_{1} + \alpha_{2} + ... + \alpha_{n}$$

with α_i non-negative integers.

If $k \ge 2$ is an integer we denote by $C^k(E_{n+1}^+)$ or briefly C^k the Banach space of real functions v(t, x) having continuous and bounded on E_{n+1}^+ partial derivatives $D_x^{\alpha}v$, $D_x^{\beta}v$, v_{tt} , $|\beta| < |\alpha| \le k$. Setting

(2.6)
$$|||v|||_{k,t} = \sup_{x \in E_n} \{ |D_x^{\alpha} v(t,x)|, |D_x^{\beta} v_t(t,x)|, |v_{tt}(t,x)|; |\beta| < |\alpha| \leq k \},$$

we define the norm in C^k by

(2.7)
$$\sup_{t\geq 0} |||v||_{k,t} = ||v||_{C^k}.$$

Further let us put

(2.8)
$$\|v\|_{k,t} = \sup_{x \in E_n} \{ |D_x^{\alpha} v(t, x)|; |\alpha| \le k \} .$$

In $(1.1^{(n)})$, $(1.2^{(n)})$ we introduce the variables

(2.9)
$$u = w \cdot \exp \{-at + (B, x)\},$$
$$v = e^{-(B,x)} \cdot u = e^{-at} \cdot w.$$

We get

$$(2.10^{(n)}) \qquad \qquad \square_{nW} + \omega^2 w = e^{at} \{ \delta(t, x) + \varepsilon f_1(t, x, w, \varepsilon) \},$$

$$w(0, x) = \varphi(x), \quad w_t(0, x) = \psi(x)$$

and

(2.11⁽ⁿ⁾)
$$\square_n v + 2av_t + (\omega^2 + a) v = \delta(t, x) + \varepsilon g(t, x, v, \varepsilon),$$
$$v(0, x) = \varphi(x), \quad v_t(0, x) = (\psi - a\varphi)(x)$$

with

(2.12)
$$\delta(t, x) = e^{-(B,x)} \cdot h(t, x),$$

$$g(t, x, v, \varepsilon) = f_1(t, x, w, \varepsilon) = e^{-(B,x)} \cdot f(t, x, u, \varepsilon),$$

$$\varphi(x) = e^{-(B,x)} \cdot p(x),$$

$$\psi(x) = e^{-(B,x)} \cdot (q + ap)(x).$$

We shall assume

 $(G^{(k,n)}, 1)$: For an arbitrary $R \ge 0$ there exists a positive constant A(R) such that for $(t, x, v, \varepsilon) \in E_{n+1}^+ \times \langle -R, R \rangle \times \langle -\varepsilon_0, \varepsilon_0 \rangle$ the partial derivatives

$$D_x^{\alpha} D_v^j g(t, x, v, \varepsilon), \quad |\alpha| + j \leq k$$

are continuous and bounded:

$$\left|D_x^{\alpha}D_v^j g(t, x, v, \varepsilon)\right| \leq A(R);$$

 $(G^{(k,n)} \cdot 2)$: For an arbitrary $R \ge 0$ the partial derivatives $D_x^{\alpha} D_v^j g(t, x, v, \varepsilon)$, $|\alpha| + j \le k$ satisfy the Lipschitz condition in v with B(R) as the Lipschitz constant; $(G^{(k,n)} \cdot 3)$: $D_x^{\alpha} \delta(t, x)$, $|\alpha| \le k$ are continuous and bounded on E_{n+1}^+ ;

 $(G^{(n)}, 4)$: The functions $\varphi(x)$ and $\psi(x)$ possess continuous and bounded on E_n partial derivatives up to the orders $\lfloor n/2 \rfloor + 2$ and $\lfloor n/2 \rfloor + 1$, respectively.

Specializing to the case n=2, $(2.10^{(2)})$ is equivalent to an integral equation (cf. [1]; $D_t = \partial/\partial t$)

$$(2.13) \quad w(t, x, y) = (2\pi)^{-1} \iiint_{G} \operatorname{ch} i\omega \Gamma \cdot \Gamma^{-1} \cdot e^{at} \cdot \left[\delta(\tau, \xi, \eta) + \varepsilon f_{1}(\tau, \xi, \eta, w(\tau, \xi, \eta), \varepsilon)\right] d\xi d\eta d\tau + (2\pi)^{-1} \iint_{B} \operatorname{ch} i\omega \Gamma_{0} \cdot \Gamma_{0}^{-1} \cdot \psi(\xi, \eta) d\xi d\eta + (2\pi)^{-1} D_{t} \left(\iint_{B} \operatorname{ch} i\omega \Gamma_{0} \cdot \Gamma_{0}^{-1} \cdot \varphi(\xi, \eta) d\xi d\eta\right).$$

G denotes the interior of the cone $\Gamma > 0$, $0 < \tau < t$ with B as the basis. The equi-

valence of $(2.10^{(2)})$ to (2.13) is meant in the sense that every solution of $(2.10^{(2)})$ solves (2.13) and vice versa. (Cf. [1].)

Introducing polar coordinates r, v in (2.13)

(2.14)
$$\xi = x + \varrho r \cdot \cos v, \quad \eta = y + \varrho r \cdot \sin v,$$
$$0 \le v < 2\pi, \quad \varrho = t - \tau \text{ or } t,$$

we get with regard to (2.9)

(2.15)
$$v(t, x, y) = (2\pi)^{-1} \cdot \int_{0}^{t} \int_{0}^{2\pi} \int_{0}^{1} \operatorname{ch} \left(i\omega \cdot (t - \tau) \sqrt{(1 - r^{2})}\right) \cdot (1 - r^{2})^{-1/2} \cdot e^{-a(t-\tau)} \cdot (t - \tau) \cdot \left[\delta(\tau, x + (t - \tau) r \cdot \cos v, y + (t - \tau) r \cdot \sin v) + \varepsilon g(\tau, x + (t - \tau) r \cdot \cos v, y + (t - \tau) r \cdot \sin v, v(\tau, x + (t - \tau) r \cdot \cos v, y + (t - \tau) r \cdot \sin v), \varepsilon)\right] \cdot r \, dr \, dv \, d\tau + (2\pi)^{-1} t e^{-at} \int_{0}^{2\pi} \int_{0}^{1} \operatorname{ch} i\omega t \sqrt{(1 - r^{2}) \cdot (1 - r^{2})^{-1/2}} \cdot (\psi + a\varphi) \cdot (x + tr \cdot \cos v, y + tr \cdot \sin v) \cdot r \, dr \, dv + (2\pi)^{-1} D_{t} \left(t e^{-at} \int_{0}^{2\pi} \int_{0}^{1} \operatorname{ch} i\omega t \sqrt{(1 - r^{2}) \cdot (1 - r^{2})^{-1/2}} \cdot (1 - r^{2})^{-1/2} \cdot (1 - r^{$$

In the case n = 3 the analogue of (2.15) is (cf. [1])

$$(2.16) \quad v(t,X) = \int_{0}^{t} e^{-at} (t-\tau)^{-1} D_{t} \left(\int_{0}^{t-\tau} r^{2} J_{0}(\omega\Gamma) \cdot Q(\delta + \varepsilon g(v)) (\tau, X, r) dr \right).$$

$$\cdot d\tau + e^{-at} \cdot t^{-1} \cdot D_{t} \left(\int_{0}^{t} r^{2} J_{0}(\omega\Gamma_{0}) Q(\psi + a\varphi) (X, r) dr \right) +$$

$$+ D_{t} \left(e^{-at} \cdot t^{-1} \cdot D_{t} \left(\int_{0}^{t} r^{2} J_{0}(\omega\Gamma_{0}) Q(\varphi) (X, r) dr \right) \right)$$

where X = [x, y, z],

(2.17)
$$Q(\delta + \varepsilon g(v))(\tau, X, r) = (4\pi)^{-1} \int_0^{2\pi} \int_0^{\pi} e^{at} \cdot \left[\delta(\tau, x + r \cdot \sin \theta \cdot \cos v, y + r \cdot \sin \theta \cdot \sin v, z + r \cdot \cos \theta\right] + \varepsilon g(\tau, x + r \cdot \sin \theta \cdot \cos v, y + r \cdot \sin \theta \cdot \sin v, z + r \cdot \cos \theta, v(\tau, x + r \cdot \sin \theta \cdot \cos v, y + r \cdot \sin \theta \cdot \sin v, z + r \cdot \cos \theta, \varepsilon\right] \cdot \sin \theta d\theta dv,$$

$$Q(\mu)(X, r) = (4\pi)^{-1} \int_0^{2\pi} \int_0^{\pi} \mu(r, \theta, \nu) \cdot \sin \theta \, d\theta \, d\nu \,,$$

$$\mu(r, \theta, \nu) = \mu(x + r \cdot \sin \theta \cdot \cos \nu, y + r \cdot \sin \theta \cdot \sin \nu, z + r \cdot \cos \theta) \,.$$

In what follows we make use of the standard contraction mapping principle:

Let $v \to V(v)$ be a mapping of a Banach space X into itself such that (i) $\|v\| \le R$, $R \ge 0$ implies $\|V(v)\| \le R$; (ii) there exists a real number θ , $0 < \theta < 1$ such that for $\|v_1\|$, $\|v_2\| \le R$ the inequality $\|V(v_1) - V(v_2)\| \le \theta$. $\|v_1 - v_2\|$ is valid. Then there exists a unique $v \in X$ such that v = V(v) and $\|v\| \le R$.

For the proof cf. e.g. [4].

3. n = 2. The starting point is the equation (2.15). We introduce special ad hoc integral operators similarly as in [2]. For $s(t, x, y) \in C^3(E_3^+)$ put

(3.1)
$$H_{2}(s)(\tau, t, x, y) = (2\pi)^{-1} (t - \tau) e^{-a(t - \tau)} \cdot \int_{0}^{2\pi} \int_{0}^{1} s(\tau, x + (t - \tau) r) \cdot d\tau$$

$$\cdot \cos v, y + (t - \tau) r \cdot \sin v \cdot ch i\omega(t - \tau) \sqrt{(1 - r^{2})} \cdot (1 - r^{2}) \cdot (1 - r^{2})^{-1/2} r dr dv \cdot dt$$

$$M_{2}(s)(t, x, y) = \int_{0}^{t} H_{2}(s)(\tau, t, x, y) d\tau \cdot dt$$

$$\Psi_{2}(\psi)(t, x, y) = H_{2}(s)(0, t, x, y), \psi(\xi, \eta) = s(0, \xi, \eta) \cdot dt$$

$$\Psi_{2}(\psi)(t, x, y) = (D_{t}\Psi_{2} + a\Psi_{2})(\psi)(t, x, y) \cdot dt$$

Then (2.15) may be written as

(3.2)
$$v(t, x, y) = (\varepsilon M_2(g(v)) + M_2(\delta) + \Phi_2(\varphi) + \Psi_2(\psi))(t, x, y) \equiv$$
$$\equiv V_2(v)(g, \delta, \varphi, \psi)(t, x, y).$$

To this equation we apply the contraction mapping principle. We choose $X = C^2(E_3^+)$ and

$$||v|| = ||v||_{C^{2}} = \sup_{(t,x,y)\in E_{3}^{+}} \{ |D_{1}^{\alpha_{1}}D_{2}^{\alpha_{2}}v(t,x,y)|, |D_{1}^{\beta_{1}}D_{2}^{\beta_{2}}v_{t}(t,x,y)|, |v_{tt}(t,x,y)|; |\beta| < |\alpha| \le 2 \}.$$

Using the inequalities

(3.3a)
$$\operatorname{ch}(i\omega(t-\tau)\sqrt{(1-r^2)}) = \operatorname{ch}(-\Omega(t-\tau)\sqrt{(1-r^2)}) = \operatorname{ch}(\Omega(t-\tau)\sqrt{(1-r^2)}) \le \exp(\Omega(t-\tau)\sqrt{(1-r^2)}) \le \exp(\Omega(t-\tau)),$$

 $|\operatorname{sh}(\Omega(t-\tau)\sqrt{(1-r^2)})| \le \exp(\Omega(t-\tau))$

for the case $\omega^2 < 0$,

(3.3b)
$$\left| \operatorname{ch} \left(i\omega(t - \tau) \sqrt{(1 - r^2)} \right) \right| = \left| \operatorname{ch} \left(i\Omega(t - \tau) \sqrt{(1 - r^2)} \right) \right| =$$

$$= \left| \cos \left(\Omega(t - \tau) \sqrt{(1 - r^2)} \right) \right| \le 1 ,$$

$$\left| \sin \left(\Omega(t - \tau) \sqrt{(1 - r^2)} \right) \right| \le 1$$

for the case $\omega^2 \ge 0$ and the identity

$$\int_{0}^{1} r(1-r^{2})^{-1/2} dr = 1$$

we get

(3.4)
$$|H_{2}(s)(\tau, t, x, y)| \leq ||s||_{0,\tau} \cdot (t - \tau) \cdot e^{-\lambda(t - \tau)},$$

$$|D^{\alpha} H_{2}(s)(\tau, t, x, y)| = |H_{2}(D^{\alpha}s)(\tau, t, x, y)| \leq$$

$$\leq ||D^{\alpha}s||_{0,\tau} \cdot (t - \tau) \cdot e^{-\lambda(t - \tau)} \leq ||s||_{2,\tau} \cdot (t - \tau) \cdot e^{-\lambda(t - \tau)}, \quad |\alpha| \leq 2.$$

(The symbol D^{α} denotes $D_x^{\alpha_1} D_y^{\alpha_2}$.) Differentiating with respect to t in (3.1) we get

(3.5)
$$|D^{\beta}D_{t} H_{2}(s) (\tau, t, x, y)| = |D_{t} H_{2}(D^{\beta}s) (\tau, t, x, y)| \leq$$

$$\leq ||D^{\beta}s||_{1,\tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{1,1}(t-\tau) \leq$$

$$\leq ||s||_{2,\tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{1,1}(t-\tau) , \quad |\beta| \leq 1 .$$

Here $P_{1,1}(\zeta)$ denote a polynomial of the first degree with non-negative coefficients. Similar estimates hold for $D_t^2 H_2(s)$ and $D_t^3 H_2(s)$.

From (3.4-5) it results

(3.6)
$$|M_2(s)(t, x, y)| \le \int_0^t |H_2(s)(\tau, t, x, y)| d\tau \le$$

$$\le \int_0^t ||s||_{0,\tau} \cdot e^{-\lambda(t-\tau)} \cdot (t-\tau) d\tau \le ||s||_0 \cdot m_0,$$

where

$$m_0 = \sup_{t>0} (\lambda^{-2} - \lambda^{-2}(1 + \lambda t) e^{-\lambda t}) > 0.$$

Similarly for $D^{\beta}D_t M_2(s)$ and $D_t^2 M_2(s)$. Finally, we thus obtain

$$||M_2(s)||_{C^2} \leq m \cdot ||s||_2,$$

m being a positive constant. For Ψ_2 and Φ_2 we have

(3.8)
$$\|\Psi(\psi)\|_{C^2} \leq K \cdot \|\psi\|_2 ,$$

$$\|\Phi_{2}(\varphi)\|_{C^{2}} \leq L \cdot \|\varphi\|_{3}.$$

These estimates used in (3.2) give

(3.10)
$$\|V_2(v)(g, \delta, \varphi, \psi)\|_{C^2} \le |\varepsilon| \cdot m \|g(v)\|_2 + m \|\delta\|_2 + L \|\varphi\|_3 + K \|\psi\|_2 \le$$

 $\le |\varepsilon| m A_1(R) + m \|\delta\|_2 + L \|\varphi\|_3 + K \|\psi\|_2.$

In the last inequality the estimate $||g(v)||_2 \le A_1(R)$ has been used based on the $(G^{(2,2)}, 1)$ assumption; $A_1(R) = A(R), (1+3R+R^2)$. Proceeding further we have

$$(3.11) ||V_2(v_1)(g, \delta, \varphi, \psi) - V_2(v_2)(g, \delta, \varphi, \psi)||_{C_2} \le |\varepsilon| m ||g(v_1) - g(v_2)||_2.$$

With the aid of $(G^{(2,2)}, 1, 2)$ we get first of all

$$||g(v_1) - g(v_2)||_{2,t} \le E(R) \cdot |||v_1 - v_2|||_{2,t}$$

with

(3.13)
$$E(R) = B(R)(1 + 3R + R^2) + A(R)(3 + 2R).$$

Taking the supremum over $t \ge 0$ and inserting the result into (3.11) we have

$$(3.14) ||V_2(v_1)(g, \delta, \varphi, \psi) - V_2(v_2)(g, \delta, \varphi, \psi)||_{C^2} \le |\varepsilon| m E(R) ||v_1 - v_2||_{C^2}.$$

If we now choose the numbers R, ε_1 , ε from the conditions ($\varepsilon_0 > 0$ given)

(3.15)
$$R > m \|\delta\|_2 + L \|\phi\|_3 + K \|\psi\|_2$$
, $0 < \varepsilon_1 < \min(\varepsilon_0, (m E(R))^{-1})$, $|\varepsilon| \le \min(\varepsilon_1, (R - m \|\delta\|_2 - L \|\phi\|_3 - K \|\psi\|_2) (m A_1(R))^{-1})$

then from (3.10) and (3.14) we see that the (i) and (ii) assumptions of the Contraction lemma are fulfilled. Hence we conclude

Theorem 1. (Existence and uniqueness of the solution.) Let the functions g, δ defined in (2.12) satisfy ($G^{(2,2)}$. 1-3). Let φ and ψ satisfy ($G^{(2)}$. 4). Assume further (3.15). Then the Cauchy problem (1.1⁽²⁾), (1.2⁽²⁾) has exactly one solution $u = u(t, x, y) \in C^2(E_3^+)$, $\|e^{-b_1x-b_2y} \cdot u\|_{C^2} \leq R$.

Remark. Later on we shall need the following modification of Theorem 1. Assume that g, δ satisfy $(G^{(3,2)} \cdot 1-3)$. Assume $\varphi = \psi = 0$. Let

(3.16)
$$R > m \|\delta\|_{3},$$

$$0 < \varepsilon'_{1} < \min(\varepsilon_{0}, (m E_{1}(R))^{-1}),$$

$$|\varepsilon| \le \varepsilon_{2} = \min(\varepsilon'_{1}, (R - m \|\delta\|_{3}) / m A_{2}(R)),$$

where

$$A_2(R) = A(R) \cdot (1 + 7R + 6R^2 + R^3), \quad E_1(R) = A(R) (7 + 12R + 3R^2) + B(R) (1 + 7R + 6R^2 + R^3).$$

Then the Cauchy problem $(1.1^{(2)})$ with zero initial conditions has exactly one solution

$$u \in C^3(E_3^+), \|e^{-b_1x-b_2y} \cdot u\|_{C^3} \le R.$$

The proof closely follows the ideas developed in the proof of Theorem 1 and is therefore omitted.

Theorem 2. (Continuous dependence of the solution on the initial conditions and on the right-hand side.) Let $R \ge 0$ and (3.15) hold. Let the functions $v_r(t, x, y)$, $||v_r||_{C^2} \le R$ (r = 1, 2, ...) satisfy the equation

$$(3.17) v_r = V_2(v_r) \left(g, \delta_r, \varphi_r, \psi_r\right).$$

Suppose that there exist functions δ , φ , ψ such that

(3.18)
$$\|\delta_r - \delta\|_2 \to 0$$
, $\|\varphi_r - \varphi\|_3 \to 0$, $\|\psi_r - \psi\|_2 \to 0$

when $r \to +\infty$. Then there exists a function v(t, x, y) which is the solution of (3.2) such that $||v_r - v||_{C^2} \to 0$ for $r \to +\infty$.

Proof. From (3.17), (3.7)-(3.9) and (3.12) we get

$$||v_{r} - v_{r'}||_{C^{2}} \leq |\varepsilon| \cdot m E(R) ||v_{r} - v_{r'}||_{C^{2}} + m||\delta_{r} - \delta_{r'}||_{C^{2}} + L||\varphi_{r} - \varphi_{r'}||_{3} + K||\psi_{r} - \psi_{r'}||_{2}$$

or

$$(1 - |\varepsilon| \cdot m E(R)) \cdot ||v_{r} - v_{r'}||_{C^{2}} \le m ||\delta_{r} - \delta_{r'}||_{2} + L ||\varphi_{r} - \varphi_{r'}||_{3} + K ||\psi_{r} - \psi_{r'}||_{2}.$$

Now $1 - |\varepsilon| m E(R)$ is positive by the assumption and the right-hand side tends to zero because of (3.18). Therefore $\{v_r\}_{r=1}^{\infty}$ is a Cauchy sequence in C^2 and there exists a function v(t, x, y), the C^2 -limit of $\{v_r\}_{r=1}^{\infty}$,

as $r \to +\infty$. Moreover, $||v||_{C^2} \le R$. On the other hand, letting $r \to +\infty$ in (3.17) and taking into account (3.18), (3.19) we see that v(t, x, y) is a solution of (3.2), q.e.d.

Theorem 3. (Asymptotic stability.) Assume $(G^{(2,2)}, 1, 2)$. Let the functions $v_i(t, x, y) \in C^2(E_3^+)$, i = 1, 2, $||v_i||_{C^2} \le R$ satisfy (3.2) with the initial conditions $\varphi_i \in C^3$, $\psi_i \in C^2$,

(3.20)
$$\|\varphi_i\|_3 \leq B, \|\psi_i\|_2 \leq B,$$

B > 0 a constant. Let E(R) be defined as in (3.13). Then there exist positive constants K, α such that for

$$|\varepsilon| < \min\left(\varepsilon_0, \alpha K^{-1}, (2E(R))^{-1}\right)$$

the inequality

(3.22)
$$|||v_1 - v_2|||_{2,t} \le K \cdot \exp(-(\alpha - |\varepsilon| K) t)$$

is valid.

Proof. Performing similar estimates which led to Theorem 1 (cf. (3.12)) we have with respect to (3.20)

$$(3.23) ||v_{1} - v_{2}||_{2,t} \leq 2|\varepsilon| \int_{0}^{t} ||v_{1} - v_{2}||_{2,\tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{1}(t-\tau) d\tau + + 2\{||\varphi_{1} - \varphi_{2}||_{3} \cdot e^{-\lambda t} \cdot R_{1}(t) + ||\psi_{1} - \psi_{2}||_{2} \cdot e^{-\lambda t} \cdot P_{1}(t)\} \leq \leq |\varepsilon| \cdot K \int_{0}^{t} |||v_{1} - v_{2}||_{2,t} \cdot e^{-\lambda(t-\tau)} d\tau + Ke^{-\alpha t}.$$

Applying Gronwall's lemma we get (3.22).

Theorem 4. (Existence and uniqueness of a T-periodic solution.) Assume $a \neq 0$, $b_1^2 + b_2^2 + c > 0$, $0 < T < +\infty$. Let h, f in $(1.1^{(2)})$ be T-periodic in t and let $(G^{(3,2)}: 1-3)$ be valid. Then for $|\varepsilon| \leq \varepsilon_2$ and for $R > m \|\delta\|_3$ (cf. (3.16)) the equation $(1.1^{(2)})$ has exactly one T-periodic in t solution $u(t, x, y) \in C^2$, $\|e^{-b_1x-b_2y}: u\|_{C^2} \leq R$.

Proof. Assume firstly a > 0. By Remark following Theorem 1, Eq. $(1.1^{(2)})$ has exactly one solution $u(t, x, y) \in C^3(E_3^+)$, $\|e^{-b_1x-b_2y} \cdot u\|_{C^3} \le R$, $u(0, x, y) = u_t(0, x, y) = 0$. Put

$$v(t, x, y) = e^{-b_1 x - b_2 y} \cdot u(t, x, y)$$
.

We have $v \in C^3$, $||v||_{C^3} \le R$ and v solves (3.2) with zero initial conditions. For j = 1, 2, ... let us define

(3.24)
$$\varphi_{j}(x, y) = v(jT, x, y),$$

$$(\psi_{j} - a\varphi_{j})(x, y) = v_{t}(jT, x, y),$$

$$v_{j}(t, x, y) = v(t + jT, x, y).$$

We have

$$||v_j||_{C^3} \le R.$$

Besides, φ_j , ψ_j satisfy $(G^{(2)} \cdot 4)$, $\|\varphi_j\|_3 \le B$, $\|\psi_j\|_2 \le B$. Because $v_j(t, x, y)$ is a solution of (3.2) with the initial conditions (3.24) we get by Theorem 3

$$|||v_i - v|||_{2,t} \to 0$$
 as $t \to +\infty$.

In other words to every $\eta \ge 0$ there exists $t_{\eta} \ge 0$ in such a way that for $t \ge t_{\eta}$ and

for j = 1, 2, ... we have $|||v_j - v|||_{2,t} \le \eta$. Choosing $mT \ge t_{\eta}$, m an integer and $t \ge 0$ we have $mT + t \ge t_{\eta}$. Therefore

$$||v_i - v||_{2, mT + t} \le \eta.$$

Choosing further integers p > q, pT, $qT \ge t_n$ we get by (3.26)

$$|||v_{p} - v_{q}||_{2,t} = |||v_{p-q} - v||_{2,qT+t} \le \eta.$$

Therefore

$$||v_p - v_q||_{C^2} = \sup_{t \ge 0} |||v_p - v_q||_{2,t} \le \eta$$

which proves that $\{v_j(t,x,y)\}_{j=1}^\infty$ forms a Cauchy sequence in C^2 . In this way the existence of a C^2 -limit function $\bar{v}(t,x,y)$ is established, $\|\bar{v}-v_j\|_{C^2}\to 0$ as $j\to +\infty$. From here and (3.25) it results further $\|\bar{v}\|_{C^2}\le R$. By (3.27) and (3.23) applied with $v_p, v_q, \varphi_p, \varphi_q, \psi_p, \psi_q$ instead of $v_1, v_2, \varphi_1, \varphi_2, \psi_1, \psi_2$, respectively, we get also the Cauchy character of $\{\varphi_p\}_{p=1}^\infty$, $\{\psi_p\}_{p=1}^\infty$ respectively in C^3 , C^2 . I. e., there exist $\bar{\varphi}\in C^3$ and $\bar{\varphi}\in C^2$ such that $\|\bar{\varphi}-\varphi_p\|_3\to 0$ as $p\to +\infty$ and $\|\bar{\psi}-\psi_p\|_2\to 0$ as $p\to +\infty$. Besides, $\|\bar{\varphi}\|_3\le B$, $\|\bar{\psi}\|_2\le B$. Further

$$\bar{v}(0, x, y) = \lim_{p \to +\infty} v_p(0, x, y) = \lim_{p \to +\infty} v(pT, x, y) = \lim_{p \to +\infty} \varphi_p(x, y) = \bar{\varphi}(x, y)$$

and similarly $\bar{v}_t(0, x, y) = (\bar{\psi} - a\bar{\varphi})(x, y)$. On the other hand, $\bar{v}(t, x, y)$ is a solution of (3.2) with $\bar{\varphi}$ and $\bar{\psi}$ as initial conditions. This follows from the identity

$$\bar{v} - V_2(\bar{v})(g, \delta, \bar{\varphi}, \bar{\psi}) = (\bar{v} - v_p) + (V_2(v_p)(g, \delta, \varphi_p, \psi_p) - V_2(\bar{v})(g, \delta, \bar{\varphi}, \bar{\psi}))$$

letting $p \to +\infty$ and using previous estimates.

T-periodicity of \bar{v} is guaranteed by the relation

$$\bar{v}(t+T, x, y) = \lim_{n \to +\infty} v_p(t+T, x, y) = \lim_{n \to +\infty} v_{p+1}(t, x, y) = \bar{v}(t, x, y).$$

Uniqueness: Would there exist another T-periodic solution of (3.2), say p(t, x, y), $||p||_{C^2} \le R$, then reasoning as in the case of $v_j(t, x, y)$ we get that the initial conditions for $p_j(t, x, y) \stackrel{\text{df}}{=} p(t + jT, x, y)$, viz. $p(0, x, y) = p(jT, x, y) = p_j(0, x, y)$ do converge towards $\overline{\varphi}(x, y)$ as $j \to +\infty$. We conclude that $p(0, x, y) = \overline{v}(0, x, y)$ and similarly $p_j(0, x, y) = \overline{v}(0, x, y)$.

On the other hand, by Theorem 3

as $t \to +\infty$. Because the initial conditions of \bar{v} and p coincide (3.28) cannot hold unless \bar{v} and p are identical in the whole range of t. From $\bar{v}(t,x,y) = e^{-b_1x-b_2y}$. $\bar{u}(t,x,y)$ we get a unique T-periodic in t solution $\bar{u}(t,x,y) \in C^2$ of $(1.1^{(2)})$, $\|e^{-b_1x-b_2y} \cdot \bar{u}\|_{C^2} \le R$.

II. The case a < 0 is changed to the former one by defining $f(t, x, y, u, \varepsilon)$ for negative values of t via T-periodicity. The substitution $t = -\tau$ in $(1.1^{(2)})$ then changes the u_t -coefficient sign and so the proof is reduced to part I.

4. n = 3. Following the ideas developed in Sec. 3 we define: for $s(t, X) \in C^2(E_4^+)$ (X = [x, y, z])

(4.1)
$$H_3(s)(\tau, t, X) = e^{-at}(t - \tau)^{-1} \cdot D_t \left(\int_0^{t-\tau} r^2 \cdot J_0(\omega \Gamma) \cdot Q(s)(\tau, X, r) dr \right)$$

with

$$Q(s)(\tau, X, r) = (4\pi)^{-1} \int_0^{2\pi} \int_0^{\pi} e^{a\tau} \cdot s(r) \cdot \sin \theta \, d\theta \, dv,$$

 $s(r) = s(\tau, x + r \cdot \sin \theta \cdot \cos \nu, y + r \cdot \sin \theta \cdot \sin \nu, z + r \cdot \cos \theta);$

(4.2)
$$M_{3}(s)(t, X) = \int_{0}^{t} H_{3}(s)(\tau, t, X) d\tau,$$

$$\Psi_{3}(\psi)(t, X) = H_{3}(s)(0, t, X) \text{ with } \psi(X) = s(0, X),$$

$$\Phi_{3}(\varphi)(t, X) = (D_{t}\Psi_{3} + a\Psi_{3})(\varphi)(t, X).$$

After that (2.16) may be rewritten in the form

$$(4.3) v(t,X) = (\varepsilon M_3(g(v)) + M_3(\delta) + \Phi_3(\varphi) + \Psi_3(\psi))(t,X) \equiv$$

$$\equiv V_3(v)(g,\delta,\varphi,\psi)(t,X).$$

Theorem 1'. Assume $(G^{(2,3)}, 1-3)$ and $(G^{(2)}, 4)$. Let

(4.4)
$$R > K \|\delta\|_2 + \varrho \|\varphi\|_3 + \mu \|\psi\|_2$$

with suitable positive constants K, ϱ , μ . Let

$$(4.5) |\varepsilon| \leq \min(\varepsilon_3, (R - K \|\delta\|_2 - \varrho \|\varphi\|_3 - \mu \|\psi\|_2) / K A_1(R)),$$

where $0 < \varepsilon_3 < \min(\varepsilon_0, (K E(R))^{-1})$, E(R) and $A_1(R)$ is defined as in Sec. 3. Then the problem $(1.1^{(3)})$, $(1.2^{(3)})$ has exactly one solution $u = u(t, X) \in C^2(E_4^+)$, $||e^{-b_1x - b_2y - b_3z}|$. $u||_{C^2} \le R$.

Proof. Similarly as in the proof of Theorem 1, estimates of the integral operators are required in order to verify the (i) and (ii) assumptions of the Contraction mapping principle.

1) The counterpart of the inequalities (3.3) is the inequality

(4.6)
$$2^{n} n! \left| J_{n}(\omega \Gamma) \cdot (\omega \Gamma)^{-n} \right| \leq \frac{1}{e^{\Omega(t-\tau)}} \quad \text{if} \quad \omega^{2} \geq 0 ,$$

$$n = 1, 2, \dots \text{ (Cf. [5].)}$$

- 2) For the total derivatives $D_t^k s$, k = 1, 2, 3 we get $|D_t^k s| \le 3^k$. $||s||_{k,t}$.
- 3) For $Q(D_t^k s)(\tau, X, t \tau)$ we have $|Q(D_t^k s)(\tau, X, t \tau)| \le 3^k \cdot e^{a\tau} \cdot ||s||_{k,\tau}, k = 0, 1, 2, 3.$
 - 4) Using the known relations (cf. [5])

$$J_n(\omega\Gamma) \cdot (\omega\Gamma)^{-n} \to (n! \cdot 2^n)^{-1}$$
 as $\Gamma \to 0$,
 $J'_n(z) = n \cdot J_n(z) \cdot z^{-1} - J_{n+1}(z)$, $n = 0, 1, 2, ...$

we get

$$|D_t^k J_0(\omega \Gamma)| \leq R_k(t-\tau)$$
 if $\omega^2 \geq 0$

and

$$|D_t^k J_0(\omega \Gamma)| \le R_k(t-\tau) \cdot e^{\Omega(t-\tau)}$$
 if $\omega^2 < 0$,

k = 1, 2, 3, 4. $R_k(\tau)$ are polynomials of the k-th degree with non-negative coefficients.

5) Denoting (k = 1, 2, 3, 4)

$$I_k = \int_0^{t-\tau} r^2 D_t^k J_0(\omega \Gamma) Q(s) (\tau, X, r) dr$$

we have

$$|I_k| \le S_{k+3}(t-\tau) \cdot e^{a\tau} \cdot ||s||_{2,\tau} \text{ if } \omega^2 \ge 0$$

and

$$\left|I_k\right| \leq S_{k+3}(t-\tau) \cdot e^{a\tau + \Omega(t-\tau)} \cdot \|s\|_{2,\tau} \quad \text{if} \quad \omega^2 < 0.$$

 $S_{k+3}(t-\tau)$ are polynomials of the (k+3)-rd degree with non-negative coefficients. With the aid of the above inequalities we easily deduce the estimates for the integral operators:

a)
$$|D^{\alpha} H_{3}(s) (\tau, t, X)| \leq ||D^{\alpha} s||_{0,\tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{3}(t-\tau) \leq$$

$$\leq ||s||_{2,\tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{3}(t-\tau), \quad |\alpha| \leq 2;$$

$$|D^{\beta} D_{t} H_{3}(s) (\tau, t, X)| \leq ||s||_{2,\tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{4}(t-\tau), \quad |\beta| \leq 1;$$

$$|D^{2} H_{3}(s) (\tau, t, X)| \leq ||s||_{2,\tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{5}(t-\tau);$$

$$|D^{3} H_{3}(s) (\tau, t, X)| \leq ||s||_{3,\tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{6}(t-\tau).$$

$$(P_{k} \text{ are polynomials of the } k\text{-th degree}, k = 3, 4, 5, 6.)$$

b) Using the inequalities stated in a) we get for M_3

$$|D^{\alpha} M_{3}(s)(t, X)| \leq \int_{0}^{t} ||s||_{2, \tau} \cdot e^{-\lambda(t-\tau)} \cdot P_{3}(t-\tau) d\tau \leq$$

$$\leq ||s||_{2} \cdot (r_{3} + e^{-\lambda t} Q_{3}(t)) \leq ||s||_{2} \cdot K_{0}, \quad |\alpha| \leq 2.$$

Here

$$K_0 = \sup_{t>0} (r_3 + e^{-\lambda t} Q_3(t));$$

 r_3 is a positive constant, $Q_3(t)$ a polynomial of the third degree. Further we have

$$|D^{\beta}D_{t} M_{3}(s)(t, X)| \leq ||s||_{2} \cdot K_{1},$$

$$|D^{2}_{t} M_{3}(s)(t, X)| \leq ||s||_{2} \cdot K_{2}.$$

Summarizing we have

$$||M(s)||_{C^2} \le ||s||_2 . K$$

with $K = \max(K_0, K_1, K_2)$.

c)
$$|D^{\alpha} \Psi_{3}(\psi)(t, X)| \leq ||\psi||_{2} \cdot e^{-\lambda t} \cdot P_{3}(t) \leq ||\psi||_{2} \cdot \mu_{0}, \quad |\alpha| \leq 2;$$

 $|D^{\beta}D_{t} \Psi_{3}(\psi)(t, X)| \leq ||\psi||_{2} \cdot e^{-\lambda t} \cdot P_{4}(t) \leq ||\psi||_{2} \cdot \mu_{1}, \quad |\beta| \leq 1;$
 $|D_{t}^{2} \Psi_{3}(\psi)(t, X)| \leq ||\psi||_{2} \cdot e^{-\lambda t} \cdot P_{5}(t) \leq ||\psi||_{2} \cdot \mu_{2};$

$$\|\Psi_{3}(\psi)\|_{C^{2}} \leq \mu \cdot \|\psi\|_{2}$$

with $\mu = \max (\mu_0, \mu_1, \mu_2)$.

d)
$$|D^{\alpha} \Phi_{3}(\varphi)(t, X)| \leq ||\varphi||_{3} \cdot e^{-\lambda t} \cdot (P_{4}(t) + a P_{3}(t)) \leq ||\varphi||_{3} \cdot (\mu_{1} + a\mu_{0}),$$
 $|\alpha| \leq 2;$
 $|D^{\beta} D_{t} \Phi_{3}(\varphi)(t, X)| \leq ||\varphi||_{3} \cdot e^{-\lambda t} \cdot (P_{5}(t) + a P_{4}(t)) \leq ||\varphi||_{3} \cdot (\mu_{2} + a\mu_{1}),$
 $|\beta| \leq 1;$
 $|D_{t}^{2} \Phi_{3}(\varphi)(t, X)| \leq ||\varphi||_{3} \cdot e^{-\lambda t} \cdot (P_{6}(t) + a P_{5}(t)) \leq \mu_{3} \cdot ||\varphi||_{3};$

$$(4.9)$$
 $||\Phi_{3}(\varphi)||_{C^{2}} \leq \varrho \cdot ||\varphi||_{3}$

with $\varrho = \max(\mu_1 + a\mu_0, \mu_2 + a\mu_1, \mu_3)$.

Using in (4.3) the derived estimates a) to d) we get

$$(4.10) ||V_{3}(v)(g, \delta, \varphi, \psi)||_{C^{2}} \leq |\varepsilon| \cdot K \cdot ||g(v)||_{2} + K \cdot ||\delta||_{2} + \varrho \cdot ||\varphi||_{3} + \mu \cdot ||\psi||_{2} \leq |\varepsilon| \cdot K \cdot A_{1}(R) + K ||\delta||_{2} + \varrho ||\varphi||_{3} + \mu ||\psi||_{2}$$

and

These inequalities with respect to (4.4) and (4.5) imply both (i) and (ii) conditions of the Contraction mapping principle.

In the same way it would be possible to derive analogous theorems 2', 3' and 4'. The author gratefully acknowledges Dr. O. Vejvoda for his encouragement and instruction and Dr. M. Kopáčková-Suchá for several helpful discussions.

Bibliography

- [1] R. Courant, D. Hilbert: Methoden der mathematischen Physik. Vol. 2, Berlin 1937.
- [2] F. A. Ficken, B. A. Fleishman: Initial value problems and time-periodic solutions for a non-linear wave equation. Comm. pure appl. math. 10 (1957), 331-356.
- [3] J. Havlová: Periodic solutions of a nonlinear telegraph equation. Čas. pěst. mat. 90 (1965), 273–289.
- [4] L. V. Kantorovič, G. P. Akilov: Funkcional'nyj analiz v normirovanych prostranstvach. Moskva 1959.
- [5] G. N. Watson: A treatise on the theory of Bessel functions. 2nd ed. Cambridge U.P. 1944.

Souhrn

PERIODICKÁ ŘEŠENÍ SLABĚ NELINEÁRNÍ HYPERBOLICKÉ ROVNICE V E_2 A E_3

VÁCLAV VÍTEK

Pro n=2 a 3 se dokazuje existence a jednoznačnost klasického periodického řešení rovnice

$$\Box_n u + 2au_t + 2(B, \nabla_n u) + cu = h(t, x) + \varepsilon f(t, x, u, \varepsilon)$$

 $(x = (x_1, x_2, ..., x_n))$ za předpokladu periodičnosti pravé strany.

Author's address: RNDr. Václav Vítek, CSc., INORGA (Ústav pro automatizaci řízení v průmyslu), Letenská 17, 118 06 Praha 1.