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# DUAL FINITE ELEMENT ANALYSIS FOR CONTACT PROBLEM <br> OF ELASTIC BODIES WITH AN ENLARGING CONTACT ZONE 

Tran Van Bon

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Summary. Dual finite element analysis of the contact problem of two elastic bodies with an enlarging contact zone is presented. Approximations of the solution are defined on two types of triangulations by piecewise constant stress fields. Convergence is proved in both cases.

Keywords: contact of elastic bodies, dual variational formulation, dual finite element method AMS Subject class.: 65 N 30, 73 C 99

## INTRODUCTION

Dual finite element analysis for the contact problem of two plane elastic bodies without friction was given in [1]-III by Haslinger and Hlaváček, but only for the case of a bounded contact zone. In the case of the contact zone enlarging during the deformation a primary finite element analysis was discussed in [1]-II. It is the aim of this paper to analyse the latter problem by a dual procedure. Two types of triangulations are used. The convergence of the dual finite approximations for each type of triangulations is proven.

## 1. THE DUAL VARIATIONAL FORMULATION

First of all we recall the primary variational formulation of a contact problem with an enlarging contact zone (see [1]-I, § 1 for details). We introduce the following notation: the norm and the semi-norm in the Sobolev space $H^{k}(\Omega)$ are denoted by $\|\cdot\|_{k},|\cdot|_{k}$, respectively. $P_{k}(T)$ denotes the set of all polynomials of order $k$ defined on $T$.

$$
\begin{gathered}
W=\left\{\mathbf{u} \mid \boldsymbol{u}=\left(\mathbf{u}^{\prime}, \mathbf{u}^{\prime \prime}\right) \in\left[H^{1}\left(\Omega^{\prime}\right)\right]^{2} \times\left[H^{1}\left(\Omega^{\prime \prime}\right)\right]^{2}\right\}, \\
V=\left\{\boldsymbol{u} \in W \mid \boldsymbol{u}^{\prime}=0 \text { on } \Gamma_{u}, u_{n}^{\prime \prime}=0 \text { on } \Gamma_{0}\right\}, \\
K_{\varepsilon}=\left\{\mathbf{v} \in V \mid v_{\xi}^{\prime \prime}(\eta)-v_{\xi}^{\prime}(\eta) \leqq \varepsilon(\eta) \text { for a.e. } \eta \in\langle a, b\rangle\right\},
\end{gathered}
$$

where $v_{\xi}^{M}(\eta)$ is the component of the displacement vector $\boldsymbol{v}^{M}$ in the direction $\xi$,
$\varepsilon(\eta)=f^{\prime}(\eta)-f^{\prime \prime}(\eta) ; f^{\prime}, f^{\prime \prime}$ are functions describing the arcs $\Gamma_{K}^{\prime}$ and $\Gamma_{K}^{\prime \prime}$, respectively (see Fig. 1), $f^{\prime}, f^{\prime \prime} \in C^{2}(\langle a, b\rangle)$.

$$
A(\mathbf{u}, \mathbf{v})=\int_{\Omega} c_{i j k l} e_{i j}(\mathbf{u}) e_{k l}(\mathbf{v}) \mathrm{d} x, \quad \Omega=\Omega^{\prime} \cup \Omega^{\prime \prime}
$$



Fig. 1
where $e_{i j}(\mathbf{u})=\frac{1}{2}\left(\partial u_{i} / \partial x_{j}+\partial u_{j} / \partial x_{i}\right)$ are the components of the strain tensor with respect to the displacement vector $u, c_{i j k l}$ are coefficients defined by the generalized Hooke's law. Assume that there is a positive constant $c_{0}$ such that each symmetric tensor of the 2 nd order $\tau=\left(\tau_{i j}\right), i, j=1,2$ satisfies

$$
c_{i j k l} \tau_{i j} \tau_{k l} \geqq c_{0} \tau_{i j} \tau_{i j} \text { for a.e. } x \in \Omega .
$$

Throughout this paper we shall use the adding convention, i.e., sum from 1 to 2 for any twice repeated index.

$$
L(\mathbf{v})=\int_{\Omega} F_{i} v_{i} \mathrm{~d} x+\int_{\Gamma_{\tau^{\prime} \cup \Gamma_{\tau^{\prime \prime}}}} P_{i} v_{i} \mathrm{~d} S,
$$

where $F_{i}, P_{i}$ are the components of the body forces and the surface loads, respectively.

$$
\mathscr{L}(\mathbf{v})=\frac{1}{2} A(\mathbf{v}, \mathbf{v},)-L^{\prime}(\mathbf{v}) .
$$

Then the primary problem is to find $\boldsymbol{u} \in K_{\varepsilon}$ such that

$$
\begin{equation*}
\mathscr{L}(\mathbf{u}) \leqq \mathscr{L}(\mathbf{v}) \quad \forall \mathbf{v} \in K_{\varepsilon} . \tag{1.1}
\end{equation*}
$$

The existence and uniqueness of a solution were discussed in [1]-I.
We now consider the sets

$$
\begin{gathered}
S=\left\{\mathscr{N}=\left(\mathscr{N}_{i j}\right), i, j=1,2 \mid \mathscr{N}_{i j} \in L_{2}(\Omega), \mathscr{N}_{i j}=\mathscr{N}_{j i}\right\}, \\
\mathscr{W}_{\varepsilon}=S \times K_{\varepsilon},
\end{gathered}
$$

and the functional

$$
\mathscr{H}([\mathscr{N}, \mathbf{v}] ; \lambda)=\mathscr{L}_{1}(\mathcal{N}, \mathbf{v})+\int_{\Omega} \lambda_{i j}\left(e_{i j}(\mathbf{v})-\mathscr{N}_{i j}\right) \mathrm{d} x ; \quad[\mathcal{N}, \mathbf{v}] \in \mathscr{W}_{\varepsilon}, \quad \lambda \in S,
$$

where

$$
\mathscr{L}_{1}(\mathcal{N}, \mathbf{v})=\frac{1}{2} \int_{\Omega} c_{i j k l} \mathscr{N}_{i j} \mathscr{N}_{k l} \mathrm{~d} x-L^{\prime}(\mathbf{v})
$$

Lemma 1.1.1.If $\left\{\left[\mathscr{N}^{*}, \mathbf{v}^{*}\right] ; \lambda^{*}\right\}$ is a saddle-point of the functional $\mathscr{H}$ on $\mathscr{W}_{\varepsilon} \times S$, i.e.,

$$
\mathscr{H}\left(\left[\mathscr{N}^{*}, \mathbf{v}^{*}\right] ; \mu\right) \leqq \mathscr{H}\left(\left[\mathscr{N}^{*}, \mathbf{v}^{*}\right] ; \lambda^{*}\right) \leqq \mathscr{H}\left([\mathcal{N}, \mathbf{v}] ; \lambda^{*}\right)
$$

holds for all $[\mathcal{N}, \mathbf{v}] \in \mathscr{W}_{\varepsilon}$ and $\mu \in S$, then a solution $\mathbf{u}$ of the primary problem (1.1) exists and

$$
\mathscr{N}^{*}=e(\mathbf{u}), \quad \mathbf{v}^{*}=\mathbf{u}, \quad \lambda^{*}=\tau(\mathbf{u}),
$$

where $e^{\prime}(\mathbf{u})$ and $\tau(\mathbf{u})$ are the strain and the stress tensor, respectively.
2. If $\mathbf{u}$ is a solution of the primary problem (1.1), then $\{[e(\mathbf{u}), \mathbf{u}] ; \tau(\mathbf{u})\}$ is a saddlepoint of $\mathscr{H}$ on $\mathscr{W}_{\varepsilon} \times S$.

Proof. 1. From the properties of the saddle-point we deduce that

$$
\begin{align*}
& \delta_{\lambda} \mathscr{H}\left(\left[\mathscr{N}^{*}, \mathbf{v}^{*}\right] ; \lambda^{*}\right)=0 \Leftrightarrow \mathcal{N}_{i j}^{*}=e_{i j}(\mathbf{v}),  \tag{1.2}\\
& \delta_{\mathcal{N}} \mathscr{H}\left[\left[\mathscr{N}^{*}, \mathbf{v}^{*}\right] ; \lambda^{*}\right)=0 \Leftrightarrow \lambda_{i j}^{*}=c_{i j k l} \mathscr{N}_{k l}^{*},  \tag{1.3}\\
& \delta_{v} \mathscr{H}\left(\left[\mathscr{N}^{*}, \mathbf{v}^{*}\right] ; \lambda^{*}\right)\left(\mathbf{v}-\mathbf{v}^{*}\right) \geqq 0 \quad \forall \mathbf{v} \in K_{\varepsilon}, \tag{1.4}
\end{align*}
$$

where e.g. $\delta_{\lambda} \mathscr{H}$ denotes the partial Gâteaux differential of $\mathscr{H}$ with respect to $\lambda$.
From (1.4) it follows that

$$
\begin{equation*}
\int_{\Omega} \lambda_{i j}^{*} e_{i j}\left(\mathbf{v}-\mathbf{v}^{*}\right) \mathrm{d} x \geqq L\left(\mathbf{v}-\mathbf{v}^{*}\right) \quad \forall \mathbf{v} \in K_{\varepsilon} . \tag{1.5}
\end{equation*}
$$

Making use of (1.5), (1.3) and (1.2), we obtain

$$
\left.\int_{\Omega} c_{i j k l} e_{i j}\left(\mathbf{v}-\mathbf{v}^{*}\right) e_{k l}(\mathbf{v}) \mathrm{d} x \geqq L^{\prime} \mathbf{v}-\mathbf{v}^{*}\right) \quad \forall \mathbf{v} \in K_{\varepsilon},
$$

i.e., $\boldsymbol{v}^{*}=\boldsymbol{u}$ is a solution of (1.1). Furthermore,

$$
\left.\mathscr{N}^{*}=e^{\prime} \boldsymbol{u}\right), \quad \lambda^{*}=\tau(\boldsymbol{u})
$$

2. Let $\boldsymbol{u}$ be a solution of (1.1). We have to verify that

$$
\begin{align*}
\mathscr{H}([e(\mathbf{u}), \mathbf{u}] ; \mu) & \left.\leqq \mathscr{H}\left(\left[e^{\prime}(\mathbf{u}), \mathbf{u}\right] ; \tau^{\prime} \mathbf{u}\right)\right) \leqq \mathscr{H}([\mathscr{N}, \mathbf{v}] ; \tau(\mathbf{u}))  \tag{1.6}\\
& \forall \mu \in S, \quad[\mathscr{N}, \mathbf{v}] \in \mathscr{W}_{\varepsilon} .
\end{align*}
$$

It is easy to show that the left inequality holds even with the equality sign. We have

$$
\begin{gather*}
\mathscr{H}([\mathscr{N}, \mathbf{v}] ; \tau(\mathbf{u}))-\mathscr{H}([e(\mathbf{u}), \mathbf{u}] ; \tau(\mathbf{u}))=  \tag{1.7}\\
=\mathscr{L}_{1}(\mathscr{N}, \mathbf{v})-\mathscr{L}_{1}(e(\mathbf{u}), \boldsymbol{u})+\int_{\Omega} \tau_{i j}(\mathbf{u})\left[e_{i j}(\mathbf{v})-\mathscr{N}_{i j}\right] \mathrm{d} x= \\
\left.=\frac{1}{2} \int_{\Omega} c_{i j k l}\left[\mathscr{N}_{i j}-e_{i j}(\mathbf{u})\right]\left[\mathscr{N}_{k l}-e_{k l}(\mathbf{u})\right] \mathrm{d} x+A(\mathbf{u}, \mathbf{v}-\mathbf{u})-L^{( } \mathbf{v}-\mathbf{u}\right) .
\end{gather*}
$$

Since $u$ is a solution of (1.1), we have

$$
\begin{equation*}
\left.A(\mathbf{u}, \mathbf{v}-\boldsymbol{u})-L^{\prime} \mathbf{v}-\mathbf{u}\right) \geqq 0 \quad \forall \mathbf{v} \in K_{\varepsilon} . \tag{1.8}
\end{equation*}
$$

From (1.7), (1.8) and the positive definiteness of the coefficients $c_{i j k l}$ the right inequality of (1.6) follows. Thus $\{[e(\mathbf{u}), \mathbf{u}] ; \tau(\mathbf{u})\}$ is a saddle-point of $\mathscr{H}$ on $\mathscr{W}_{\varepsilon} \times S$. Q.E.D.

Using the definition of $\mathscr{H}$ and the relation

$$
\operatorname{Sup}_{\lambda \in S} \int_{\Omega} \lambda_{i j}\left[e_{i j}(\mathbf{v})-\mathscr{N}_{i j}\right] \mathrm{d} x=\left\langle\begin{array}{ll}
0 & \text { if } \quad \mathcal{N}=e(\mathbf{v}), \\
+\infty & \text { if } \quad \mathscr{N} \neq e(\mathbf{v}),
\end{array}\right.
$$

we arrive at

$$
\begin{equation*}
\mathscr{L}(\boldsymbol{u})=\operatorname{Inf}_{\mathbf{v} \in K_{\varepsilon}} \mathscr{L}(\mathbf{v})=\operatorname{Inf}_{\substack{\mathbf{v} \in K_{s} \\ \mathcal{V} \in S}} \operatorname{Sup}_{\lambda \in S} \mathscr{H}([\mathscr{N}, \mathbf{v}] ; \lambda), \tag{1.9}
\end{equation*}
$$

where $\boldsymbol{u}$ is an arbitrary solution of (1.1).
The problem: to find

$$
\begin{equation*}
\operatorname{Sup}_{\lambda \in S} \operatorname{Inf}_{[\mathscr{N}, \mathbf{v}] \in \mathscr{W}_{\varepsilon}} \mathscr{H}([\mathscr{N}, \mathbf{v}] ; \lambda) \tag{1.10}
\end{equation*}
$$

will be called dual to the primary problem (1.1).
We shall reformulate the dual problem (1.10) in a simpler form. To this end, we introduce the decomposition

$$
\mathscr{H}([\mathscr{N}, \mathbf{v}] ; \lambda)=\mathscr{H}_{1}(\mathscr{N}, \lambda)+\mathscr{H}_{2}(\lambda, \mathbf{v})
$$

where

$$
\begin{gather*}
\mathscr{H}_{1}(\mathscr{N}, \lambda)=\frac{1}{2} \int_{\Omega} c_{i j k l} \mathcal{N}_{i j} \mathscr{N}_{k l} \mathrm{~d} x-\int_{\Omega} \lambda_{i j} \mathscr{N}_{i j} \mathrm{~d} x \\
\left.\mathscr{H}_{2}(\lambda, \mathbf{v})=\int_{\Omega} \lambda_{i j} e_{i j}(\mathbf{v}) \mathrm{d} x-L^{( } \mathbf{v}\right) \tag{1.11}
\end{gather*}
$$

The definitions of $\mathscr{H}_{1}, \mathscr{H}_{2}$ imply that

$$
\begin{equation*}
\operatorname{Inf}_{[\mathscr{N}, \mathbf{v}] \in \mathscr{W}_{\varepsilon}} \mathscr{H}([\mathscr{N}, \mathbf{v}] ; \lambda)=\operatorname{Inf}_{\mathscr{N} \in S} \mathscr{H}_{1}(\mathscr{N}, \lambda)+\operatorname{Inf}_{v \in K_{\varepsilon}} \mathscr{H}_{2}(\lambda, \mathbf{v}) . \tag{1.12}
\end{equation*}
$$

It is readily seen that

$$
\begin{equation*}
\operatorname{Inf}_{\mathcal{N} \in S} \mathscr{H}_{1}(\mathscr{N}, \lambda)=\mathscr{H}_{1}\left(\mathcal{N}^{*}, \lambda\right)=-\frac{1}{2} \int_{\Omega} a_{i j k l} \lambda_{i j} \lambda_{k l} \mathrm{~d} x \tag{1.13}
\end{equation*}
$$

where $a_{i j k l}$ are the coefficients of the inverse generalized Hooke's law, and $\mathscr{N}_{i j}^{*}=$ $=a_{i j k l} \lambda_{k l}$.

Since $K_{0} \subset K_{\varepsilon}$ and $K_{0}$ is a convex cone ( $K_{0}$ denotes $K_{\varepsilon}$ with $\varepsilon=0$ ), the infimum of $\mathscr{H}_{2}$ on $K_{\varepsilon}$ is finite only if

$$
\begin{equation*}
\mathscr{H}_{2}(\lambda, \mathbf{v}) \geqq 0 \quad \forall \boldsymbol{v} \in K_{0} . \tag{1.14}
\end{equation*}
$$

We set

$$
S_{F, P}^{+}=\left\{\lambda \in S \mid \mathscr{H}_{2}(\lambda, \mathbf{v}) \geqq 0 \quad \forall \mathbf{v} \in K_{0}\right\} .
$$

It is possible to show that for any $\lambda \in S_{F, P}^{+}$smooth enough we have

$$
\begin{gather*}
\frac{\partial \lambda_{i j}}{\partial x_{j}}+F_{\iota}=0 \text { in } \Omega, \quad i=1,2,  \tag{1.15}\\
\lambda_{i j} n_{j}=P_{i} \text { on } \Gamma_{\tau}, \quad i=1,2,  \tag{1.16}\\
T_{t}^{\prime \prime}\left(\lambda^{\prime \prime}\right)=0 \text { on } \Gamma_{0},  \tag{1.17}\\
T_{\eta}^{\prime}\left(\lambda^{\prime}\right)=0 \text { on } \Gamma_{k}^{\prime}, T_{\eta}^{\prime \prime}\left(\lambda^{\prime \prime}\right)=0 \text { on } \Gamma_{K}^{\prime \prime},  \tag{1.18}\\
-T_{\xi}^{\prime}\left(\lambda^{\prime}\right)\left(\cos \alpha^{\prime}\right)^{-1}=T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1} \leqq 0, \quad \eta \in\langle a, b\rangle \tag{1.19}
\end{gather*}
$$

where $\alpha^{M}, M={ }^{\prime}$," is the angle between the axis $\eta$ and the tangent to $\Gamma_{K}^{M}$ and $\left(\cos \alpha^{M}\right)^{-1}=\left[1-\left(\partial f^{M} / \partial \eta\right)^{2}\right]^{1 / 2}$.

In fact, from the condition $\lambda \in S_{F, P}^{+}$it follows that any $\boldsymbol{v} \in K_{0}$ satisfies

$$
-\int_{\Omega} v_{i} \frac{\partial \lambda_{i j}}{\partial x_{j}} \mathrm{~d} x+\int_{\partial \Omega^{\prime} \cup \partial \Omega^{\prime \prime}} v_{i} \lambda_{i j} n_{j} \mathrm{~d} S \geqq \int_{\Omega} F_{i} v_{i} \mathrm{~d} x+\int_{\Gamma_{\tau}} P_{i} v_{i} \mathrm{~d} S .
$$

Inserting $v_{i}^{M}= \pm \varphi_{i} \in C_{0}^{\infty}\left(\Omega^{M}\right), M==^{\prime \prime}, "$ we obtain (1.15). Then we arrive at

$$
\int_{\partial \Omega^{\prime} \cup \partial \Omega^{\prime \prime}} v_{i} \lambda_{i j} n_{j} \mathrm{~d} S \geqq \int_{r_{\tau}} P_{i} v_{i} \mathrm{~d} S .
$$

Choosing $v_{i}= \pm \psi_{i}$ such that the trace of $\psi_{i}$ has its support in $\Gamma_{\tau}$, we obtain (1.16).
Consequently, we deduce

$$
\begin{aligned}
& 0 \leqq \int_{\partial \Omega \div \Gamma_{\tau}} v_{i} \lambda_{i j} n_{j} \mathrm{~d} S=\int_{\Gamma_{\mathrm{K}^{\prime}}}\left[T_{\xi}^{\prime}\left(\lambda^{\prime}\right) v_{\xi}^{\prime}+T_{\eta}^{\prime}\left(\lambda^{\prime}\right) v_{\eta}^{\prime}\right] \mathrm{d} S+ \\
+ & \int_{\Gamma_{\mathrm{K}^{\prime \prime}}}\left[T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right) v_{\xi}^{\prime \prime}+T_{\eta}^{\prime \prime}\left(\lambda^{\prime \prime}\right) v_{\eta}^{\prime \prime}\right] \mathrm{d} S+\int_{\Gamma_{0}} T_{t}^{\prime \prime}\left(\lambda^{\prime \prime}\right) v_{t}^{\prime \prime} \mathrm{d} S \quad \forall \mathbf{v} \in K_{0} .
\end{aligned}
$$

Choosing $\boldsymbol{v}^{\prime} \equiv 0, \boldsymbol{v}^{\prime \prime}$ with $v_{n}^{\prime \prime}=0, v_{t}^{\prime \prime}= \pm \psi$ on $\Gamma_{0}$, where the support of $\psi$ is contained in $\Gamma_{0}$, we obtain (1.17). Thus we are left with the following inequality

$$
\begin{equation*}
\int_{\Gamma_{K^{\prime}}}\left[T_{\xi}^{\prime} v_{\xi}^{\prime}+T_{\eta}^{\prime} v_{\eta}^{\prime}\right] \mathrm{d} S+\int_{\Gamma_{K^{\prime \prime}}}\left[T_{\xi}^{\prime \prime} v_{\xi}^{\prime \prime}+T_{\eta}^{\prime \prime} v_{\eta}^{\prime \prime}\right] \mathrm{d} S \geqq 0 \quad \forall v \in K_{0} . \tag{1.20}
\end{equation*}
$$

Taking $\mathbf{v} \in V$ such that $v_{\eta}^{\prime}=v_{\eta}^{\prime \prime}=0$ and $v_{\xi}^{\prime}=v_{\xi}^{\prime \prime}= \pm \varphi$ for points on $\Gamma_{K}^{\prime} \cup \Gamma_{K}^{\prime \prime}$
having the same coordinates $\eta$ and with $\varphi \in C_{0}^{\infty}(\langle a, b\rangle)$ we have

$$
\begin{gathered}
0=\int_{\Gamma_{K^{\prime}}}\left[T_{\xi}^{\prime}\left(\lambda^{\prime}\right) \varphi\right] \mathrm{d} S+\int_{\Gamma_{K^{\prime \prime}}} T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right) \varphi \mathrm{d} S= \\
=\int_{a}^{b}\left[T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1}+T_{\xi}^{\prime}\left(\lambda^{\prime}\right)\left(\cos \alpha^{\prime}\right)^{-1}\right] \varphi \mathrm{d} \eta .
\end{gathered}
$$

Consequently, we deduce

$$
\begin{equation*}
-T_{\xi}^{\prime}\left(\lambda^{\prime}\right)\left(\cos \alpha^{\prime}\right)^{-1}=T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1} \quad \forall \eta \in\langle a, b\rangle . \tag{1.21}
\end{equation*}
$$

Choosing $v_{\xi}^{\prime}=v_{\xi}^{\prime \prime}=0, v_{\eta}^{\prime \prime}=0, v_{\eta}^{\prime}= \pm \varphi$ on $\Gamma_{K}^{\prime}$, we obtain from (1.20) that $T_{\eta}^{\prime}\left(\lambda^{\prime}\right)=0$ on $\Gamma_{K}^{\prime}$. The condition $T_{\eta}^{\prime}\left(\lambda^{\prime \prime}\right)=0$ on $\Gamma_{K}^{\prime \prime}$ can be deduced in the same way. Thus, (1.20) yields that

$$
\int_{a}^{b}\left[T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1}\right]\left(v_{\xi}^{\prime \prime}-v_{\xi}^{\prime}\right) \mathrm{d} \eta \geqq 0 \quad \forall \mathbf{v} \in K_{0}
$$

By virtue of the condition $v_{\xi}^{\prime \prime}-v_{\xi}^{\prime} \leqq 0 \forall \mathbf{v} \in K_{0}$, (1.19) follows.
From (1.15) up to (1.18) and (1.21) it is readily seen that for any $v \in K_{\varepsilon}, \lambda \in S_{F, P}^{+}$ we have

$$
\begin{equation*}
\mathscr{H}_{2}(\lambda, \boldsymbol{v})=\int_{a}^{b} T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1}\left(v_{\xi}^{\prime \prime}-v_{\xi}^{\prime}\right) \mathrm{d} \eta . \tag{1.22}
\end{equation*}
$$

We can show that for any $\lambda \in S_{F, P}^{+}$,

$$
\begin{equation*}
\operatorname{Inf}_{\mathbf{v} \in \mathrm{K}_{\varepsilon}} \mathscr{H}_{2}(\lambda, \mathbf{v})=\int_{a}^{b} T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1} \varepsilon(\eta) \mathrm{d} \eta . \tag{1.23}
\end{equation*}
$$

In fact, from (1.22), (1.19) and the definition of $K_{\varepsilon}$, it follows that

$$
\begin{equation*}
\mathscr{H}_{2}(\lambda, \mathbf{v}) \geqq \int_{a}^{b} T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1} \varepsilon(\eta) \mathrm{d} \eta \quad \forall \mathbf{v} \in K_{\varepsilon}, \quad \lambda \in S_{F, P}^{+} . \tag{1.24}
\end{equation*}
$$

Let the functions $f^{\prime}$ and $f^{\prime \prime}$ describing $\Gamma_{K}^{\prime}$ and $\Gamma_{K}^{\prime \prime}$, respectively, be twice continuously differentiable on $\langle a, b\rangle$. There exists a function $\boldsymbol{u}_{0}^{\prime \prime} \in\left[H^{1}\left(\Omega^{\prime \prime}\right)\right]^{2}$ such that $u_{0 n}^{\prime \prime}=0$ on $\Gamma_{0}, u_{0 \xi}^{\prime \prime}=\varepsilon(\eta)$ on $\langle a, b\rangle$. Choosing $\mathbf{u}_{0}=\left(\mathbf{u}_{0}^{\prime}, \mathbf{u}_{0}^{\prime \prime}\right)$, where $\mathbf{u}_{0}^{\prime} \equiv 0$, we have $u_{0} \in K_{\varepsilon}$ and

$$
\begin{equation*}
\mathscr{H}_{2}\left(\lambda, u_{0}\right)=\int_{a}^{b} T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1} \varepsilon(\eta) \mathrm{d} \eta \tag{1.25}
\end{equation*}
$$

Then (1.23) follows from (1.24) and (1.25).
Combining (1.23), (1.12) with (1.13) we arrive at

$$
\begin{equation*}
\operatorname{Sup}_{\lambda \in S} \operatorname{Inf}_{[\cdot \mathcal{V}, \mathbf{v}] \in \mathscr{W}_{\varepsilon}} \mathscr{H}([\mathcal{N}, \mathbf{v}] ; \lambda)=-\operatorname{Inf}_{\lambda \in S_{F, P}^{+}} \mathscr{S}_{1}(\lambda), \tag{1.26}
\end{equation*}
$$

where

$$
\mathscr{S}_{1}(\lambda)=\frac{1}{2} \int_{\Omega} a_{i j k l} \lambda_{i j} \lambda_{k l} \mathrm{~d} x-\int_{a}^{b} T_{\xi}^{\prime \prime}\left(\lambda^{\prime \prime}\right)\left(\cos \alpha^{\prime \prime}\right)^{-1} \varepsilon(\eta) \mathrm{d} \eta .
$$

We introduce the following notation for $\sigma, \tau \in S$ :

$$
\begin{gathered}
\langle\sigma, \tau\rangle=\int_{\Omega} \sigma_{i j} \tau_{i j} \mathrm{~d} x, \quad(\sigma, \tau)=\left\langle c^{-1} \sigma, \tau\right\rangle, \\
\|\tau\|_{0}^{2}=\langle\tau, \tau\rangle, \quad\|\tau\|^{2}=(\tau, \tau)
\end{gathered}
$$

where $c: S \rightarrow S$ is the isomorphism defined by the generalized Hooke's law:

$$
\sigma=c e \Leftrightarrow \sigma_{i j}=c_{i j k l} e_{k l}
$$

Then, making use of (1.25) and (1.11), we may write

$$
\mathscr{S}_{1}(\lambda)==\frac{1}{2}\|\lambda\|^{2}-\left\langle e\left(\boldsymbol{u}_{0}\right), \lambda\right\rangle+L\left(\boldsymbol{u}_{0}\right)
$$

and the dual problem (1.10) is reformulated as follows: to find $\lambda^{*} \in S_{F, P}^{+}$such that

$$
\begin{equation*}
\mathscr{S}\left(\lambda^{*}\right) \leqq \mathscr{S}(\lambda) \quad \forall \lambda \in S_{F, P}^{+}, \tag{1.27}
\end{equation*}
$$

where $\mathscr{S}(\lambda)=\frac{1}{2}\|\lambda\|^{2}-\left\langle e\left(u_{0}\right), \lambda\right\rangle$.
Lemma 1.2. Let $\mathbf{u}$ be a solution of the primary problem (1.1). Then $\lambda^{*}=\left(\lambda_{i j}^{*}\right)=$ $=\left(c_{i j k l} e_{k l}(\mathbf{u})\right)$ is the unique solution of the dual problem (1.27).
Proof. It is easy to see that $\lambda^{*} \in S_{F, P}^{+}$since $\mathbf{v}=\boldsymbol{u}+\mathbf{w} \in K_{\varepsilon}$ for any $\boldsymbol{w} \in K_{0}$. For any $\lambda \in S_{F, P}^{+}$we have

$$
\begin{gathered}
\mathscr{S}(\lambda)-\mathscr{S}\left(\lambda^{*}\right)=\frac{1}{2} \int_{\Omega} a_{i j k l}\left(\lambda_{i j}-\lambda_{i j}^{*}\right)\left(\lambda_{k l}-\lambda_{k l}^{*}\right) \mathrm{d} x+ \\
+\int_{\Omega} a_{i j k l} \lambda_{i j}^{*}\left(\lambda_{k l}-\lambda_{k l}^{*}\right) \mathrm{d} x-\left\langle e\left(\mathbf{u}_{0}\right), \lambda-\lambda^{*}\right\rangle \geqq \\
\geqq \frac{a_{0}}{2} \int_{\Omega}\left(\lambda_{i j}-\lambda_{i j}^{*}\right)\left(\lambda_{k l}-\lambda_{k l}^{*}\right) \mathrm{d} x+\int_{\Omega} e_{i j}(\mathbf{u})\left(\lambda_{i j}-\lambda_{i j}^{*}\right) \mathrm{d} x-\left\langle e\left(\mathbf{u}_{0}\right), \lambda-\lambda^{*}\right\rangle \geqq \\
\geqq\left\langle e\left(\mathbf{u}-\mathbf{u}_{0}\right), \lambda-\lambda^{*}\right\rangle,
\end{gathered}
$$

where we have used the ellipticity of the coefficients $a_{i j k l}$. The solution $\boldsymbol{u}$ of the primary problem satisfies

$$
\left.\left\langle\lambda^{*}, e(\mathbf{v}-\mathbf{u})\right\rangle \geqq L_{( }^{\prime} \mathbf{v}-\mathbf{u}\right) \quad \forall \mathbf{v} \in K
$$

Since $u_{0} \in K$ we have

$$
\left\langle e\left(\boldsymbol{u}-\boldsymbol{u}_{0}\right),-\lambda^{*}\right\rangle+L\left(\boldsymbol{u}-\boldsymbol{u}_{0}\right) \geqq 0 .
$$

From $\lambda \in S_{F, P}^{+}$and $\boldsymbol{u}-\boldsymbol{u}_{0} \in K_{0}$ it follows that

$$
0 \leqq \mathscr{H}_{2}\left(\lambda, \boldsymbol{u}-\boldsymbol{u}_{0}\right)=\left\langle e\left(\boldsymbol{u}-\mathbf{u}_{0}\right), \lambda\right\rangle-L\left(\boldsymbol{u}-\boldsymbol{u}_{0}\right) .
$$

Adding the two inequalities we arrive at

$$
\left\langle e\left(\boldsymbol{u}-\boldsymbol{u}_{0}\right), \lambda-\lambda^{*}\right\rangle \geqq 0 \quad \forall \lambda \in S_{F, \boldsymbol{P}}^{+},
$$

i.e.,

$$
\mathscr{S}(\lambda) \geqq \mathscr{S}\left(\lambda^{*}\right) \quad \forall \lambda \in S_{F, P}^{+}
$$

This implies that $\lambda^{*}$ solves the problem (1.27). Moreover, since the functional $\mathscr{S}(\lambda)$ is strictly convex on $S_{F, P}^{+}$, the uniqueness of $\lambda^{*}$ follows. Q.E.D.

Remark 1.1. Let $\boldsymbol{u}, \lambda^{*}$ be as in Lemma 1.2. Making use of (1.9), (1.25) and the relation between $\mathscr{S}_{1}(\lambda)$ and $\mathscr{S}(\lambda)$, it is possible to show that

$$
\mathscr{L}(\mathbf{u})+L\left(\mathbf{u}_{0}\right)+\mathscr{S}\left(\lambda^{*}\right)=0 .
$$

Remark 1.2. A result analogous to that of Lemma 1.5 in the paper [1]-III can be obtained. Namely, if a solution $\mathbf{u}$ of the primary problem (1.1) exists, then the set $S_{F, P}^{+}$is non-empty, convex and closed in $S$. Existence and uniqueness of a solution of (1.27) can be shown directly by using the strict convexity and the lower weak semicontinuity of the functional $\mathscr{S}(\lambda)$.

We here emphasize the fact that the dual problem (1.27) is uniquely solvable if the primary problem possesses at least one solution. The existence of the solution of the dual problem, however, can be proved directly, if in some way we show the non-emptiness of the set $S_{F, P}^{+}$. Thus, the dual problem may have a solution even in some cases when the primary problem has none.

## 2. APPROXIMATION OF THE DUAL PROBLEM

As in [2] it is possible to approximate the solution of the dual problem by means or piecewise constant stress tensors. Here we follow the procedure suggested by Haslinger and Hlaváček [2].
Let us consider a triangulation $\mathscr{T}_{h}^{M}$ of $\Omega^{M}, M={ }^{\prime}$, " such that the triangles adjacent to the boundaries may have a curved side along the boundary, and the nodes on $\Gamma_{K}^{M}$ are on lines parallel to the $\xi$ axis. If a curved triangle $T \in \mathscr{T}_{h}=\mathscr{T}_{h}^{\prime} \cup \mathscr{T}_{h}^{\prime \prime}$ adjacent to $\Gamma_{K}^{M}$ is convex, it will be divided by the chord into the "straight" triangle $T_{0}$ and the segment $T_{s}$ such that $T=T_{0} \cup T_{s}$. If $T_{c} \in \mathscr{T}_{h}$ is not convex, one of its sides must be parallel to $\xi$ axis. We define

$$
\begin{aligned}
& V_{h}=\left\{\mathbf{v} \in V|\mathbf{v}|_{T_{\mathrm{c}}} \in\left[P_{1}\left(T_{0}\right)\right]^{2} \quad \forall T_{0} \subset T \in \mathscr{T}_{h} \quad \text { adjacent to } \Gamma_{K}^{M},\right. \\
& \left.\left(\frac{\partial \mathbf{v}}{\partial \xi}\right)\right|_{T_{s}}=0 \quad \forall T_{s} \subset T \in \mathscr{T}_{h} \quad \text { adjacent to } \Gamma_{K}^{M}, \\
& \left.\left(\frac{\partial \mathbf{v}}{\partial \xi}\right)\right|_{T_{c}}=0 \quad \text { for each non-convex triangle } T_{c} \text { adjacent to } \Gamma_{K}^{M}, \\
& \left.\left.v\right|_{T} \in\left[P_{1}(T)\right]^{2} \text { for the other triangles } T\right\} .
\end{aligned}
$$

We introduce the sets

$$
\begin{aligned}
& S_{h}=\left\{\tau \in S \mid \tau_{i j} \in P_{0}\left(T^{*}\right) ; \quad T^{*}=T, T_{0}, T_{s}, T_{c} ; \quad i, j=1,2\right\}, \\
& K_{0 h}=\left\{\mathbf{v} \in V_{h} \mid v_{\xi}^{\prime \prime}-v_{\xi}^{\prime} \leqq 0 \quad \forall \eta \in\langle a, b\rangle \quad \text { on } \quad \Gamma_{K}^{\prime} \cup \Gamma_{K}^{\prime \prime}\right\} .
\end{aligned}
$$

Then we define approximations $S_{F, P, h}^{+}$of $S_{F, P}^{+}$as follows:

$$
S_{F, P, h}^{+}=\left\{\lambda \in S_{h} \mid\left\langle\lambda, e\left(\mathbf{v}_{h}\right)\right\rangle \geqq L\left(\mathbf{v}_{h}\right) \quad \forall \mathbf{v}_{h} \in K_{0 h}\right\} .
$$

It is readily seen that the condition $v_{\xi}^{\prime \prime}-v_{\xi}^{\prime} \leqq 0, v \in V_{h}$ holds everywhere on $\langle a, b\rangle$ if only it holds for the nodes lying on $\Gamma_{K}^{M}$.
Let the function $\mathbf{u}_{0}=\left(0, \mathbf{u}_{0}^{\prime \prime}\right)$ be known (see (1.25)).
Then the approximate dual problem is formulated as follows: to find $\lambda_{h} \in S_{F, P, h}^{+}$ such that

$$
\begin{equation*}
\mathscr{P}\left(\lambda_{h}\right) \leqq \mathscr{P}(\lambda) \quad \forall \lambda \in S_{F, P, h}^{+} . \tag{2.1}
\end{equation*}
$$

We introduce the projection mapping $r_{h}: S \rightarrow S_{h}$ defined by the relation

$$
\begin{equation*}
\left\langle\tau-r_{h} \tau, \chi_{h}\right\rangle=0 \quad \forall \chi_{h} \in S_{h} . \tag{2.2}
\end{equation*}
$$

Lemma 2.1. Let $\tau \in S_{F, P}^{+}$. Then $r_{h} \tau \in S_{F, P, h}^{+}$.
Proof. It is readily seen that

$$
\begin{gather*}
\left.\left(r_{h} \tau\right)\right|_{T^{*}}=\left(\operatorname{mes} T^{*}\right)^{-1} \int_{T^{*}} \tau \mathrm{~d} x, \quad \forall T^{*}=T, T_{0}, T_{s}, T_{c} \in \mathscr{T}_{h},  \tag{2.3}\\
\left\|r_{h} \tau-\tau\right\|_{0} \rightarrow 0, \quad h \rightarrow 0 .
\end{gather*}
$$

Assume that $\mathbf{v}_{h} \in K_{0 h}$. Then $e\left(\mathbf{v}_{h}\right) \in S_{h}$, and using (2.2) we obtain

$$
\left\langle r_{h} \tau, e\left(\mathbf{v}_{h}\right)\right\rangle=\left\langle\tau, e\left(\mathbf{v}_{h}\right)\right\rangle \geqq L\left(\mathbf{v}_{h}\right)
$$

because of $K_{0 h} \subset K_{0}, \tau \in S_{F, P}^{+}$. This yields $r_{h} \tau \in S_{F, P, h}^{+}$.
Q.E.D.

Next we recall some results of [2].
Lemma 2.2. Let $\Gamma_{K}^{\prime} \cap \bar{\Gamma}_{u}=\emptyset, \Gamma_{K}^{\prime \prime} \cap \bar{\Gamma}_{0}=\emptyset$. Assume that the number of points of $\bar{\Gamma}_{\tau} \cap \bar{\Gamma}_{u}, \bar{\Gamma}_{\tau} \cap \bar{\Gamma}_{0}$ is finite, and $f^{M} \in C^{m}(\langle a-\delta, b+\delta\rangle), \delta>0, m \geqq 1, M={ }^{\prime},{ }^{\prime \prime}$. Then the set

$$
\mathscr{K}_{m}=K_{0} \cap\left[C^{m}\left(\bar{\Omega}^{\prime}\right)\right]^{2} \times\left[C^{m}\left(\bar{\Omega}^{\prime \prime}\right)\right]^{2}
$$

is dense in $K_{0}$.
Lemma 2.3. Let $\mathbf{v} \in\left[H^{2}\left(\Omega^{\prime}\right)\right]^{2} \times\left[H^{2}\left(\Omega^{\prime \prime}\right)\right]^{2}, f^{M} \in C^{2}(\langle a, b\rangle)$. Define the Lagrange linear interpolation $\mathbf{v}_{I} \in V_{h}$ as follows: If $A_{i}$ is a node of a curved non-convex triangle, $\mathbf{v}_{I}\left(A_{i}\right)=\boldsymbol{v}\left(D_{i}\right)$, where $D_{i}$ is the projection of $A_{i}$ onto $\Gamma_{K}^{M}$ in the direction $\xi$; $\mathbf{v}_{I}\left(A_{i}\right)=\boldsymbol{v}\left(A_{i}\right)$ at the other nodes. We construct $\mathbf{v}_{I}$ such that $\mathbf{v}_{I} \in V_{h}$, i.e., $\mathbf{v}_{I}$ is a piecewise linear function which is extended continuously by constants on segments $T_{s}$ and on non-convex curved triangles. Then

$$
\left\|\mathbf{v}_{I}-\mathbf{v}\right\|_{1, \Omega^{\prime} \cup \Omega^{\prime \prime}} \rightarrow 0 \quad \text { for } \quad h \rightarrow 0
$$

for any regular system of triangulations.
We can prove the following result about the convergence of $\lambda_{h}$ to $\lambda^{*}$.
Theorem 1. Let $f^{M} \in C^{2}, M={ }^{\prime}, "$ in a neighbourhood of the interval $\langle a, b\rangle$
and let the same assumptions on $\Gamma_{K}^{M}, \Gamma_{u}, \Gamma_{0}, \Gamma_{\tau}$ as in Lemma 2.2 hold. Assume that $S_{F, P}^{+} \neq \emptyset$.

Then the approximate dual problem (2.1) possesses a unique solution $\lambda_{h}$ and

$$
\begin{equation*}
\left\|\lambda_{h}-\lambda^{*}\right\|_{0} \rightarrow 0, \quad h \rightarrow 0 \tag{2.4}
\end{equation*}
$$

holds for any regular family of triangulations $\mathscr{T}_{h}$, where $\lambda^{*}$ is the solution of the dual problem (1.27).

Proof. The set $S_{F, P, h}^{+}$is convex and closed. Lemma 2.1 implies that it is non-empty. The functional $\mathscr{S}(\lambda)$ is strictly convex and lower weakly semicontinuous on $S_{F, P, h}^{+}$. The existence and uniqueness of the solution of (2.1) follows.

To show (2.4), we shall apply an abstract theorem about the convergence of the Ritz-Galerkin method (see e.g. [3], Chapt. 4). We only have to prove the following two conditions:
(i) $\exists\left\{\tau_{h}\right\}, \tau_{i h} \in S_{F, P, h}^{+}, \tau_{h} \rightarrow \lambda^{*}$ in $S$ for $h \rightarrow 0$,
(ii) $\tau_{h} \in S_{F, P, h}^{+}, \tau_{h} \rightarrow \tau$ (weakly) in $S$ implies $\tau \in S_{F, P}^{+}$.

Choosing $\tau_{h}=r_{h} \lambda^{*}$, making use of Lemma 2.1 and (2.3) we deduce that (i) holds.
Now let $\boldsymbol{v}$ be an arbitrary element of $K_{0}$. By virtue of Lemma 2.2 we find $\mathbf{v}_{\gamma} \in \mathscr{K}_{2}$ such that for any $\gamma>0$,

$$
\left\|\mathbf{v}_{\gamma}-\boldsymbol{v}\right\|_{1, \Omega^{\prime} \cup \Omega^{\prime \prime}}<\gamma .
$$

Using Lemma 2.3 we arrive at

$$
\left\|\mathbf{v}_{\gamma I}-\mathbf{v}_{\gamma}\right\|_{1, \Omega^{\prime} \cup \Omega^{\prime \prime}} \rightarrow 0, \quad h \rightarrow 0 .
$$

It is clear that $\boldsymbol{v}_{\gamma I} \in K_{0 h}$. Then we obtain

$$
\left\|\boldsymbol{v}_{\gamma I}-\mathbf{v}\right\|_{1, \Omega^{\prime} \cup \Omega^{\prime \prime}} \leqq\left\|\mathbf{v}_{\gamma I}-\mathbf{v}_{\gamma}\right\|_{1, \Omega^{\prime} \cup \Omega^{\prime \prime}}+\left\|\mathbf{v}_{\gamma}-\mathbf{v}\right\|_{1, \Omega^{\prime} \cup \Omega^{\prime \prime}} \rightarrow 0 \text { for } \gamma, h \rightarrow 0 .
$$

If $\tau_{h} \in S_{F, P, h}^{+}$, we have

$$
\left\langle\tau_{h}, e\left(\boldsymbol{v}_{\gamma I}\right)\right\rangle \geqq L\left(\boldsymbol{v}_{\gamma l}\right) .
$$

Since $\tau_{h} \rightarrow \tau$ (weakly) in $S$ and $e\left(\mathbf{v}_{y I}\right) \rightarrow e(\mathbf{v})$ (strongly) in $S$, passing to the limit for $h, \gamma \rightarrow 0$ we conclude

$$
\langle\tau, e(\mathbf{v})\rangle \geqq L(\mathbf{v}) \quad \forall \mathbf{v} \in K_{0}
$$

i.e., $\tau \in S_{F, P}^{+}$.
Q.E.D.
3. AN ALGORITHM FOR THE SOLUTION OF THE APPROXIMATE PROBLEM

As in [2] we can simplify the approximate problem (2.1) by eliminating the auxiliary functions $\boldsymbol{v}_{h}$, which appear in the definition of $S_{F, P, h}^{+}$.

We set

$$
\mathbf{v}_{h}(\xi, \eta)=\sum_{i=1}^{N} q_{i} \varphi_{i}(\xi, \eta),
$$

where $q_{i}$ are the values of the displacement components at the nodes of the triangulation $\mathscr{T}_{h}$.

If we describe the condition $v_{h \xi}^{\prime \prime}-v_{h \xi}^{\prime} \leqq 0$ at the nodes lying on $\Gamma_{K}$, then precisely two components $\left\{q_{k_{1}}, q_{k_{2}}\right\}$ intervene at any pair of the nodes $\left(A_{k}, B_{k}\right) \in \bar{\Gamma}_{K}$ (see Fig. 2). Namely, the condition $v_{h \xi}^{\prime \prime}-v_{h \xi}^{\prime} \leqq 0$ implies $q_{k_{1}}-q_{k_{2}} \leqq 0$. We now introduce that linear transformation

$$
\boldsymbol{q}=F_{K} \mathbf{Y}: R^{2} \rightarrow R^{2}
$$

defined by the relations


It is readily seen that $F_{K}$ is regular. We consider the same transformation for any pair $M_{K}=\left\{q_{k_{1}}, q_{k_{2}}\right\}, k=1, \ldots, Q$ corresponding to the pair of nodes $\left(A_{k}, B_{k}\right) \in \bar{\Gamma}_{\kappa}$. We also set $y_{p}=q_{p}$ for $q_{p} \not \underset{k=1}{Q} M_{k}, 1 \leqq p \leqq N$. Finally, we have

$$
\begin{gather*}
\boldsymbol{q}=F \mathbf{y}: R^{N} \rightarrow R^{N}, \quad \text { and }  \tag{3.1}\\
\mathbf{v}_{h} \in K_{0 h} \Leftrightarrow \mathbf{q} \in \mathscr{K}_{q} \Leftrightarrow \mathbf{y} \in \mathscr{K}_{y}=\left\{\boldsymbol{y} \in R^{N}, y_{k_{2}} \leqq 0, k=1, \ldots, Q\right\} .
\end{gather*}
$$

Let $\psi_{T^{*}}$ be the characteristic function of the figure $T^{*} \in \mathscr{T}_{h}, T^{*}=T, T_{0}, T_{s}, T_{c}$. Then we have

$$
\begin{equation*}
\tau_{h} \in S_{h} \Leftrightarrow \tau_{h}(\xi, \eta)=\sum_{T^{*} \in \mathscr{F}_{h}} \tau\left(T^{*}\right) \psi_{\tau}(\xi, \eta) . \tag{3.2}
\end{equation*}
$$

Denoting by

$$
\begin{equation*}
\boldsymbol{t}^{\mathrm{T}}=\left\{\tau_{11}\left(T_{1}\right), \tau_{22}\left(T_{1}\right), \tau_{12}\left(T_{1}\right), \tau_{11}\left(T_{2}\right), \tau_{22}\left(T_{2}\right), \tau_{12}\left(T_{2}\right), \ldots\right\} \tag{3.3}
\end{equation*}
$$

the corresponding vector in $R^{M}$, we obtain

$$
\left\langle\tau, e\left(\mathbf{v}_{h}\right)\right\rangle=\sum_{T^{*} \in \mathscr{F}_{h}} \int_{T^{*}} \tau_{k j}\left(T^{*}\right) \sum_{i=1}^{N} q_{i} e_{k j}\left(\varphi_{i}\right) \mathrm{d} \xi \mathrm{~d} \eta=(E \mathbf{t}, \boldsymbol{q}),
$$

where $E$ is a matrix of order $N \times M(N<M)$ and $(E \boldsymbol{t}, \boldsymbol{q})=\boldsymbol{q}^{\mathrm{T}} E \boldsymbol{t}$. Since

$$
\left.L^{( }\left(\boldsymbol{v}_{h}\right)=\sum_{i=1}^{N} q_{i} L^{( } \varphi_{i}\right)=(\mathbf{I}, \boldsymbol{q}),
$$

where $\boldsymbol{I}$ is a fixed vector in $R^{N}$, the condition $\tau \in S_{F, P, h}^{+}$can be written in the form

$$
(\boldsymbol{I}-E \boldsymbol{t}, \boldsymbol{q}) \leqq 0 \quad \forall \boldsymbol{q} \in \mathscr{K}_{q} .
$$

Applying the mapping $F$, we obtain an equivalent condition

$$
\begin{equation*}
(I-E t, F \mathbf{y}) \leqq 0 \quad \forall \mathbf{y} \in \mathscr{K}_{y} \tag{3.4}
\end{equation*}
$$

Denote by $J^{-}$the set of all indexes $k_{2}, k=1, \ldots, Q$, and $J^{0}=\{1, \ldots, N\} \neg J^{-}$. Since the cone $\mathscr{K}_{y}$ is defined by the vectors

$$
\left\{\mathbf{e}_{j},-\mathbf{e}_{m} ; j \in J^{0}, m \in J^{-}\right\}
$$

where $\left\{\mathbf{e}_{j}, \mathbf{e}_{m}\right\}$ create an orthonormal basis of $R^{N},(3.4)$ is equivalent to the system

$$
\begin{align*}
& g_{j}(\mathbf{t})=\left(\boldsymbol{I}-E \mathbf{t}, F \mathbf{e}_{j}\right)=0, \quad j \in J^{0},  \tag{3.5}\\
& g_{m}(\mathbf{t})=\left(I-E t, F \mathbf{e}_{m}\right) \leqq 0, \quad m \in J^{-} \tag{3.6}
\end{align*}
$$

Inserting (3.2), (3.3) into the functional $\mathscr{S}\left(\tau_{h}\right)$, we arrive at the following problem: $\mathscr{S}_{0}(\boldsymbol{t})=$ min over the set of all $\boldsymbol{t}$ satisfying the conditions (3.5) and (3.6).

Remark 3.1. When choosing global coordinates $\left(x_{1}, x_{2}\right)$ instead of $(\xi, \eta)$ the procedure is analogous, but the situation is somewhat more complicated. In fact, if we describe the condition $v_{h \xi}^{\prime \prime}-v_{h \xi}^{\prime} \leqq 0$ at the nodes on $\Gamma_{K}$, there exist precisely four components $\left\{q_{k_{1}}, q_{k_{2}}, q_{k_{3}}, q_{k_{4}}\right\}$ at any pair of nodes $\left(A_{k}, B_{k}\right) \in \bar{\Gamma}_{K}$. The condition $v_{h \xi}^{\prime \prime}-v_{h \xi}^{\prime} \leqq 0$ implies $\sum_{j=1}^{4} b_{j} q_{h_{j}} \leqq 0$, where $b_{1}=-b_{3}=-\cos \left(x_{1}, \xi\right), b_{2}=-b_{4}=$
$=\cos \left(x_{2}, \xi\right)$.

## 4. ANOTHER APPROXIMATION BY FINITE ELEMENT METHOD

In Section 2 the dual problem (1.27) was approximated by piecewise constant external approximations using the triangulation suggested by Haslinger and Hlaváček [2]. It is the aim of this section to approximate the same problem using another triangulation, analogous to that suggested by Hlaváček and Křižek in [4].

Assume that the parts $\Gamma_{K}^{M}, M={ }^{\prime}$," consist of a finite number of convex and concave arcs. Here an arc $\Gamma_{0} \subset \partial \Omega$ is called convex (concave) if there exists a convex domain $\Omega_{0} \subset \Omega\left(\Omega_{0} \subset R^{2}-\Omega\right)$ such that $\Gamma_{0} \subset \partial \Omega_{0}$.

Next we define a triangulation of $\Omega=\Omega^{\prime} \cup \Omega^{\prime \prime}$. We construct domains $\Omega_{h}^{M}$, approximating $\Omega^{M}$ and such that $\bar{\Omega}_{h}^{M} \subset \bar{\Omega}^{M}, M={ }^{\prime}$, ". Each concave arc belonging to $\Gamma_{K}^{M}$ is approximated by a polygonal curve consisting of a finite number of line segments whose lengths are not greater than $h$. Each of these line segments is tangential to the concave arc. Furthermore, each inflexion point of $\Gamma_{K}^{M}$ must be a vertex of the
polygonal curve. By $\Gamma_{K h}^{M}$ we denote the union of the convex arcs and the approximate polygonal curves. The part of $\Omega^{M}$ bounded by $\Gamma_{K}^{M}$ and $\Gamma_{K h}^{M}$ will be denoted by $D_{h}^{M}$. Then the domain $\Omega_{h}^{M}=\Omega^{M} \doteq D_{h}^{M}$ will be triangulated in the standard way used in the finite element method. The triangles adjacent to the boundary may have a curved side along the boundary. We require each end point of $\Gamma_{K}^{M}$ and each node of the approximate polygonal curves to be also a node of the triangulation. Moreover, the corresponding nodes $A_{i}, B_{i}$ lying on $\Gamma_{K h}^{\prime}$ and $\Gamma_{K h}^{\prime \prime}$ must lie on the lines $A_{i} B_{i}$ parallel to $\xi$-axis (see Fig. 3). This triangulation will be denoted by $\mathscr{T}_{h}^{M}$ and $\mathscr{T}_{h}=$ $=\mathscr{T}_{h}^{\prime} \cup \mathscr{T}_{h}^{\prime \prime}$.


We say that the system $\left\{\mathscr{T}_{h}\right\}, 0<h \leqq h_{0}$ is regular if the smallest internal angle of all triangles of $\mathscr{T}_{h}$, for any $0<h \leqq h_{0}$, is not smaller than a constant $\theta>0$, independent of $h$. Here if a triangle is curved, the internal angles are defined by the angles of the "straight" triangle with the same vertices.

If a curved triangle $T \in \mathscr{T}_{h}$ adjacent to $\Gamma_{K}^{M}$ is convex, it is divided by a chord into a "straight" triangle $T_{0}$ and a segment $T_{s}$ such that $T=T_{0} \cup T_{s}$.

We introduce

$$
V_{h}=\left\{\mathbf{v} \in V|\mathbf{v}|_{T_{0}} \in\left[P_{1}\left(T_{0}\right)\right]^{2} \quad \forall T_{0} \subset T \in \mathscr{T}_{h} \text { adjacent to } \Gamma_{K}^{M}\right.
$$

$\left.\left(\frac{\partial}{\partial \xi} \mathbf{v}\right)\right|_{T_{s}}=0 \quad \forall T_{s} \subset T \in \mathscr{T}_{h}$ adjacent to $\Gamma_{K}^{M}$,
$\left.\left(\frac{\partial}{\partial \xi} \mathbf{v}\right)\right|_{T_{c}}=0 \quad \forall T_{c} \subset D_{h}, \quad$ where $T_{c}$ is a curved triangle defined by $H_{i}, \begin{array}{r}A_{i}, H_{i+1} \\ \text { (Fig. 3), }\end{array}$ $\left.\mathbf{v}\right|_{T} \in\left[P_{1}(T)\right]^{2}$ for all the remaining triangles $\}$.

In other words, $V_{h}$ consists of piecewise linear vector-functions which are extended continuously on segments $T_{s}$ and on non-convex curved triangles, constantly in the direction $\xi$.

Let $\varepsilon>0$ be fixed. We define $G^{M} \subset \Omega^{M}$ to be the so called $\varepsilon$-skin of the curved part of $\Gamma_{K}^{M}$, i.e.,

$$
G^{M}=\left\{y \in \Omega^{M}, \exists x \in \Gamma_{K}^{M} \dot{-} \Gamma_{K}^{* M}: \operatorname{dist}(x, y)<\varepsilon\right\}, \quad M==^{\prime}, \prime
$$

where $\Gamma_{K}^{* M}=\left\{x \in \Gamma_{K}^{M}\right.$; there exists a line segment $\left.S \subset \Gamma_{K}^{M}: x \in S\right\}$, and $G=G^{\prime} \cup G^{\prime \prime}$ (Fig. 4).


Fig. 4
The following assertion holds.
Lemma 4.1. Let $\quad \mathbf{v} \in\left[H^{2}\left(\Omega^{\prime}\right) \cap W^{1, \infty}\left(G^{\prime}\right)\right]^{2} \times\left[H^{2}\left(\Omega^{\prime \prime}\right) \cap W^{1, \infty}\left(\Omega^{\prime \prime}\right)\right]^{2}, f^{M} \in$ $\in C^{2}(\langle a, b\rangle)$. We define a Lagrange linear interpolation $\mathbf{v}_{I} \in V_{h}$ as follows: if $A_{i}$ is a vertex of a side tangential to $\Gamma_{K}^{M}, \mathbf{v}_{1}\left(A_{i}\right)=\mathbf{v}\left(D_{i}\right)$, where $D_{i} \in \Gamma_{K}^{M}$ is the projection of $A_{i}$ in the direction $\xi ; \mathbf{v}_{I}\left(a_{i}\right)=\mathbf{v}\left(a_{i}\right)$ at the other nodes; then we construct $\mathbf{v}_{I}$ so that $\mathbf{v}_{I} \in V_{h}$. Then

$$
\left\|\mathbf{v}_{I}-\boldsymbol{v}\right\|_{1, \Omega^{\prime} \cup \Omega^{\prime \prime}} \rightarrow 0 \quad \text { for } \quad h \rightarrow 0
$$

holds for any regular system of triangulations.
Proof. Let $V_{h}^{0}$ be the space of all piecewise linear functions on $\mathscr{T}_{h} \cup D_{h}$, continuous on $\bar{\Omega}^{M}, M={ }^{\prime}$, ", where each convex curved triangle $T=T_{0} \cup T_{s}$ will remain undivided while each non-convex curved triangle $T_{c} \subset D_{h}$ is divided by the line segment $A_{i} D_{i}$ into two triangles $T_{c}^{m}, m=1,2$, on which the functions from $V_{h}^{0}$ are linear and defined by the values at the nodes $A_{i}, H_{i}, H_{i+1}$ (see Fig. 3). Let the functions $\boldsymbol{w}_{h}^{0} \in V_{h}^{0}$ be constant in the direction $\xi$, i.e. $\left.\left((\partial / \partial \xi) \boldsymbol{w}_{h}^{0}\right)\right|_{T_{c^{m}}}=0$ for each $\boldsymbol{w}_{h}^{0} \in V_{h}^{0}$, $m=1,2$. The interpolation $\mathbf{v}_{h}^{0} \in V_{h}^{0}$ of $\mathbf{v}$ is defined so that $\mathbf{v}_{h}^{0}=\mathbf{v}$ at all nodes of the triangulation $\mathscr{T}_{h}$, i.e., except the nodes $D_{i}$.

First of all, we shall show that

$$
\begin{equation*}
\left\|\mathbf{v}_{h}^{0}-\boldsymbol{v}\right\|_{1, \Omega^{\prime} \cup \Omega^{\prime \prime}} \rightarrow 0 \quad \text { for } \quad h \rightarrow 0 . \tag{4.1}
\end{equation*}
$$

In fact, there exists an extension $E \mathbf{v} \in\left[H^{2}\left(R^{2}\right)\right]^{2}$ of $\mathbf{v}$ such that

$$
\|E \mathbf{v}\|_{2, R^{2}} \leqq C\|\mathbf{v}\|_{2, \Omega^{M}}
$$

(see e.g. [5], Chapt. 2). For any convex curved triangle adjacent to $\Gamma_{K}^{M}$ we define $\widetilde{T}=\Delta \tilde{a}_{i} a_{j} \tilde{a}_{k}$ (i.e. twice extended $T_{0}$ ) (see Fig. 5). $\widetilde{T}=T$ if $T$ is a straight triangle. Let $\pi$ denote the linear interpolation on $\tilde{T}$ with the nodes $a_{i}, a_{j}, a_{k}$. Using the affine equivalence and the regularity of the system $\left\{\mathscr{T}_{h}\right\}$ we deduce

$$
\begin{equation*}
\|\pi E \mathbf{v}-E \mathbf{v}\|_{1, \tilde{T}} \leqq C h|E v|_{2, \tilde{T}}, \tag{4.3}
\end{equation*}
$$

where $C$ is independent of $h$ and $E \mathbf{v}$. Since $\pi E \mathbf{v}=\mathbf{v}_{h}^{0}$ on $T^{*}, \forall T^{*} \in \mathscr{T}_{h}$, making use of (4.2), (4.3) we arrive at

$$
\begin{gather*}
\left\|\mathbf{v}-\mathbf{v}_{h}^{0}\right\|_{1, \Omega_{h}{ }^{M}}^{2}=\sum_{T^{*} \in \mathscr{F}_{h}{ }^{M}}\left\|\mathbf{v}-\mathbf{v}_{h}^{0}\right\|_{1, T^{*}}^{2} \leqq \sum_{T^{*} \in \mathscr{F}_{h}{ }^{M}}\left\|\mathbf{v}-\mathbf{v}_{h}^{0}\right\|_{1, T^{*}}^{2} \leqq  \tag{4.4}\\
\leqq C h_{T^{*} \in \mathscr{F}_{h}{ }^{M}}\|E \mathbf{v}\|_{1, \tilde{T}^{*}}^{2} \leqq 2 C h^{2}\|E \mathbf{v}\|_{2, R^{2}}^{2} \leqq C_{1} h^{2}|\mathbf{v}|_{2, \Omega^{M}}^{2} .
\end{gather*}
$$



Fig. 5

Fig. 6


Furthermore, it is readily seen that for each point $(\xi, \eta) \in T_{c}$ we have

$$
\left|v_{h j}^{0}(\xi, \eta)\right| \leqq \max _{k=i-1, i, i+1}\left|v_{j}\left(A_{h}\right)\right| \leqq\|v\|_{0, \infty, G}, \quad j=1,2,
$$

because for $h$ small enough $T_{c} \subset G$ and $\boldsymbol{v} \in\left[W^{1, \infty}(G)\right]^{2}$. Consequently, we have

$$
\left|\mathbf{v}_{h}^{0}(\xi, \eta)-\mathbf{v}(\xi, \eta)\right| \leqq 2\|\mathbf{v}\|_{0, \infty, G} \quad \forall(\xi, \eta) \in T_{c}
$$

and

$$
\begin{equation*}
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}\right\|_{0, T_{c}}^{2} \leqq 4\|\mathbf{v}\|_{0, \infty, G}^{2} \operatorname{mes} T_{c} \tag{4.5}
\end{equation*}
$$

Let us consider $(\partial / \partial \eta) \mathbf{v}_{h}^{0}$ on $T_{c}^{1}$ (Fig. 6). It is readily seen that

$$
\frac{\partial}{\partial v} \mathbf{v}_{h}^{0}=\frac{\partial}{\partial \eta} \mathbf{v}_{h}^{0}\left(\cos \alpha^{M}\right)+\left(\sin \alpha^{M}\right) \frac{\partial}{\partial \xi} \mathbf{v}_{h}^{0}=\left(\cos \alpha^{M}\right) \frac{\partial}{\partial \eta} \mathbf{v}_{h}^{0}
$$

i.e.,

$$
\frac{\partial}{\partial \eta} \mathbf{v}_{h}^{0}=\left(\cos \alpha^{M}\right)^{-1} \frac{\partial}{\partial v} \mathbf{v}_{h}^{0} .
$$

But $\boldsymbol{v}_{h}^{0}$ is linear in the direction $v$, therefore we have

$$
\frac{\partial}{\partial v} \mathbf{v}_{h}^{0}=\frac{\mathbf{v}_{h}^{0}\left(A_{i+1}\right)-\mathbf{v}_{h}^{0}\left(A_{i}\right)}{\left|A_{i} A_{i+1}\right|}=\frac{\mathbf{v}\left(A_{i+1}\right)-\mathbf{v}\left(A_{i}\right)}{\left|A_{i} A_{i+1}\right|}=\frac{\partial}{\partial v} \mathbf{v}\left(N_{i}\right),
$$

where $N_{i}$ is a point lying between $A_{i}$ and $A_{i+1}$. Since

$$
\left.\frac{\partial}{\partial v} \mathbf{v}_{( }^{\prime} N_{i}\right)=\frac{\partial}{\partial \eta} \mathbf{v}\left(N_{i}\right)\left(\sin \alpha^{M}\right)+\left(\cos \alpha^{M}\right) \frac{\partial}{\partial \xi} \boldsymbol{v}^{\prime}\left(N_{i}\right)
$$

we have

$$
\frac{\partial}{\partial \eta} \mathbf{v}_{h}^{0}=\frac{\partial}{\partial \eta} \mathbf{v}^{\prime}\left(N_{i}\right)+\left(\operatorname{tg} \alpha^{M}\right) \frac{\partial}{\partial \xi} \mathbf{v}^{\prime}\left(N_{i}\right)
$$

i.e.,

$$
\left|\frac{\partial}{\partial \eta} \mathbf{v}_{h}^{0}\right| \leqq C\|\mathbf{v}\|_{1, \infty, G}
$$

Here $C$ is independent of $T_{c}$ because of the fact that

$$
\left|\operatorname{tg} \alpha^{M}\right| \leqq\left\|f^{M}\right\|_{C^{1}(\langle a, b\rangle)} .
$$

A similar result holds for $T_{c}^{2}$.
Finally, we obtain

$$
\begin{aligned}
& \left|\frac{\partial}{\partial \xi}\left(\mathbf{v}_{h}^{0}-\mathbf{v}\right)\right|=\left|\frac{\partial}{\partial \xi} \mathbf{v}\right| \leqq\|\mathbf{v}\|_{1, \infty, G}, \\
& \left|\frac{\partial}{\partial \eta}\left(\mathbf{v}_{h}^{0}-\mathbf{v}\right)\right| \leqq C\|\mathbf{v}\|_{1, \infty, G}
\end{aligned}
$$

Thus we deduce

$$
\begin{equation*}
\left|\mathbf{v}_{h}^{0}-\mathbf{v}\right|_{1, T_{c}}^{2} \leqq C\|\mathbf{v}\|_{1, \infty, G}^{2} \operatorname{mes} T_{c} . \tag{4.6}
\end{equation*}
$$

Combining (4.5) with (4.6) we arrive at

$$
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}\right\|_{1, T_{c}}^{2} \leqq C\|\boldsymbol{v}\|_{1, \infty, G}^{2} \text { mes } T_{c}
$$

and

$$
\begin{equation*}
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}\right\|_{1, D_{h}^{M}}^{2} \leqq \sum_{T_{c} \in D_{h} M}\left\|\mathbf{v}_{h}^{0}-\mathbf{v}\right\|_{1, T_{c}}^{2} \leqq C\|\mathbf{v}\|_{1, \infty, G}^{2} \text { mes } D_{h}^{M} \rightarrow 0 \tag{4.7}
\end{equation*}
$$

for $h \rightarrow 0$ since mes $D_{h}^{M} \rightarrow 0$ for $h \rightarrow 0$. Then (4.1) is a consequence of (4.4) and (4.7).

Next we show that

$$
\begin{equation*}
\left\|\mathbf{v}_{I}-\mathbf{v}_{h}^{0}\right\|_{1, \Omega^{M}} \rightarrow 0 \text { for } h \rightarrow 0 . \tag{4.8}
\end{equation*}
$$

In fact, it is clear that

$$
\operatorname{Supp}\left(\mathbf{v}_{I}-\mathbf{v}_{h}^{0}\right) \subset \cup T_{s} \cup D_{h} \cup E_{h},
$$

where $E_{h}$ is the set of all "straight" triangles having at least one vertex at the point of intersection of two tangents to $\Gamma_{K}^{M}$.

In every $T_{s}$ we may write

$$
v_{h j}^{0}(\xi, \eta)-v_{I j}(\xi, \eta)=v_{h j}^{0}(\xi, \eta)-v_{I j}(\xi(s), \eta)=[\xi-\xi(s)] \frac{\partial}{\partial \xi} v_{h j}^{0}
$$

for the $j$-th component, $j=1,2$, where $(\xi(s), \eta)$ is the projection of the point $(\xi, \eta)$ on the chord in the direction $\xi$. Since $f^{M} \in C^{2}(\langle a, b\rangle)$ we have $|\xi-\xi(s)| \leqq h$. Thus we deduce

$$
\begin{equation*}
\int_{T_{s}}\left(v_{h j}^{0}-v_{I j}\right)^{2} \mathrm{~d} \xi \mathrm{~d} \eta \leqq h^{2} \int_{T_{s}}\left(\frac{\partial}{\partial \xi} v_{h j}^{0}\right)^{2} \mathrm{~d} \xi \mathrm{~d} \eta . \tag{4.9}
\end{equation*}
$$

For $T_{s}$ we also have

$$
\begin{gather*}
\frac{\partial}{\partial \xi}\left(v_{h j}^{0}-v_{I j}\right)=\frac{\partial}{\partial \xi} v_{h j}^{0}, \quad \frac{\partial}{\partial \eta}\left(v_{h j}^{0}-v_{I j}\right)=-\frac{\partial}{\partial \xi} v_{h j}^{0} \frac{\partial}{\partial \eta} \xi(s),  \tag{4.10}\\
\left|\frac{\partial}{\partial \eta} \xi(s)\right|=|\operatorname{tg} \alpha| \leqq\left\|f^{M}\right\|_{C^{1}(\langle a, b\rangle)}, \quad j=1,2
\end{gather*}
$$

where $\alpha$ is the angle between $\eta$-axis and the chord. From (4.9), (4.10) we obtain

$$
\begin{equation*}
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{I}\right\|_{1, T_{s}}^{2} \leqq C\left\|\mathbf{v}_{h}^{0}\right\|_{1, T_{s}}^{2} . \tag{4.11}
\end{equation*}
$$

Let us consider a $T \in E_{h}$. Then $T \subset G$ for $h$ small enough. We know that $T$ has either the position of $T^{*}$ or $T^{* *}$ (see Fig. 7).


Fig. 7


Fig. 8

If $T$ has the position of $T^{*}$, then $v_{h j}^{0}\left(a_{i+1}\right)=v_{I j}\left(a_{i+1}\right)$, i.e.,

$$
\begin{equation*}
\left|v_{h j}^{0}-v_{I j}\right| \leqq \max _{k=i, i+1}\left|v_{h j}^{0}\left(A_{k}\right)-v_{1 j}\left(A_{k}\right)\right|, \quad j=1,2 . \tag{4.12}
\end{equation*}
$$

If $T$ has the position of $T^{* *}, v_{h j}^{0}\left(a_{k}\right)=v_{I j}\left(a_{k}\right), k=i, i+1$, and for each $(\xi, \eta) \in T^{* *}$ we have

$$
\begin{equation*}
\left|v_{h j}^{0}(\xi, \eta)-v_{1 j}(\xi, \eta)\right| \leqq\left|v_{h j}^{0}\left(A_{i}\right)-v_{I j}\left(A_{i}\right)\right|, \quad j=1,2 . \tag{4.13}
\end{equation*}
$$

Moreover, for the $j$-th component, $j=1,2$, we have

$$
\begin{equation*}
\left|v_{h j}^{0}\left(A_{i}\right)-v_{I j}\left(A_{i}\right)\right|=\left|v_{j}\left(A_{i}\right)-v_{j}\left(D_{i}\right)\right|=\left|\int_{A_{i}}^{D_{i}} \frac{\partial}{\partial \xi} v_{j} \mathrm{~d} \xi\right| \leqq \int_{A_{i}}^{D_{i}}\left|\frac{\partial}{\partial \xi} v_{j}\right| \mathrm{d} \xi \tag{4.14}
\end{equation*}
$$

If $h$ is sufficiently small, $\overline{A_{i} D_{i}} \subset G,\left|A_{i} D_{i}\right| \leqq C h_{T}^{2}$, where $C$ is independent of $A_{\boldsymbol{i}}$ and $h_{T}$. Furthermore, we have

$$
\left|\frac{\partial}{\partial \xi} v_{j}\right| \leqq|\mathbf{v}|_{1, \infty, G}, \quad(\xi, \eta) \in \overline{A_{i} D_{i}} .
$$

Thus, from (4.14) we get

$$
\begin{equation*}
\left|v_{h j}^{0}\left(A_{i}\right)-v_{1 j}^{\prime}\left(A_{i}\right)\right| \leqq C h_{T}^{2}|\mathbf{v}|_{1, \infty, G} \tag{4.15}
\end{equation*}
$$

Combining (4.12) $\dot{\sim}(4.15)$, for every $T \in E_{h}$ we obtain

$$
\begin{gather*}
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{\boldsymbol{I}}\right\|_{0, \infty, T} \leqq C h^{2}|\boldsymbol{v}|_{1, \infty, G}  \tag{4.16}\\
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{I}\right\|_{0, T}^{2} \leqq C h_{T}^{4}|\boldsymbol{v}|_{1, \infty, G} \operatorname{mes} T \tag{4.17}
\end{gather*}
$$

Since $\boldsymbol{v}_{h}^{0}$ and $\mathbf{v}_{I}$ are linear on $T$ and the system $\left\{\mathscr{T}_{h}\right\}$ is regular, setting $\boldsymbol{w}=\mathbf{v}_{h}^{0}-\mathbf{v}_{r}$ we have (see Fig. 8)

$$
\frac{\partial w_{j}}{\partial x_{1}}=\frac{w_{j}\left(0_{1}\right)-w_{j}\left(0_{2}\right)}{\left|0_{1} 0_{2}\right|} \leqq \frac{2\|\boldsymbol{w}\|_{0, \infty, T}}{2 \varrho_{T}} \leqq \frac{1}{\alpha h_{T}}\|\boldsymbol{w}\|_{0, \infty, T}, \quad j=1,2
$$

where $\alpha>0$ is the number from the definition of the regularity of $\left\{\mathscr{T}_{h}\right\}$. Making use of (4.16) we obtain

$$
\begin{equation*}
\int_{T}\left|\frac{\partial w_{j}}{\partial x_{1}}\right|^{2} \mathrm{~d} x \leqq \frac{1}{\alpha^{2}} h_{T}^{2}|\boldsymbol{v}|_{1, \infty, G} \text { mes } T, \quad j=1,2 \tag{4.18}
\end{equation*}
$$

and a similar estimate for $\left.\int_{T}\left|\partial w_{j}\right| \partial x_{2}\right|^{2} \mathrm{~d} x$.
Finally, (4.17), (4.18) imply

$$
\begin{equation*}
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{I}\right\|_{1, T}^{2} \leqq C h^{2}|\mathbf{v}|_{1, \infty, G}^{2} \operatorname{mes} T . \tag{4.19}
\end{equation*}
$$

Let us consider the last case, when $T_{c} \subset D_{h}$ is defined by sides $A_{i} H_{i}, A_{i} H_{i+1}$ and the $\operatorname{arc} \overparen{H}_{i} H_{i+1}$ (see Fig. 6). By definitions of $\boldsymbol{v}_{h}^{0}, \boldsymbol{v}_{I}$ we get

$$
v_{h j}^{0}\left(A_{i}\right)=v_{j}\left(A_{i}\right), \quad v_{I j}\left(A_{i}\right)=v_{j}\left(D_{i}\right),
$$

i.e.,

$$
\begin{equation*}
v_{h j}^{0}\left(A_{i}\right)-v_{I j}\left(A_{i}\right)=v_{j}\left(A_{i}\right)-v_{j}\left(D_{i}\right) . \tag{4.20}
\end{equation*}
$$

Furthermore, for each $(\xi, \eta) \in T_{c}$ we can show

$$
\begin{equation*}
\left|v_{h j}^{0}(\xi, \eta)-v_{I j}(\xi, \eta)\right| \leqq \max _{k=i-1, i, i+1}\left|v_{j}\left(A_{k}\right)-v_{j}\left(D_{k}\right)\right| \tag{4.21}
\end{equation*}
$$

because of the fact that

$$
\left|v_{h j}^{0}\left(H_{s}\right)-v_{I j}\left(H_{s}\right)\right| \leqq \max _{k=s, s-1}\left|v_{j}\left(A_{k}\right)-v_{j}\left(D_{k}\right)\right|, \quad s=i, i+1 ; \quad j=1,2
$$

and the functions $v_{h j}^{0}$ are linear on $T_{c}^{m}, m=1,2$ (see Fig. 6).
From (4.14), (4.15) and (4.21) we deduce

$$
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{I}\right\|_{0, T_{c}}^{2} \leqq c h^{4}|\mathbf{v}|_{1, \infty, G}^{2} \operatorname{mes} T_{c} .
$$

Moreoves, for each $T_{c}$ we have

$$
\frac{\partial}{\partial \xi}\left(\mathbf{v}_{h}^{0}-\mathbf{v}_{I}\right)=0, \quad\left|\frac{\partial}{\partial \eta} \mathbf{v}_{h}^{0}\right| \leqq C|\mathbf{v}|_{1, \infty, G}, \quad\left|\frac{\partial}{\partial \eta} \mathbf{v}_{1}\right| \leqq C|\mathbf{v}|_{1, \infty, G} .
$$

Consequently,

$$
\left|\mathbf{v}_{h}^{0}-\mathbf{v}_{1}\right|_{1, T_{c}}^{2} \leqq C \mid \mathbf{v}_{1, \infty, G}^{2} \operatorname{mes} T_{c}
$$

holds, i.e.,

$$
\begin{equation*}
\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{I}\right\|_{1, T_{c}}^{2} \leqq\left(C h^{4}+C_{1}\right)|\boldsymbol{v}|_{1, \infty, G}^{2} \text { mes } T_{c} . \tag{4.22}
\end{equation*}
$$

Using (4.11), (4.19) and (4.22) we arrive at

$$
\begin{gather*}
\quad\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{I}\right\|_{1, \Omega}^{2}=\sum_{T_{s}}\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{I}\right\|_{1, T_{s}}^{2}+\sum_{T_{t \in E_{h}}}\left\|\mathbf{v}_{h}^{0}-\mathbf{v}_{l}\right\|_{1, T}^{2}+  \tag{4.23}\\
+\sum_{T_{c} \in D_{h}}\left\|v_{h}^{0}-\mathbf{v}_{I}\right\|_{1, T_{c}}^{2} \leqq C\left\|\mathbf{v}_{h}^{0}\right\|_{1, \cup T_{s}}^{2}+C_{1}\|\boldsymbol{v}\|_{1, \alpha, G} \operatorname{mes}\left(E_{h} \cup D_{h}\right) \rightarrow 0, \quad h \rightarrow 0 .
\end{gather*}
$$

Here we have also used the fact that mes $\left(E_{h} \cup D_{h}\right) \rightarrow 0$ for $h \rightarrow 0$, (using (4.4)) $\operatorname{mes}\left(\cup T_{s}\right) \rightarrow 0$ for $h \rightarrow 0$ and $\left\|\mathbf{v}_{h}^{0}\right\|_{1, \cup T_{s}} \leqq\|\boldsymbol{v}\|_{1, \cup T_{s}}+\left\|\mathbf{v}_{h}^{0}-\boldsymbol{v}\right\|_{1, \cup T_{s}}$.

Combining (4.1), (4.23) with the triangle inequality we obtain the assertion of Lemma 4.1.
Q.E.D.

Theorem 2. Assume that $f^{M} \in C^{2}, M={ }^{\prime}$," in a neighbourhood of the interval $\langle a, b\rangle$. Let the assumptions on $\Gamma_{K}^{M}, \Gamma_{0}, \Gamma_{u}, \Gamma_{\tau}$ from Lemma 2.2 hold. Assume that $\mathbf{u}$ satisfies the conditions of Lemma 4.1 and $S_{F, P}^{+} \neq \emptyset$.

Then the approximate problem 2.1 possesses a unique solution and

$$
\left\|\lambda_{h}-\lambda^{*}\right\| \rightarrow 0, \quad h \rightarrow 0
$$

holds for any regular system of triangulations $\left\{\mathscr{T}_{h}\right\}$, where $\lambda^{*}$ is the solution of the dual problem (1.27).

Proof is parallel to that of Theorem 1, but instead of Lemma 2.3 we use Lemma 4.1.

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## Souhrn

## DUÁLNÍ ANALÝZA KONTAKTU DVOU PRUŽNÝCH TĚLES S PROMĚNNÝM ROZSAHEM KONTAKTU METODOU KONEČNÝCH PRVKU゚

Tran Van Bon

V práci je provedena duální analýza kontaktního problému dvou pružných těles s proměnným rozsahem kontaktu. Aproximace řešení metodou konečných prvkủ jsou definovány na dvou typech triangulací po částech konstantními poli napětí. Dokazuje se konvergence obou typủ aproximací.

## Резюме

## ДВОЙСТВЕННЫЙ АНАЛИЗ КОНТАКТНОЙ ЗАДАЧИ УПРУГИХ ТЕЛ С РАСШИРЯЮЩЕЙСЯ ЗОНОЙ КОНТАКТА

Tran Van Bon

Определяется двойственная вариационная формулировка контактной задачи двух упругих тел. Аппроксимации постросны при помощи двух типов триангуляции и по частям постоянных полей напряжений. В обоих случаях доказывается сходимость аппроксимаций.

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