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# SOME NATURAL OPERATORS ON VECTOR FIELDS 

Jiníí Tomás


#### Abstract

We determine all natural operators transforming vector fields on a manifold $M$ to vector fields on $T^{*} T_{1}^{2} M, \operatorname{dim} M \geq 2$, and all natural operators transforming vector fields on $M$ to functions on $T^{*} T T_{1}^{2} M, \operatorname{dim} M \geq 3$. We describe some relations between these two kinds of natural operators.


## 0. Preliminaries

We present a contribution to the theory of natural operators and we follow the basic terminology used in [6]. Our starting point was a paper by Kobak, [2], in which all natural operators $T \rightarrow T T^{*} T$ were determined. In Section 1 we find all natural operators $T \rightarrow T T^{*} T_{1}^{2}$, where $T_{1}^{2}$ denotes the bundle of (1,2)velocities. In Section 2 we solve a related problem of finding of all natural operators transforming vector fields into functions on $T^{*} T T_{1}^{2}$. Our approach is heavily based on the technique of Weil bundles, [6].

All natural bundles and operators are considered on $\mathcal{M} f_{m}$, the category of smooth $m$-dimensional manifolds and their local diffeomorphisms. Let $\mathcal{M} f$ be the category of smooth manifolds and smooth maps and $\mathcal{F M}$ be the category of fibered manifolds.

Let $A=E(k) / I$ be a Weil algebra, where $E(k)$ is the algebra of germs of smooth functions $\mathbb{R}^{k} \rightarrow \mathbb{R}$ at zero and $I$ is an ideal of finite codimension. We remind the covariant definition of the Weil bundle functor $T^{A}: \mathcal{M} f \rightarrow \mathcal{F} \mathcal{M}$, [6],[3]. Two maps $f, g: \mathbb{R}^{k} \rightarrow M$ satisfying $f(0)=g(0)=x$ are said to be $I$-equivalent, if for every germ $h: M \rightarrow \mathbb{R}$ at $x$ it holds $h \circ f-h \circ g \in I$. Classes of such an equivalence relation are called $A$-velocities and are denoted by $j^{A} f$. They are the elements of $T^{A} M$. For a smooth map $f: M \rightarrow N$ we define $T^{A} f: T^{A} M \rightarrow T^{B} M$ by $T^{A} f\left(j^{A} g\right)=j^{A}(f \circ g)$ for all $j^{A} g \in T^{A} M$.

Given two Weil algebras $A, B$, we denote by $\operatorname{Hom}(A, B)$ the set of all algebra homomorphisms. A classical result reads there is a bijection between the elements

[^0]of $\operatorname{Hom}(A, B)$ and natural transformations $T^{A} \rightarrow T^{B}$. We shall need the following form of the result. Let $A=E(k) / I, B=E(p) / J$ and $f:\left(\mathbb{R}^{k}, 0\right) \rightarrow\left(\mathbb{R}^{p}, 0\right)$ be a smooth map. Then $j^{A} f \in T_{0}^{A} \mathbb{R}^{p}$ is said to be a $B$-admissible $A$-velocity iff $j^{A}(g \circ f)=0_{A}$ for all $g \in J$. It can be easily seen, that if $j^{A} f \in T_{0}^{A} \mathbb{R}^{p}$ is a $B$ admissible $A$-velocity, then $j^{A}(g \circ f)$ depends only on $j^{B} g$ for every $g: \mathbb{R}^{p} \rightarrow M$. The main result of [3] is that every $B$-admissible $A$-velocity $X=j^{A} f$ defines a natural transformation $i^{X}: T^{B} M \rightarrow T^{A} M$ by $j^{B} g \mapsto j^{A}(g \circ f)$. Moreover, every natural transformation $T^{B} \rightarrow T^{A}$ is of this type. It is proved in [6], that all those results remain valid if we restrict ourselves to the category $\mathcal{M} f_{m}$.

The group $\operatorname{Aut}(A)$ of all algebra automorphisms is a closed subgroup in GL $(A)$, so it is a Lie subgroup. Every element $D$ of its Lie algebra $\mathcal{A} u t(A)$ is tangent to a one-parameter subgroup $d(t)$ and determines a vector field $D(M)$ tangent to $(d(t))_{M}$ in $t=0$ on $T^{A} M$. Thus we have an absolute natural operator $T \rightarrow T T^{A}$ such that $X \mapsto D(M)$ for every vector field $X$. This operator is denoted by $\operatorname{op}(D)$, [6] ,[5].

Furthermore, for every natural bundle $F$ we have the flow operator $\mathcal{F}$, defined by $\mathcal{F}(X)=\left.\frac{\partial}{\partial t}\right|_{0} F\left(F l_{t}^{X}\right)$.

According to [6],[5], we have the following action of $A$ on tangent vectors of $T^{A} M$. If $m: \mathbb{R} \times T M \rightarrow T M$ is the multiplication of the tangent vectors on $M$ by reals, applying the functor $T^{A}$ we obtain $T^{A} m: T^{A} \mathbb{R} \times T^{A} T M \rightarrow T^{A} T M$. Since $T^{A} T M=T^{A \otimes \mathbb{D}} M$ and $T^{A} \mathbb{R}=A$, where $\mathbb{D}$ is the algebra of dual numbers, we have constructed a map $A \times T T^{\boldsymbol{A}} M \rightarrow T T^{\boldsymbol{A}} M$. The coordinate expression of the action of $c \in A$ is $c\left(a_{1}, \ldots, a_{m}, b_{1}, \ldots, b_{m}\right)=\left(a_{1}, \ldots, a_{m}, c b_{1}, \ldots, c b_{m}\right)$ for all $a_{1}, \ldots, a_{m}, b_{1}, \ldots, b_{m} \in A$. This is a natural affinor [5] and we denote it by $a f_{M}(c): T T^{A} M \rightarrow T T^{A} M$.
Proposition 1 ([6]). All natural operators $T \rightarrow T T^{A}$ are of the form af(c)。 $\mathcal{T}^{A}+\operatorname{op}(D)$ for all $c \in A, D \in \mathcal{A} u t(A)$.

In the special case $A=\mathbb{R}[x] /\left\langle x^{r+1}\right\rangle=\mathbb{D}_{1}^{r}$ we have $T^{A} M=T_{1}^{r} M=J_{0}^{r}(\mathbb{R}, M)$. Using the standard coordinates $\left(x^{i}, y_{1}^{i}, \ldots, y_{r}^{i}, X^{i}, Y_{1}^{i}, \ldots, Y_{r}^{i}\right)$ on $T T_{1}^{r} M$, we find $a f_{M}\left(x+\left\langle x^{r+1}\right\rangle\right)\left(X^{i}, Y_{1}^{i}, \ldots, Y_{r}^{i}\right)=\left(0, X^{i}, Y_{1}^{i}, \ldots, Y_{r-1}^{i}\right)$. Let $Q_{M}$ denote $a f_{M}(x+$ $\left.\left\langle x^{r+1}\right\rangle\right)$.
Proposition 2 ([6]). All natural operators $T \rightarrow T T_{1}^{r}$ are linearly generated by $\mathcal{T}_{1}^{r}, Q \circ \mathcal{T}_{1}^{r}, \ldots, Q^{r} \circ \mathcal{T}_{1}^{r}, L, Q \circ L, \ldots, Q^{r-1} \circ L$, where $L$ is the generalized Liouville vector field having the coordinate form $X^{i}=0, Y_{s}^{i}=s y_{s}^{i}$.

## 1. Natural Operators Transforming Vector Fields to $T^{*} T_{1}^{2}$

According to Proposition 1 we have five generating natural operators $T \rightarrow T T_{1}^{2}$ and according to [2] we have two generating natural operators $T \rightarrow T T^{*}$, the flow operator $\mathcal{T}^{*}\left(x^{i}, p_{i}\right)=X^{i} \frac{\partial}{\partial x_{i}}-X_{i}^{j} p_{j} \frac{\partial}{\partial p_{i}}$ and the Liouville field $\mathcal{L}\left(x^{i}, p_{i}\right)=p_{i} \frac{\partial}{\partial p_{i}}$, where $\left(x^{i}, p_{i}\right)$ are the standard coordinates on $T^{*} M$.

Composing these two sets of generators we obtain the following natural operators $T \rightarrow T T^{*} T_{1}^{2}: A_{1}=\mathcal{T}^{*} \circ \mathcal{T}_{1}^{2}, A_{2}=\mathcal{T}^{*} \circ\left(Q \circ \mathcal{T}_{1}^{2}\right), A_{3}=\mathcal{T}^{*} \circ\left(Q^{2} \circ \mathcal{T}_{1}^{2}\right)$ and absolute operators $A_{4}=\mathcal{T}^{*} \circ L, A_{5}=\mathcal{T}^{*} \circ(Q \circ L), A_{6}=\mathcal{L}$.

Let the canonical coordinates $x^{i}$ on $\mathbb{R}^{m}$ induce the coordinates $y^{i}=\frac{\partial x^{i}}{\partial \tau}, z^{i}=$ $\frac{\partial^{2} x^{i}}{\partial \tau^{2}}$ on $T_{1}^{2} \mathbb{R}^{m}$, while the additional coordinates on $T^{*} T_{1}^{2} \mathbb{R}^{m}$ are defined by $p_{i} d x^{i}+$ $q_{i} d y^{i}+r_{i} d z^{i}$. Further, let $x^{i}$ induce the additional coordinates $\omega_{i}$ on $T^{*} \mathbb{R}^{m}$ and $u^{i}=\frac{\partial x^{i}}{\partial \tau}, \gamma_{i}=\frac{\partial \omega_{i}}{\partial \tau}, w^{i}=\frac{\partial^{2} x^{i}}{\partial \tau^{2}}, \delta_{i}=\frac{\partial^{2} \omega_{i}}{\partial \tau^{2}}$ on $T_{1}^{2} T^{*} \mathbb{R}^{m}$.

We have the natural equivalence $s: T_{1}^{2} T^{*} \rightarrow T^{*} T_{1}^{2}$ of Cantrijn et al [1]

$$
\begin{align*}
& \left(x^{i}, \omega_{i}, v^{i}, \gamma_{i}, w^{i}, \delta_{i}\right) \mapsto\left(x^{i}, y^{i}, z^{i}, p_{i}, q_{i}, r_{i}\right)  \tag{1}\\
& \quad y^{i}=v^{i}, z^{i}=w^{i}, p_{i}=\delta_{i}, q_{i}=2 \gamma_{i}, r_{i}=\omega_{i}
\end{align*}
$$

Thus we have two other natural operators: $A_{7}=T s\left(\left(Q \circ \mathcal{T}_{1}^{2}\right) \circ \mathcal{L} \circ s^{-1}\right)$ and $A_{8}=T s\left(\left(Q^{2} \circ \mathcal{T}_{1}^{2}\right) \circ \mathcal{L} \circ s^{-1}\right)$.

Then the coordinate expressions of our operators are

$$
\begin{aligned}
& A_{1}(X)=X^{i} \frac{\partial}{\partial x^{i}}+X_{j}^{i} y^{j} \frac{\partial}{\partial y^{i}}+\left(X_{j}^{i} z^{j}+X_{j k}^{i} y^{j} y^{k}\right) \frac{\partial}{\partial z^{i}}-\left(X_{i}^{j} p_{j}+X_{i k}^{j} y^{k} q_{j}+\right. \\
& \\
& \left.+X_{i k}^{j} z^{k} r_{j}+X_{i k l}^{j} y^{k} y^{l} r_{j}\right) \frac{\partial}{\partial p_{i}}-\left(X_{i}^{j} q_{j}+2 X_{i k}^{j} y^{k} r_{j}\right) \frac{\partial}{\partial q_{i}}-X_{i}^{j} r_{j} \frac{\partial}{\partial r_{i}} \\
& A_{2}(X)=X^{i} \frac{\partial}{\partial y^{i}}+2 X_{j}^{i} y^{j} \frac{\partial}{\partial z^{i}}-\left(X_{i}^{j} q_{j}+2 X_{i k}^{j} y^{k} r_{j}\right) \frac{\partial}{\partial p_{i}}-2 X_{i}^{j} r_{j} \frac{\partial}{\partial q_{i}} \\
& A_{3}(X)=2 X^{i} \frac{\partial}{\partial z^{i}}-2 X_{i}^{j} r_{j} \frac{\partial}{\partial p_{i}} \quad A_{4}=y^{i} \frac{\partial}{\partial y^{i}}+2 z^{i} \frac{\partial}{\partial z^{i}}-q_{i} \frac{\partial}{\partial q_{i}}-2 r_{i} \frac{\partial}{\partial r_{i}} \\
& A_{5}=2 y^{i} \frac{\partial}{\partial z^{i}}-2 r_{i} \frac{\partial}{\partial q_{i}} \\
& A_{6}=p_{i} \frac{\partial}{\partial p_{i}}+q_{i} \frac{\partial}{\partial q_{i}}+r_{i} \frac{\partial}{\partial r_{i}} \\
& A_{7}=2 r_{i} \frac{\partial}{\partial q_{i}}+q_{i} \frac{\partial}{\partial p_{i}}
\end{aligned} \quad A_{8}=2 r_{i} \frac{\partial}{\partial p_{i}} .
$$

Let $p_{M}: F M \rightarrow M$ be a natural bundle of order r. According to the general theory, [6], there is a bijective correspondence between natural operators $A_{M}: T \rightarrow$ $T F M$ and natural transformations $\mathcal{A}_{M}: J^{r} T M \times_{M} F M \rightarrow T F M$ over the identity of $F M$, which is given by $\mathcal{A}_{M}\left(j_{x}^{r} X, y\right)=A_{M} X(y), x=p_{M}(y)$. Furthermore, there is a bijection between these natural transformations and equivariant maps of the standard fibers in question. Since $T^{*} T_{1}^{2}$ is a natural bundle of order three, we are searching for equivariant maps $\left(J^{3} T\right)_{0} \mathbb{R}^{m} \times\left(T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m} \rightarrow\left(T T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m}$. Let the additional coordinates on $T T^{*} T_{1}^{2}$ be

$$
\begin{equation*}
W^{i}=d x^{i}, Y^{i}=d y^{i}, Z^{i}=d z^{i}, P_{i}=d p_{i}, Q_{i}=d q_{i}, R_{i}=d r_{i} \tag{2}
\end{equation*}
$$

We evaluate the necessary transformation laws of the action of $G_{m}^{4}$ on the standard fibers. Denote by $\left(a_{j_{1}}^{i}, \ldots, a_{j_{1} \ldots j_{r}}^{i}\right)$ the canonical coordinates on $G_{m}^{r}$ and indicate by tilde the coordinates of the inverse element. The action of $G_{m}^{4}$ on $\left(T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m}$ looks as follows

$$
\begin{equation*}
\bar{y}^{i}=a_{j}^{i} y^{j} \quad \bar{z}^{i}=a_{j}^{i} z^{j}+a_{j k}^{i} y^{j} y^{k} \quad \bar{r}_{i}=\tilde{a}_{i}^{j} r_{j} \quad \bar{q}_{i}=\tilde{a}_{i}^{j} q_{j}+2 \tilde{a}_{i k}^{j} a_{l}^{k} y^{l} r_{j} \tag{3}
\end{equation*}
$$

3) $\quad \bar{p}_{i}=\tilde{a}_{i}^{j} p_{j}+\tilde{a}_{i k}^{j} a_{l}^{k} y^{l} q_{j}+\tilde{a}_{i k}^{j} a_{l}^{k} z^{l} r_{j}+\tilde{a}_{i k}^{j} a_{l m}^{k} y^{l} y^{m} r_{j}+\tilde{a}_{i k l}^{j} a_{m}^{k} a_{n}^{l} y^{m} y^{n} r_{j}$

Let $B_{m}^{r+1}=\left\{j_{0}^{r+1} f ; j_{0}^{r} f=j_{0}^{r} i d_{\mathbb{R}_{m}}\right\}$. Then

$$
\begin{array}{ll}
\bar{q}_{i}=q_{i}-2 a_{i k}^{j} y^{k} r_{j} & \text { for the action of } B_{m}^{2} \\
\bar{p}_{i}=p_{i}-a_{i k l}^{j} y^{k} y^{l} r_{j} & \text { for the action of } B_{m}^{3} \tag{4}
\end{array}
$$

and

$$
\begin{array}{cc}
\bar{X}_{j_{1} \ldots j_{r}}^{i}=X^{i}+a_{j_{1} \ldots j_{r} k}^{i} X^{k} \quad \text { for the action of } B_{m}^{r+1}  \tag{5}\\
\bar{X}_{j_{1} j_{2}}^{i}=X_{j_{1} j_{2}}^{i}-a_{k l}^{i} a_{j_{1} j_{2}}^{l} X^{k}-a_{j_{1} j_{2}}^{k} X_{k}^{i}+a_{j_{1} k}^{i} X_{j_{2}}^{k}+a_{j_{2} k}^{i} X_{j_{1}}^{k} \\
& \text { for the action of } B_{m}^{2},
\end{array}
$$

where $X_{j_{1} \ldots j_{r}}^{i}$ indicates the $r$-jets of a vector field $X$. Furthermore

$$
\begin{array}{llll}
\bar{W}^{i}=a_{j}^{i} W^{j} & \bar{R}_{i}=\tilde{a}_{i}^{j} R_{j} & \text { and it holds } &  \tag{6}\\
\bar{Y}^{i}=a_{j}^{i} Y^{j} & \bar{Q}_{i}=\tilde{a}_{i}^{j} Q_{j} & \bar{Z}^{i}=a_{j}^{i} Z^{j} & \bar{P}_{i}=\tilde{a}_{i}^{j} R_{j}
\end{array}
$$

whenever all the previous coordinates are zeros. Moreover, only $P_{i}$ are changed by $B_{m}^{4}$ and it holds

$$
\begin{equation*}
\bar{P}_{i}=P_{i}-a_{i k l m}^{j} y^{k} y^{l} W^{m} r_{j} \tag{7}
\end{equation*}
$$

Finally we need the following lemma. Let

$$
V_{p, q}=\underbrace{V \times \ldots \times V}_{p-\text { times }} \times \overbrace{V^{*} \times \ldots \times V^{*}}^{q-\text { times }}
$$

where $V$ denotes the vector space $\mathbb{R}^{m}$ with the standard action of $G_{m}^{1}$.
Lemma 3 ([6]). (a) All smooth $G_{m}^{1}$-equivariant maps $V_{p, q} \rightarrow V$ are of the form

$$
\sum_{j=1}^{p} g_{j}\left(\left\langle x_{k}, y_{l}\right\rangle\right) x_{j}
$$

where $g_{j}: \mathbb{R}^{p q} \rightarrow \mathbb{R}$ are any smooth functions, $j, k=1, \ldots, p, l=1, \ldots, q$.
(b) All smooth $G_{m}^{1}$-equivariant maps $V_{p, q} \rightarrow V^{*}$ are of the form

$$
\sum_{l=1}^{q} h_{l}\left(\left\langle x_{k}, y_{h}\right\rangle\right) y_{l}
$$

where $h_{l}: \mathbb{R}^{p q} \rightarrow \mathbb{R}$ are any smooth functions, $k=1, \ldots, p, h, l=1, \ldots, q$.
The proof of the main result essentially uses the following two lemmas.

Lemma 4. Let $h:\left(J^{3} T\right)_{0} \mathbb{R}^{m} \times\left(T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m} \rightarrow \mathbb{R}^{m}$ be an equivariant smooth mapping, $m \geq 2$. Then it holds

$$
\begin{equation*}
W^{i}=g_{1}\left(I_{1}, \ldots, I_{5}\right) X^{i}+g_{2}\left(I_{1}, \ldots, I_{5}\right) y^{i} \tag{8}
\end{equation*}
$$

where $g_{1}, g_{2}$ are any smooth functions $\mathbb{R}^{5} \rightarrow \mathbb{R}$ and $I_{1}, \ldots, I_{5}$ are invariants of the form

$$
\begin{array}{ll}
I_{1}=X^{i} p_{i}+X_{j}^{i} y^{j} q_{i}+\left(X_{j}^{i} z^{j}+X_{j k}^{i} y^{j} y^{k}\right) r_{i} & I_{2}=X^{i} q_{i}+2 X_{j}^{i} y^{j} r_{i}  \tag{9}\\
I_{3}=X^{i} r_{i} & I_{4}=y^{i} q_{i}+2 z^{i} r_{i}
\end{array} I_{5}=y^{i} r_{i} .
$$

Proof. The first formula from (5) implies, that $W^{i}=h^{i}\left(j_{0}^{3} X, y^{i}, z^{i}, p_{i}, q_{i}, r_{i}\right)$ does not depend on $X_{j_{1} j_{2} j_{3}}^{i}$. Therefore we are searching equivariant maps $\left(J^{2} T\right)_{0} \mathbb{R}^{m} \times$ $\left(T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m} \rightarrow \mathbb{R}^{m}$.

Let $S_{0}$ be $C_{0} \times\left(T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m}$, where $C_{0}$ is the set of all 2 -jets of constant vector fields on $\mathbb{R}^{m}$ at zero. Since $S_{0}$ is $G_{m}^{1}$-invariant and $W^{i}=a_{j}^{i} W^{j}$, the equivariance and Lemma 1 yield $W^{i}=\alpha_{1} X^{i}+\alpha_{2} y^{i}+\alpha_{3} z^{i}$ on $S_{0}$, where $\alpha_{1}, \alpha_{2}, \alpha_{3}$ are some functions of $X^{i} p_{i}, X^{i} q_{i}, X^{i} r_{i}, y^{i} p_{i}, y^{i} q_{i}, y^{i} r_{i}, z^{i} p_{i}, z^{i} q_{i}, z^{i} r_{i}$. Since $X^{i} p_{i}, X^{i} q_{i}, X^{i} r_{i}$ coincide with $I_{1}, I_{2}, I_{3}$ on $S_{0}, \alpha_{1}, \alpha_{2}, \alpha_{3}$ can be considered as functions of arguments $I_{1}, I_{2}, I_{3}, I_{4}, I_{5}$ and $y^{i} p_{i}, z^{i} p_{i}, z^{i} q_{i}, z^{i} r_{i}$.

Let $S_{1} \subseteq S_{0}$ be the subset of all elements of $S_{0}$ satisfying the following conditions: $X^{i}$ and $y^{i}$ as well as $X^{i}$ and $z^{i}$ as well as $y^{i}$ and $z^{i}$ are linearly independent vectors and $r_{i}$ is a non-zero vector. Obviously, $S_{1}$ is a dense subset of $S_{0}$. Let $i: G_{m}^{1} \rightarrow G_{m}^{3}$ be the canonical injection. Fixing $X^{i}, y^{i}, z^{i}, p_{i}, q_{i}, r_{i}$ we can find some $j_{0}^{3} f \in i\left(G_{m}^{1}\right)$ transforming $X^{i}$ to $\delta_{1}^{i}$, $z^{i}$ to $\delta_{2}^{i}$, while the other values are transformed to the bared ones. This is possible on $S_{1}$ due to the conditions from its definition.

Let $\ell$ denote, in general, the left action of the $r$-th order differential group on the standard fiber of an $r$-th order natural bundle. We have $h^{i}\left(j_{0}^{2} X, y^{i}, z^{i}, p_{i}, q_{i}, r_{i}\right)=$ $\ell\left(j_{0}^{3} f^{-1}, \ell\left(j_{0}^{3} f, h^{i}\left(j_{0}^{2} X, y^{i}, z^{i}, p_{i}, q_{i}, r_{i}\right)\right)\right)=\ell\left(j_{0}^{3} f^{-1}, \alpha_{1} \delta_{1}^{i}+\alpha_{2} \bar{y}^{i}+\alpha_{3} \delta_{2}^{i}\right)$, where the arguments of $\alpha_{1}, \alpha_{2}, \alpha_{3}$ are $I_{1}, \ldots I_{5}$ and $\bar{y}^{i} \bar{p}_{i}, \bar{z}^{i} \bar{p}_{i}, \bar{z}^{i} \bar{q}_{i}, \bar{z}^{i} \bar{r}_{i}$ satisfying $\bar{z}^{i}=\delta_{2}^{i}$. It follows from the equivariance of $h$ and the fact, that the last four arguments of $\alpha_{1}, \alpha_{2}, \alpha_{3}$ are $G_{m}^{1}$-invariants.

The definition of $S_{1}$ implies, that there is $j_{0} \geq 2$ such that $y^{j_{0}} \neq 0$. Let $i_{1}: B_{m}^{2} \rightarrow G_{m}^{3}$ denote the canonical inclusion. Taking $j_{0}^{3} f_{1} \in i_{1}\left(B_{m}^{2}\right)$ with all $a_{j k}^{i}=0$ except $a_{j_{0} j_{0}}^{2}$ we can annihilate all expresions with $z_{i}$. It follows from (5) that $j_{0}^{3} f_{1}$ stabilizes $j_{0}^{2}\left(\frac{\partial}{\partial x^{1}}\right)$. But we changed the value of $\bar{y}^{i} \bar{p}_{i}$, which can be annihilated by taking a suitable $j_{0}^{3} f_{2} \in B_{m}^{3}$ with all $a_{j k l}^{i}=0$ except $a_{j_{0} j_{0} j_{0}}^{k_{0}}$, where $k_{0}$ is an index such that $\bar{r}_{k_{0}} \neq 0$. It follows directly from (4) and (5), that $j_{0}^{3} f_{2}$ stabilizes $j_{0}^{2} \frac{\partial}{\partial x^{1}}$.

Thus we obtain, that $W^{i}=\ell\left(j_{0}^{3} f^{-1}, \alpha_{1} \delta_{1}^{i}+\alpha_{2} \bar{y}^{i}\right)$ on $S_{1}$, where the last four arguments of $\alpha_{1}, \alpha_{2}$ are zeros, while the invariants are not changed. So we have

$$
\begin{aligned}
W^{i} & =\ell\left(j_{0}^{3} f^{-1}, \alpha_{1}\left(I_{1}, \ldots, I_{5}, 0,0,0,0\right) \delta_{1}^{i}+\alpha_{2}\left(I_{1}, \ldots, I_{5}, 0,0,0,0\right) \bar{y}^{i}\right)= \\
& =\alpha_{1}\left(I_{1}, \ldots, I_{5}, 0,0,0,0\right) X^{i}+\alpha_{2}\left(I_{1}, \ldots, I_{5}, 0,0,0,0\right) y^{i}
\end{aligned}
$$

which follows from the equivariance of the map $h$. Substituting $g_{i}\left(I_{1}, \ldots, I_{5}\right)$ for $\alpha_{i}\left(I_{1}, \ldots, I_{5}, 0,0,0,0\right), i=1,2$, we have

$$
\begin{equation*}
W^{i}=g_{1}\left(I_{1}, \ldots, I_{5}\right) X^{i}+g_{2}\left(I_{1}, \ldots, I_{5}\right) y^{i} \text { on } S_{1} \tag{10}
\end{equation*}
$$

Since $S_{1}$ is dense in $S_{0}$, this holds on $S_{0}$ as well. Taking into account the equivariance of $h$, (10) can be extended to $\left(J^{2} T\right)_{0} \mathbb{R}^{m} \times\left(T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m}$, which completes the proof.

The following lemma is the dualization of Lemma 4 and since its proof is almost the same as that of Lemma 4, we omit it.

Lemma 5. Let $h:\left(J^{3} T\right)_{0} \mathbb{R}^{m} \times\left(T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m} \rightarrow \mathbb{R}^{m *}$ be an equivariant smooth mapping, $m \geq 2$. Then

$$
\begin{equation*}
R_{i}=g_{1}\left(I_{1}, \ldots, I_{5}\right) r_{i} \tag{11}
\end{equation*}
$$

where $g: \mathbb{R}^{5} \rightarrow \mathbb{R}$ is a smooth function.

Proposition 6. For $m \geq 2$, every natural operator $A: T \rightarrow T T^{*} T_{1}^{2}$ is of the form $A=\sum_{j=1}^{8} h_{j}\left(I_{1}, \ldots, I_{5}\right) \bar{A}_{j}$, where $h_{j}: \mathbb{R}^{5} \rightarrow \mathbb{R}$ are some smooth functions and

$$
\begin{array}{lll}
A_{1}=\mathcal{T}^{*} \circ \mathcal{T}_{1}^{2} & A_{2}=\mathcal{T}^{*} \circ\left(Q \circ \mathcal{T}_{1}^{2}\right) & A_{3}=\mathcal{T}^{*} \circ\left(Q^{2} \circ \mathcal{T}_{1}^{2}\right) \\
A_{4}=\mathcal{T}^{*} \circ L & A_{5}=\mathcal{T}^{*} \circ(Q \circ L) & A_{6}=\mathcal{L}  \tag{12}\\
A_{7}=T s\left(\left(Q \circ \mathcal{T}_{1}^{2}\right) \circ \mathcal{L} \circ s^{-1}\right) & A_{8}=T s\left(\left(Q^{2} \circ \mathcal{T}_{1}^{2}\right) \circ \mathcal{L} \circ s^{-1}\right)
\end{array}
$$

Proof. In the whole proof we use the coordinates (2). Let $A: T \rightarrow T T^{*} T_{1}^{2}$ be a natural operator and $h$ be the corresponding equivariant map. Since $\bar{W}^{i}=a_{j}^{i} W^{j}$, applying Lemma 4 we get $W^{i}=g_{1}\left(I_{1}, \ldots, I_{5}\right) X^{i}+g_{2}\left(I_{1}, \ldots, I_{5}\right) y^{i}$. Taking the natural operator $B_{1}=A-g_{1}\left(I_{1}, \ldots, I_{5}\right) \mathcal{T}^{*} \circ \mathcal{T}_{1}^{2}$ we get its equivariant map in the form $W^{i}=g_{2}\left(I_{1}, \ldots, I_{5}\right) y^{i}$.

First of all we prove, that $g_{2}$ is the zero function. Let $\alpha=\left(j_{0}^{3}\left(\frac{\partial}{\partial x^{1}}\right), \delta_{2}^{i}, z^{i}, p_{i}, q_{i}, r_{i}\right)$ be an element of $\left(J^{3} T\right)_{0} \mathbb{R}^{m} \times\left(T^{*} T_{1}^{2}\right)_{0} \mathbb{R}^{m}$ satisfying the existence of a non-zero $r_{i}$. Let $j_{0}$ be the least index, for which $r_{j_{0}} \neq 0$, and let $j_{0}^{4} f \in B_{m}^{4}$ satisfy $a_{j k l m}^{i}=0$ except $a_{2222}^{j_{0}}$. Then the formula (7) implies, that we can change the value of $P_{2}$ stabilizing $\alpha$, whenever $g_{2}\left(p_{1}, q_{1}, r_{1}, q_{2}+2 z^{i} r_{i}, r_{2}\right) \neq 0$. Thus we obtain, that $g_{2}$ is the zero function on $\mathbb{R}^{5}$.

Now, put $h_{1}=g_{1}$ and consider the natural operator $B_{1}$. Since its equivariant map satisfies $W^{i}=0$, the formula (6) and Lemma 4 yield $Y^{i}=g_{3}\left(I_{1}, \ldots, I_{5}\right) X^{i}+$ $g_{4}\left(I_{1}, \ldots I_{5}\right) y^{i}$. We can subtract $g_{3}\left(I_{1}, \ldots, I_{5}\right) A_{2}+g_{4}\left(I_{1}, \ldots I_{5}\right) A_{4}$ and write $h_{2}=$ $g_{3}$ and $h_{4}=g_{4}$. We can iterate these steps using the formula (6), Lemmas 4 and 5 . This way we prove our claim.

## 2. Natural Operators $T \rightarrow C^{\infty}\left(T^{*} T T_{1}^{2}, \mathbb{R}\right)$

In this part we are searching all natural operators transforming vector fields to functions on $T^{*} T T_{1}^{2}$. We use essentially the following result by Kolář, [4]. Let $F$ be a natural bundle, $Y: F M \rightarrow T F M$ be a vector field and $\widetilde{Y}$ denote the function $T^{*} F M \rightarrow \mathbb{R}$ defined by $\tilde{Y}(w)=\langle Y(p(w))$, $w\rangle$, where $p$ is the cotangent bundle projection. Let $F$ have the following properties I,II,III.
I. The set $N_{F}$ of all natural operators $T \rightarrow T F M$ is a finite dimensional vector space. (This property is satisfied for every Weil bundle.)

Let $N_{F}^{*}$ be the dual vector space and $\operatorname{Nop}\left(T, T^{*} F \times \mathbb{R}\right)$ denote the set of all natural operators $T \rightarrow C^{\infty}\left(T^{*} F, \mathbb{R}\right)$. For every smooth function $h: N_{F}^{*} \rightarrow \mathbb{R}$ Kolár̆ constructed the following natural operator $\mathrm{Dh}: T \rightarrow C^{\infty}\left(T^{*} F, \mathbb{R}\right)$. Fixing a basis $A_{1}, \ldots, A_{n}$ of $N_{F}$, its dual vector space $N_{F}^{*}$ can be identified with $\mathbb{R}^{n}$ and we can put $(\mathrm{Dh})_{M} X=h\left(\widehat{A_{1, M} X}, \ldots, \widehat{A_{n, M} X}\right): T^{*} F M \rightarrow \mathbb{R}$. Thus we obtain a mapping $C^{\infty}\left(N_{F}^{*}, \mathbb{R}\right) \rightarrow \operatorname{Nop}\left(T, T^{*} F \times \mathbb{R}\right)$.
II. There exists a smooth function $j: N_{F}^{*} \rightarrow\left(T^{*} F\right)_{0} \mathbb{R}^{m}$ satisfying

$$
\begin{equation*}
\langle A, u\rangle=\widehat{A\left(\frac{\partial}{\partial x^{1}}\right)(j u)} \tag{13}
\end{equation*}
$$

for every $A \in N_{F}, u \in N_{F}^{*}$.
Let Diff $f_{0}^{1} \mathbb{R}^{m}$ denote the stability group of the origin and the vector field $\frac{\partial}{\partial x^{1}}$.
III. The orbit of $j\left(N_{F}^{*}\right)$ with respect to $D i f_{0}^{1} \mathbb{R}^{m}$ is dense in $\left(T^{*} F\right)_{0} \mathbb{R}^{m}$.

Proposition 7 ([4]). If the assumptions I, II, III are satisfied, then all natural operators $T \rightarrow C^{\infty}\left(T^{*} F, \mathbb{R}\right)$ are of the form Dh for all $h \in C^{\infty}\left(N_{F}^{*}, \mathbb{R}\right)$.

This result enables searching for natural operators $T \rightarrow C^{\infty}\left(T^{*} F, \mathbb{R}\right)$, where $F=T^{A}$ is a Weil bundle. Let $T^{A}$ be of order $r$. In order to find all the natural operators $T \rightarrow T T^{*} T^{A}$ we can use the following procedure consisting of four steps.
(a) We find a base $B_{1}, \ldots B_{k}$ of all natural operators $T \rightarrow T T^{A}$.
(b) We take some immersion element $i \in T_{0}^{A} \mathbb{R}^{m}$. Over the element $i$ we have a space $P$ in $\left(T^{*} T^{\boldsymbol{A}}\right)_{0} \mathbb{R}^{m}$, on which the stabilizing group $H$ of $i$ and $j_{0}^{r}\left(\frac{\partial}{\partial x^{1}}\right)$ acts.
(c) We compute $\left.I_{i}=\widetilde{B_{i}}\left(\frac{\partial}{\partial x^{1}}\right) \right\rvert\, P$. If possible, we choose coordinates $w_{1}, \ldots, w_{k}$, $z_{1}, \ldots, z_{l}$ on $P$ such that $w_{i}=I_{i}$.
(d) We prove, that we can annihilate $z_{1}, \ldots, z_{l}$ on a dense subset of $P$ by the group $H$.

Then every natural operator $T \rightarrow C^{\infty}\left(T^{*} F, \mathbb{R}\right)$ is smoothly generated by $\widetilde{B_{1}}, \ldots, \widetilde{B_{k}}$. Indeed we can define

$$
\begin{equation*}
j: N_{F}^{*} \rightarrow\left(T^{*} F_{0}\right) \mathbb{R}^{m}, \quad b_{1} B^{1}+\cdots+b_{k} B^{k} \mapsto\left(b_{1}, \ldots b_{k}, 0, \ldots, 0\right) \tag{14}
\end{equation*}
$$

which clearly satisfies (13). The denseness of the orbit $j\left(N_{F}^{*}\right)$ is guaranteed by (d).

Now we use this procedure for the bundle $T T_{1}^{2}$. First of all we find all natural operators $T \rightarrow T T T_{1}^{2}$. Since $T T_{1}^{2}=T^{\mathbb{D} \otimes \mathbb{D}_{1}^{2}}$, where $\mathbb{D} \otimes \mathbb{D}_{1}^{2}=\mathbb{R}[t, \tau] /\left\langle t^{2}, \tau^{3}\right\rangle$, every element from $T T_{1}^{2} M$ is of the form $x^{i}+z_{1}^{i} \tau+\frac{1}{2} z_{2}^{i} \tau^{2}+y^{i} t+w_{1}^{i} t \tau+\frac{1}{2} w_{2}^{i} t \tau^{2}$, where $\left(x^{i}, z_{1}^{i}, z_{2}^{i}, y^{i}, w_{1}^{i}, w_{2}^{i}\right)$ are the canonical coordinates on $T T_{1}^{2} M$.

Lemma 8. All natural operators $T \rightarrow T T T_{1}^{2}$ are linearly generated by the following ones

$$
\begin{array}{ll}
N_{1}=\mathcal{T} \circ \mathcal{T}_{1}^{2} & N_{2}=a f\left(\tau+\left\langle t^{2}, \tau^{3}\right\rangle\right)\left(\mathcal{T} \circ \mathcal{T}_{1}^{2}\right) \\
N_{3}=a f\left(t+\left\langle t^{2}, \tau^{3}\right\rangle\right)\left(\mathcal{T} \circ \mathcal{T}_{1}^{2}\right) & N_{4}=a f\left(\tau^{2}+\left\langle t^{2}, \tau^{3}\right\rangle\right)\left(\mathcal{T} \circ \mathcal{T}_{1}^{2}\right) \\
N_{5}=a f\left(t \tau+\left\langle t^{2}, \tau^{3}\right\rangle\right)\left(\mathcal{T} \circ \mathcal{T}_{1}^{2}\right) & N_{6}=a f\left(t \tau^{2}+\left\langle t^{2}, \tau^{3}\right\rangle\right)\left(\mathcal{T} \circ \mathcal{T}_{1}^{2}\right) \\
N_{7}=y^{i} \frac{\partial}{\partial y^{i}}+w_{1}^{i} \frac{\partial}{\partial w_{1}^{i}}+w_{2}^{i} \frac{\partial}{\partial w_{2}^{i}} & N_{8}=z_{1}^{i} \frac{\partial}{\partial z_{1}^{i}}+2 z_{2}^{i} \frac{\partial}{\partial z_{2}^{i}}+w_{1}^{i} \frac{\partial}{\partial w_{1}^{i}}+2 w_{2}^{i} \frac{\partial}{\partial w_{2}^{i}} \\
N_{9}=y^{i} \frac{\partial}{\partial w_{1}^{i}}+2 w_{1}^{i} \frac{\partial}{\partial w_{2}^{i}} & N_{10}=2 y^{i} \frac{\partial}{\partial w_{2}^{i}} \\
N_{11}=2 z_{1}^{i} \frac{\partial}{\partial z_{2}^{i}}+2 z_{2}^{i} \frac{\partial}{\partial w_{2}^{i}} & N_{12}=z_{1}^{i} \frac{\partial}{\partial w_{1}^{i}}+2 z_{2}^{i} \frac{\partial}{\partial w_{2}^{i}} \\
N_{13}=2 z_{1}^{i} \frac{\partial}{\partial w_{2}^{i}} &
\end{array}
$$

Proof. By Proposition 1 we have to determine the absolute operators. In our case $A=\mathbb{D} \otimes \mathbb{D}_{1}^{2}$. Every $A$-velocity in question is of the form

$$
\begin{align*}
& a t+b \tau+c \tau^{2}+d t \tau+e t \tau^{2}  \tag{15}\\
& f t+g \tau+h \tau^{2}+j t \tau+k t \tau^{2}
\end{align*}
$$

Taking into account the conditions of admissibility we obtain $b=0, a c=0$ and $3 f g^{2}=0$. Since every $A$-admissible $A$-velocity induces a homomorphism $A \rightarrow A$ and we are searching for curves in $\operatorname{Aut}(A)$ in a neighbourhood of the unit, we can restrict ourselves to the connected component of the unit in $\operatorname{Aut}(A)$. Then we have $c=0$ and $f=0$. Renaming the parameters in (15), all considered automorphisms $A \rightarrow A$ are given by

$$
\begin{align*}
& t \mapsto a t+b t \tau+c t \tau^{2}  \tag{16}\\
& \tau \mapsto d \tau+e \tau^{2}+f t \tau+g t \tau^{2}
\end{align*}
$$

By Proposition 1 we find the operators $N_{7}, \ldots, N_{13}$ in the form of the curves in Aut $(A)$ defined by reparametrization, e.g. $N_{7}$ by reparametrization $t \mapsto a t, \tau \mapsto \tau$ or $N_{8}$ by reparametrization $\tau \mapsto b \tau, t \mapsto t$.

Now we prove the main result of this Section.

Proposition 9. All natural operators $T \mathbb{R}^{m} \rightarrow C^{\infty}\left(T^{*} T T_{1}^{2} \mathbb{R}^{m}, \mathbb{R}\right), m \geq 3$, are of the form

$$
\begin{equation*}
h\left(\widetilde{N_{1}}, \widetilde{N_{2}}, \ldots, \widetilde{N_{13}}\right) \tag{17}
\end{equation*}
$$

where $h: \mathbb{R}^{13} \rightarrow \mathbb{R}$ is an arbitrary smooth function and $N_{1}, \ldots, N_{13}$ are the natural operators from Lemma 8.

Proof. We apply the procedure explained before Lemma 8. According to the immersion theorem, we can consider $i$ in the form

$$
y^{i}=\delta_{2}^{i}, z_{1}^{i}=\delta_{3}^{i}, z_{2}^{i}=w_{1}^{i}=w_{2}^{i}=0
$$

for all $i=1, \ldots, m$. Let $q_{i} d x^{i}+r_{i}^{1} d z_{1}^{i}+r_{i}^{2} d z_{2}^{i}+p_{i} d y^{i}+s_{i}^{1} d w_{1}^{i}+s_{i}^{2} d w_{2}^{i}$ define the additional coordinates on $T^{*} T T_{1}^{2} M$. Taking the space $P$ over the element $i$, we obtain the following values of $\left.I_{i}=\widetilde{N_{i}}\left(\frac{\partial}{\partial x^{1}}\right) \right\rvert\, P$

$$
\begin{aligned}
& I_{1}=q_{1}, I_{2}=r_{1}^{1}, I_{3}=p_{1}, I_{4}=r_{1}^{2}, I_{5}=s_{1}^{1}, I_{6}=s_{1}^{2} \\
& \quad I_{7}=p_{2}, I_{8}=r_{3}^{1}, I_{9}=s_{2}^{1}, I_{10}=s_{2}^{2}, I_{11}=r_{3}^{2}, I_{12}=s_{3}^{1}, I_{13}=s_{3}^{2}
\end{aligned}
$$

The stabilizing group $H \subseteq G_{m}^{4}$ of the element $i$ and $\frac{\partial}{\partial x^{1}}$ can be considered as a subgroup of $i d_{\mathbb{R}} \times$ Diff $\mathbb{R}^{m-1}$. The group $H$ acts in the following way:

$$
\begin{align*}
& \bar{z}_{1}^{i}=a_{j}^{i} z_{1}^{j} \quad \bar{z}_{2}^{i}=a_{j}^{i} z_{2}^{j}+a_{j k}^{i} z_{1}^{j} z_{1}^{k} \quad \bar{y}^{i}=a_{j}^{i} y^{j}  \tag{18}\\
& \bar{w}_{1}^{i}=a_{j}^{i} w_{1}^{j}+a_{j k}^{i} z_{1}^{j} y^{k} \quad \bar{w}_{2}^{i}=a_{j}^{i} w_{2}^{j}+a_{j k}^{i} z_{2}^{j} y^{k}+2 a_{j k}^{i} z_{1}^{j} w_{1}^{k}+a_{j k l}^{i} z_{1}^{j} z_{1}^{k} y^{l}
\end{align*}
$$

for $i, j \geq 2$. It is useful to annihilate the excessive coordinates extra for $m=3$ and $m \geq 4$.
$m=\overline{3}$ : We must annihilate $p_{3}, r_{2}^{1}, r_{2}^{2}$ and $q_{2}, q_{3}$. It follows from the action of $H$, that $a_{j}^{i}=\delta_{j}^{i}$, and for $i, j \geq 2$ it holds $a_{33}^{i}=a_{23}^{i}=a_{233}^{i}=0$. Taking into account the action of $B_{m}^{4} \cap H$ on $T^{*} T T_{1}^{2}$, we have $\bar{q}_{2}=q_{2}-a_{2233}^{j} s_{j}^{2}, \quad \bar{q}_{3}=q_{3}-a_{2333}^{j} s_{j}^{2}$, so we can annihilate $q_{2}, q_{3}$ by means of $a_{2233}^{2}, a_{2333}^{2}$ in the case $s_{2}^{2} \neq 0$. Furthermore $B_{m}^{3} \cap H$ turns $p_{3}$ to $\bar{p}_{3}=p_{3}-a_{333}^{j} s_{j}^{2}$ and $r_{2}^{1}$ to $\bar{r}_{2}^{1}-2 a_{223}^{j} s_{j}^{2}$. Thus we can annihilate $p_{3}$ and $r_{2}^{1}$ by means of $a_{333}^{2}$ and $a_{223}^{2}$ if $s_{2}^{2} \neq 0$. It remains to annihilate $r_{2}^{2}$. Since $B_{m}^{2} \cap H$ turns $r_{2}^{2}$ to $\bar{r}_{2}^{2}=r_{2}^{2}-a_{22}^{j} s_{j}^{2}$, we can achieve $r_{2}^{2}=0$ by means of $a_{22}^{2}$ in the case of non-zero $s_{2}^{2}$. Since the condition $s_{2}^{2} \neq 0$ determines a dense subset in $P$, our claim is proved for $m=3$.

In the case $m \geq 4$ we put $a_{j}^{i}=\delta_{j}^{i}$. Analogously to the case $m=3$ we obtain $a_{33}^{i}=a_{23}^{i}=a_{233}^{i}=0$ from (18). We can annihilate $q_{i}$ for $i \geq 2$ by means of $a_{i 233}^{2}$ in the case $s_{2}^{2} \neq 0, p_{i}$ by $a_{i 33}^{2}$ for $i \geq 3$ and $r_{i}^{1}$ by $a_{i 23}^{2}$ for $i=2$ or $i \geq 4$ in the case $s_{2}^{2} \neq 0$. It remains to annihilate $r_{i}^{2}$ for $i=2$ or $i \geq 4$, which can be done by means of $a_{i 2}^{2}$ in the case $s_{2}^{2} \neq 0$. Since the condition $s_{2}^{2} \neq 0$ defines a dense subset of $P$, our claim is proved for the case $m \geq 4$ too.

Now we show, how the generating operators $T \rightarrow T T^{*} T_{1}^{2}$ can be found by means of the natural operators $T \rightarrow C^{\infty}\left(T^{*} T T_{1}^{2}, \mathbb{R}\right)$. Let $G$ be a natural bundle. A natural operator $T \rightarrow C^{\infty}\left(T^{*} G, \mathbb{R}\right)$ is called a natural $T$-function. Every natural operator $D: T \rightarrow T G$ determines a natural $T$-function $\tilde{D}_{M}: T^{*} G M \rightarrow \mathbb{R}$, defined by $\tilde{D}_{M}(w)=\left\langle D_{M}(q w), w\right\rangle, w \in T^{*} G M, q: T^{*} G \rightarrow G$, which is linear on fibers. Conversely, let $f_{M}$ be a natural $T$-function linear on fibers $T^{*}(G M)$. Then $f_{M} \mid T_{z}^{*}(G M)$, where $z \in G M$, is identified with an element $\tilde{f}_{M}(z)$ from the dual vector space $T_{z}(G M)$. Thus we obtain a natural operator $\tilde{f}_{M}: T \rightarrow T G$ and a canonical bijection between natural operators $T \rightarrow T G$ and natural $T$-functions, which are linear on fibers of $T^{*}(G M)$.

Let $x^{i}$ be the standard coordinates on $\mathbb{R}^{m}$ and $p_{i} d x^{i}$ define the additional coordinates $p_{i}$ on $T^{*} \mathbb{R}^{m}$. Let $x^{i}, p_{i}$ induce the coordinates $X_{1}^{i}=d x^{i}, P_{i}=d p_{i}$ on $T T^{*} \mathbb{R}^{m}$. We can also define the additional coordinates $\xi_{i}, \eta^{i}$ on $T^{*} T^{*} \mathbb{R}^{m}$ by $\xi_{i} d x^{i}+\eta^{i} d p_{i}$. Furthermore, let $x^{i}$ induce the coordinates $Y^{i}=d x^{i}$ on $T \mathbb{R}^{m}$ and the additional coordinates $\alpha_{i}, \beta_{i}$ on $T^{*} T \mathbb{R}^{m}$ be defined by $\alpha_{i} d x^{i}+\beta_{i} d Y^{i}$.

We have the natural equivalence $s: T T^{*} \rightarrow T^{*} T$ by Modugno, Stefani, [8], and the natural equivalence $t: T T^{*} \rightarrow T^{*} T^{*}$ by Kolář, Radziszewski, [7],

$$
\begin{align*}
& s\left(x^{i}, p_{i}, X_{1}^{i}, P_{i}\right)=\left(x^{i}, Y^{i}, \alpha_{i}, \beta_{i}\right), \text { where } Y^{i}=X_{1}^{i}, \alpha_{i}=P_{i}, \beta_{i}=p_{i}  \tag{19}\\
& t\left(x^{i}, p_{i}, X_{1}^{i}, P_{i}\right)=\left(x^{i}, p_{i}, \xi_{i}, \eta^{i}\right), \quad \text { where } \xi_{i}=P_{i}, \eta^{i}=-X_{1}^{i}
\end{align*}
$$

Let the standard coordinates $x^{i}$ on $\mathbb{R}^{m}$ induce the coordinates $z_{1}^{i}=\frac{\partial x^{i}}{\partial \tau}$, $z_{2}^{i}=\frac{\partial^{2} x^{i}}{\partial \tau^{2}}$ on $T_{1}^{2} \mathbb{R}^{m}$ and the additional coordinates on $T^{*} T_{1}^{2} \mathbb{R}^{m}$ be defined by $p_{i} d x^{i}+s_{i}^{1} d z_{1}^{i}+s_{i}^{2} d z_{2}^{i}$. Further, define the additional coordinates on $T^{*} T^{*} T_{1}^{2} \mathbb{R}^{m}$ by $q_{i} d x^{i}+r_{i}^{1} d z_{1}^{i}+r_{i}^{2} d z_{2}^{i}-y^{i} d p_{i}-w_{1}^{i} d s_{i}^{1}-w_{2}^{i} d s_{i}^{2}$.

Clearly, $N: T \rightarrow C^{\infty}\left(T^{*} T T_{1}^{2}, \mathbb{R}\right)$ is a natural operator if and only if $A=$ $N \circ s \circ t^{-1}$ is a natural operator $T \rightarrow C^{\infty}\left(T^{*} T^{*} T_{1}^{2}, \mathbb{R}\right)$.

Transforming all the generating natural operators $T \rightarrow C^{\infty}\left(T^{*} T T_{1}^{2}, \mathbb{R}\right)$ into the generating natural operators $T \rightarrow C^{\infty}\left(T^{*} T^{*} T_{1}^{2}, \mathbb{R}\right)$ and among the transformed ones selecting those, which are linear on fibers over $T^{*} T_{1}^{2}$, we finally obtain the natural operators $A_{1}, A_{2}, A_{3}, A_{4}, A_{5}, A_{6}, A_{7}, A_{8}$ from Section 1.

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