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ON A CODIMENSION THREE BIFURCATION

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In this paper we study unfoldings of the vector field

$$\dot{x} = X_0(x) = Ax + G(x),$$

where $x = (x_1, x_2, x_3)$, $G \in C^{\infty}$, G(x) = o(||x||), the matrix $A = (a_{ij})$ is equivalent to the nilpotent matrix S with 1's just above the diagonal and 0's elsewhere. Some results contained in this paper have been announced in [15]. The unfoldings of the above vector field, possessing symmetry under the change of sign, $X_0(x) = -X_0(-x)$, are studied in [16].

Under generic hypotheses on the quadratic terms, we derive a normal form for unfoldings of X_0 which enables us to find the bifurcation diagram of the critical points. We show that generically there is a curve $Z_2(Z_{1c})$ in the parameter space, where the linear part of the corresponding vector field, computed at a critical point, has zero as an eigenvalue of multiplicity two (a couple of pure imaginary eigenvalues and one zero eigenvalue). Using Bogdanov's results [3] we describe the bifurcations near the curve Z_2 . The case of the codimension two singularity, which occurs on the curve Z_{1c} , is more complicated. It has been partially solved by several authors [5], [7-10]. There are a number of different cases of very complicated bifurcations near the curve Z_{1c} . The problem of global bifurcations of the phase portraits when the parameter goes from a nieghbourhood of Z_2 to a neighbourhood of Z_{1c} remains open.

1. PRELIMINARY LEMMAS

Consider an unfolding of the vector field X_0 , represented by the three-parameter family of vector fields

(1.1) $\dot{x} = f(x, \varepsilon),$

where $f = (f_1, f_2, f_3) \in C^{\infty}$, $x = (x_1, x_2, x_3)$, $\varepsilon = (\varepsilon_1, \varepsilon_2, \varepsilon_3)$. We also write $f_{\varepsilon}(x)$ instead of $f(x, \varepsilon)$.

The vector field $f_0 = X_0$ may be rewitten as

(1.2)
$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \end{bmatrix} = A \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix} + \begin{bmatrix} (Px, x) + h_1(x) \\ (Qx, x) + h_2(x) \\ (Rx, x) + h_3(x) \end{bmatrix}$$

where $P = (p_{ij})$, $Q = (q_{ij})$, $R = (r_{ij})$ are symmetric matrices, $A = (a_{ij})$, (\cdot, \cdot) is the scalar product in R^3 , $h_i(x) = o(||x||^2)$, i = 1, 2, 3.

There exists a linear change of coordinates y = Nx such that (1.2) becomes

(1.3)

$$\begin{aligned}
\dot{y}_{1} &= y_{2} + (\tilde{P}y, y) + g_{1}(y), \\
Y_{0} &: \dot{y}_{2} &= y_{3} + (\tilde{Q}y, y) + g_{2}(y), \\
\dot{y}_{3} &= (\tilde{R}y, y) + g_{3}(y), \\
\begin{bmatrix} (\tilde{P}y, y) \\ (\tilde{Q}y, y) \\ (\tilde{R}y, y) \end{bmatrix} &= N \begin{bmatrix} ((N^{-1})' PN^{-1}y, y) \\ ((N^{-1})' QN^{-1}y, y) \\ ((N^{-1})' RN^{-1}y, y) \end{bmatrix}, \quad g_{i}(y) = o(||y||^{2}), \quad i = 1, 2, 3, \\
\end{aligned}$$

 $(N^{-1})'$ is the transpose of N^{-1} .

Lemma 1. There exists a smooth local diffeomorphism Φ transforming the vector field Y_0 to

(1.4)
$$\Phi_*Y_0: \dot{x}_1 = x_2, \dot{x}_2 = x_3, \dot{x}_3 = (Tx, x) + h(x)$$

where $T = \tilde{R} + T_0$, $h(x) = o(||x||^2)$, $T_0 = (t_{ij}^0)$ is a symmetric matrix with $t_{11}^0 = 0$, $t_{12}^0 = \tilde{q}_{11}$, $t_{13}^0 = \tilde{p}_{11} + \tilde{q}_{12}$, $\tilde{P} = (\tilde{p}_{ij})$, $\tilde{Q} = (\tilde{q}_{ij})$, $\tilde{R} = (\tilde{r}_{ij})$.

Proof. The diffeomorphism $H: z_1 = y_1, z_2 = y_2 + (\tilde{P}y, y) + g_1(y), z_3 = y_3$ transforms (1.3) to the vector field $H_* Y_0(z) = z_2 \partial/\partial z_1 + (z_3 + (\hat{Q}z, z) + \hat{g}_2(z)) \partial/\partial z_2 + (\hat{R}z, z) + \hat{g}_3(z)) \partial/\partial z_3$, where $\hat{g}_i(z) = o(||z||^2)$, $i = 2, 3, \hat{Q} = \tilde{Q} + \hat{P}, \hat{R} = \tilde{R}, \hat{P} = (\hat{p}_{ij})$ is a symmetric matrix with $\hat{p}_{11} = 0, \hat{p}_{12} = \tilde{p}_{11}, \hat{p}_{13} = \tilde{p}_{12}, \hat{p}_{22} = 2\tilde{p}_{12}, \hat{p}_{23} = \tilde{p}_{13} + \tilde{p}_{22}, \hat{p}_{33} = 2\tilde{p}_{23}$. Then $\Phi = K \circ H$, where $K: x_1 = z_1, x_2 = z_2, x_3 = z_3 + (\hat{Q}z, z) + \hat{g}_2(z)$.

Lemma 2. Let $T = (t_{ij})$ be the matrix from (1.4). Then the numbers $q = t_{1j}/t_{11}$, j = 2, 3, are invariant with respect to regular transformations of coordinates in the phase space that keep the origin fixed.

Proof. Consider a diffeomorphism of the form $R: y_i = x_i + X_i(x) + o(||x||^2)$, i = 1, 2, 3, where the functions X_i are homogeneous polynomials of degree 2. We assume that R maps (1.4) to a vector field of the same form. Any diffeomorphism transforming (1.2) to the form (1.4) and keeping the origin fixed, is composed of the mapping $\Psi = \Phi \circ N$, where Φ , N are as above, of a mapping of the form R and of a linear mapping ϱ , which does not change the linear part of the vector field (1.4). The mapping ϱ must be of the form $\varrho(x) = Dx$, where $D = (d_{ij}), d_{kk} = \lambda, k = 1, 2, 3,$ $d_{12} = d_{23} = \varepsilon, d_{13} = \delta, d_{ij} = 0$ for $i > j, \lambda, \varepsilon, \delta$ are real numbers, $\lambda \neq 0$. It suffices to prove the invariance of q with respect to the mappings R and ϱ .

By using the fact that the mapping R preserves the form (1.4) it is easy to check that

$$X_{i+1}(x) = \frac{\partial X_i}{\partial x_1} x_2 + \frac{\partial X_i}{\partial x_2} x_3, \quad i = 1, 2,$$

and therefore

$$X_{3}(x) = \frac{\partial^{2} X_{1}}{\partial x_{1}^{2}} x_{2}^{2} + 2 \frac{\partial^{2} X_{1}}{\partial x_{1} \partial x_{2}} x_{2} x_{3} + \frac{\partial^{2} X_{1}}{\partial x_{2}^{2}} x_{3}^{2} + \frac{\partial X_{1}}{\partial x_{1}} x_{3}$$

This implies that the vector field $R_*(\Phi_*Y_0)$ has the form (1.4) with a matrix $T' = (t'_{ij})$ instead of the matrix $T = (t_{ij})$ and $t'_{1j} = t_{1j}$, j = 1, 2, 3, i.e. the mapping R does not change the numbers t_{1j} , j = 1, 2, 3.

It remains to prove the invariance of q with respect to the mapping ϱ . This mapping transforms the vector field (1.4) to the form (1.3), where $\tilde{P} = (\tilde{p}_{ij}) = \delta \tilde{T}$, $\tilde{Q} =$ $= (\tilde{q}_{ij}) = \varepsilon \tilde{T}$, $\tilde{R} = (\tilde{r}_{ij}) = \lambda \tilde{T}$, $\tilde{T} = (\tilde{t}_{ij}) = (D^{-1})' TD^{-1}$, $\tilde{p}_{11} = \delta \lambda^{-2} t_{11}$, $\tilde{q}_{11} =$ $= \varepsilon \lambda^{-2} t_{11}$, $\tilde{q}_{12} = -\varepsilon^2 \lambda^{-3} t_{11} + \varepsilon \lambda^{-2} t_{12}$, $\tilde{r}_{12} = -\varepsilon \lambda^{-2} t_{11} + \lambda^{-1} t_{12}$, $r_{11} = \lambda^{-1} t_{11}$, $\tilde{r}_{13} = (\varepsilon^2 \lambda^{-3} - \delta \varepsilon \lambda^{-2}) t_{11} - \varepsilon \lambda^{-2} t_{12} + \lambda^{-1} t_{13}$. By Lemma 1 there exists a smooth local diffeomorphism transforming the vector field to the form (1.4), with a matrix $\tilde{T} = (\tilde{t}_{ij}) = \tilde{R} + T_0$ instead of the matrix T, and the first row of the matrix T_0 is $(0, \tilde{q}_{11}, \tilde{q}_{12} + \tilde{p}_{11})$. Therefore $\tilde{t}_{11} = \lambda^{-1} t_{11}$, $\tilde{t}_{12} = \tilde{r}_{12} + \tilde{q}_{11} = \lambda^{-1} t_{12}$, $\tilde{t}_{13} = \tilde{r}_{13} + \tilde{q}_{12} + \tilde{p}_{11} = \lambda^{-1} t_{13}$ and the proof is complete.

2. NORMAL FORM

By Lemma 1 the family (1.1) may be written in the form

(2.1)
$$\dot{x}_1 = x_2 + v_1(x, \varepsilon),$$

 $\dot{x}_2 = x_3 + v_2(x, \varepsilon).$
 $\dot{x}_3 = t_{11}x_1^2 + t_{12}x_1x_2 + t_{13}x_1x_3 + t_{23}x_2x_3 + t_{22}x_2^2 + t_{33}x_3^2 + v_3(x, \varepsilon),$

where $v_i(x, 0) \equiv 0$, $i = 1, 2, v_3(x, 0) = o(||x||^2)$.

Assuming $t_{11} \neq 0$, we may introduce new coordinates $y = t_{11}x$ and then (2.1) becomes

(2.2)
$$\dot{y}_1 = y_2 + \tilde{v}_1(y, \varepsilon),$$

 $\dot{y}_2 = y_3 + \tilde{v}_2(y, \varepsilon),$
 $\dot{y}_3 = y_1^2 + \omega_1 y_1 y_2 + \omega_2 y_1 y_3 + \tilde{t}_{23} y_2 y_3 + \tilde{t}_{22} y_2^2 + \tilde{t}_{33} y_3^3 + \tilde{v}_3(y, \varepsilon),$

where $\tilde{v}_i(y, 0) \equiv 0$, i = 1, 2, $\tilde{v}_3(y, 0) = o(||y||^2)$, $\omega_j = t_{1j+1}/t_{11}$, j = 1, 2, are invariants of the germ, represented by the family (1.1).

Introducing again new coordinates $u_1 = y_1$, $u_2 = y_2 + \tilde{v}_1(y, \varepsilon)$, $u_3 = y_3$, we obtain a family of the form (2.2) with $\tilde{v}_1 \equiv 0$. Transforming the resulting family by the diffeomorphism $z_1 = u_1$, $z_2 = u_2$, $z_3 = u_3 + \tilde{v}_2(y, \varepsilon)$, we get a family of the form (2.2) with $\tilde{v}_1 \equiv 0$. This family may be written in the form

$$\begin{aligned} (2.3) \qquad \dot{z}_1 &= z_2 ,\\ \dot{z}_2 &= z_3 ,\\ \dot{y}_3 &= \tilde{F}(z_1,\varepsilon) + z_2 \tilde{\mathcal{Q}}_1(z_1,\varepsilon) + z_3 \tilde{\mathcal{Q}}_2(z_1,\varepsilon) + z_2 \tilde{\mathcal{Q}}_3(z_3,\varepsilon) + \\ &+ z_2^2 \tilde{\Psi}_1(z,\varepsilon) + z_3^2 \tilde{\Psi}_2(z,\varepsilon) , \end{aligned}$$

where \tilde{F} , \tilde{Q}_i , $\tilde{\Psi}_j$, i = 1, 2, 3, j = 1, 2, are C^{∞} -functions,

$$\tilde{F}(0,0) = \frac{\partial \tilde{F}(0,0)}{\partial z_1} = 0, \quad \frac{\partial^2 \tilde{F}(0,0)}{\partial z_1^2} = 2, \quad \frac{\partial \tilde{Q}_i(0,0)}{\partial z_1} = \omega_i,$$
$$i = 1, 2, \quad \tilde{Q}_k(0,0) = 0, \quad k = 1, 2, 3.$$

Lemma 3. If $\omega_1 \neq 0$, $\omega_2 \neq 0$, then there exists a smooth regular mapping $y = y(z, \varepsilon)$, y(0, 0) = 0, transforming the family (2.3) to the form

(2.4)
$$\dot{y}_1 = y_2,$$

 $\dot{y}_2 = y_3,$
 $\dot{y}_3 = F(y_1, \varepsilon) + \beta(\varepsilon) y_2 + y_1 y_2 G_1(y_1, \varepsilon) + y_1 y_3 G_2(y_1, \varepsilon) +$
 $+ y_2 G_3(y_3, \varepsilon) + y_2^2 \Psi_1(y, \varepsilon) + y_3^2 \Psi_2(y, \varepsilon),$

where F, G_1 , G_2 , G_3 , β , Ψ_1 , Ψ_2 are smooth functions,

$$F(0,0) = \frac{\partial F(0,0)}{\partial y_1} = 0, \quad \frac{\partial^2 F(0,0)}{\partial y_1^2} = 2,$$

$$G_i(0,0) = \omega_i, \quad i = 1, 2, \quad \beta(0) = 0, \quad G_3(y_3,0) = 0(|y_3|)$$

Proof. Let $y_1 = z_1 - \alpha(\varepsilon)$, $y_2 = z_2$, $y_3 = z_3$, where α is any smooth function. Then the family (2.3) becomes $\dot{y}_1 = y_2$, $\dot{y}_2 = y_3$, $\dot{y}_3 = \tilde{F}(y_1 + \alpha(\varepsilon), \varepsilon) + y_2 \tilde{Q}_1(y_1 + \alpha(\varepsilon), \varepsilon) + y_3 \tilde{Q}_2(y_1 + \alpha(\varepsilon), \varepsilon) + y_2 \tilde{Q}_3(y_3, \varepsilon) + y_2^2 \hat{\Psi}_1(y, \varepsilon) + y_3^2 \hat{\Psi}_2(y, \varepsilon)$, where $\hat{\Psi}_i(y, \varepsilon) = \tilde{\Psi}_i(y_1 + \alpha(\varepsilon), y_2, y_3, \varepsilon)$, $\tilde{Q}_2(y_1 + \alpha(\varepsilon), \varepsilon) = \tilde{Q}_2(\alpha(\varepsilon), \varepsilon) + y_1 \hat{Q}_2(y_1, \varepsilon)$, $\tilde{Q}_2(0, 0) = 0$, $\partial \tilde{Q}_2(0, 0)/\partial y_1 = \omega_2$, $\hat{Q}_2(0, 0) = \omega_2$. Since $\omega_2 = 0$, the implicit function theorem implies that there exists a neighbourhood U of $0 \in R^3$ and a smooth function $\alpha : U \to R^1$ such that $\alpha(\sigma) = 0$, $\tilde{Q}_2(\alpha(\varepsilon), \varepsilon) = 0$ for all $\varepsilon \in U$. From Taylor's expansion of the function \tilde{Q}_1 we have $\tilde{Q}_1(y_1 + \alpha(\varepsilon), \varepsilon) = \beta(\varepsilon) + y_1 G_1(y_1, \varepsilon) + o(|y_1|)$, where β , $G_1 \in C^{\infty}$, $\beta(0) = 0$, $G_1(0, 0) = \partial \tilde{Q}_1(0, 0)/\partial y_1 = \omega_1$ and so the family obtained has the form (2.4).

If F is the function from Lemma 3, then by the Malgrange-Weierstrass preparation theorem (see [14]) there exist smooth functions $\varphi_i(\varepsilon)$, $\varphi_i(0) = 0$, i = 1, 2, $\Theta(y_1, \varepsilon)$, $\Theta(0, 0) = 1$, such that $F(y_1, \varepsilon) = (y_1^2 + \varphi_2(\varepsilon) y_1 + \varphi_1(\varepsilon)) \Theta(y_1, \varepsilon)$ and therefore the family (2.4) may be written as

$$\begin{split} \dot{y}_1 &= y_2 , \\ \dot{y}_2 &= y_3 , \\ \dot{y}_3 &= (\varphi_1(\varepsilon) + \varphi_2(\varepsilon) y_1 + y_1^2 + \varphi_3(\varepsilon) y_2 + y_1 y_2 Q_1(y_1, \varepsilon) + \\ &+ y_1 y_3 Q_2(y_1, \varepsilon) + y_2 Q_3(y_3, \varepsilon) + y_2^2 \Phi_1(y, \varepsilon) + y_3^2 \Phi_2(y, \varepsilon)) \Theta(y_1, \varepsilon) , \end{split}$$

where Θ , $Q_i \varphi_i$, $\Phi_j \in C^{\infty}$, $\varphi_i(0) = 0$, $\Phi_j(0, 0) = 0$, $i = 1, 2, 3, j = 1, 2, Q_k(0, 0) = \omega_k$, $k = 1, 2, \Theta(0, 0) = 1$.

We have assumed $\omega_1 \neq 0$, $\omega_2 \neq 0$ in the previous lemmas. Now we show that these conditions are generically satisfied in the space of three-parameter families of vector fields of the form (1.1). To this aim, we define some algebraic manifolds.

Let S be the nilpotent matrix with 1's just above the diagonal and 0's elsewhere. For the matrix $A = (a_{ij})$ of the linear part of the vector field X_0 there exists a regular matrix $N = (c_{ij})$ such that $NAN^{-1} = S$. Since rank A = 2, there exists at least one nonzero minor of order 2. There is no loss of generality to assume

$$A_3 = \det \begin{bmatrix} a_{12} & a_{13} \\ a_{22} & a_{23} \end{bmatrix} \neq 0 \,.$$

This condition corresponds to a stratum of some algebraic manifold, which will be specified later. Under the assumption that some other minor of order 2 is nonzero, we have to do with another stratum of this algebraic manifold and all computations for this case are similar to those for $A_3 \neq 0$.

First let us express the elements of the matrix N as functions of elements of the matrix A. If $A_3 \neq 0$, then

$$c_{11} = \frac{a_{23}}{A_3} \det \begin{bmatrix} a_{12} & A_2 \\ a_{13} & A_3 \end{bmatrix}, \quad c_{12} = \frac{a_{13}}{A_3} \det \begin{bmatrix} a_{12} & A_2 \\ a_{13} & A_3 \end{bmatrix}, \quad c_{13} = 0,$$

$$c_{21} = \frac{1}{A_3} \det \begin{bmatrix} A_2 & a_{23} \\ A_3 & a_{33} \end{bmatrix}, \quad c_{22} = \frac{1}{A_3} \det \begin{bmatrix} a_{12} & A_2 \\ a_{13} & A_3 \end{bmatrix}, \quad c_{23} = 0,$$

 $c_{31} = A_1, c_{32} = A_2, c_{33} = A_3$, where

$$A_1 = \det \begin{bmatrix} a_{22} & a_{32} \\ a_{23} & a_{33} \end{bmatrix}, \quad A_2 = \det \begin{bmatrix} a_{32} & a_{12} \\ a_{33} & a_{13} \end{bmatrix}.$$

The characteristic equation of the matrix $A = (a_{ij})$ is

$$-\lambda^3 + c_2 \lambda^2 + c_1 \lambda + c_0 = 0,$$

where $c_0 = \det A$, $c_1 = a_{23}a_{32} + a_{11}a_{33} + a_{12}a_{21} - a_{22}a_{33} - a_{13}a_{31} - a_{11}a_{22}$ $c_2 = \operatorname{Sp} A = a_{11} + a_{22} + a_{33}$. Therefore A has zero as an eigenvalue of multiplicity 3 if and only if det A = 0, $\operatorname{Sp} A = 0$, $c_1 = 0$.

Denote by Γ_3^{∞} the set of all smooth vector fields on \mathbb{R}^3 . Let $j^k v(x)$ be the k-jet of $v \in \Gamma_3^{\infty}$ at a point x and let $J_3^k(x)$ be the set of all such k-jets. For $v \in \Gamma_3^{\infty}$, $j^2 v(x) = (v(x), D v(x), D^2 v(x))$, we may identify D v(x) with

$$\left(\frac{\partial v_1(x)}{\partial x_1}, \frac{\partial v_1(x)}{\partial x_2}, \frac{\partial v_1(x)}{\partial x_3}, \dots, \frac{\partial v_3(x)}{\partial x_3}\right) \in R^9$$

and because of the symmetry of the matrices $D^2 v_k(x)$ we may identify $D^2 v(x)$ with

$$\left(\frac{\partial^2 v_1(x)}{\partial x_1^2}, \frac{\partial^2 v_1(x)}{\partial x_2^2}, \frac{\partial^2 v_1(x)}{\partial x_3^2}, \frac{\partial^2 v_1(x)}{\partial x_1 \partial x_2}, \frac{\partial^2 v_1(x)}{\partial x_1 \partial x_3}, \frac{\partial^2 v_1(x)}{\partial x_2 \partial x_3}, \frac{\partial^2 v_3(x)}{\partial x_2 \partial x_3}\right) \in \mathbb{R}^{18}.$$

This means that the 2-jet $j^2 v(x)$ may be identified with

$$\left(v(x),\frac{\partial v_1(x)}{\partial x_1}, \ldots, \frac{\partial v_3(x)}{\partial x_3}, \frac{\partial^2 v_1(x)}{\partial x_1^2}, \ldots, \frac{\partial^2 v_3(x)}{\partial x_2 \partial x_3}\right) \in \mathbb{R}^{30}.$$

Let us define the following sets:

 $T_k = \{(a, A, B) \in J_3^2: F_i(a, A, B) = 0, i = 1, 2, 3, F_4(a, A, B) = t_{1k} = 0, a = 0,$ rank $A = 2\}, k = 1, 2, 3$, where $F_1 = \text{Sp } A$, $F_2 = \det A, F_3 = a_{23}a_{32} + a_{11}a_{33} + a_{12}a_{21} - a_{22}a_{33} - a_{13}a_{31} - a_{11}a_{22}, A = (a_{ij})$ and t_{11}, t_{12}, t_{13} are the elements of the matrix T from (1.4), which are functions of the elements of the matrices A, B (the elements of B are in fact the elements of the matrices P, Q, R from (1.2)). Direct computations show that

$$\begin{split} t_{11} &= c_{31} \big(\alpha_{11} c_{11}' + \alpha_{12} c_{21}' + \alpha_{13} c_{31}' \big) + c_{32} \big(\beta_{11} c_{11}' + \beta_{12} c_{21}' + \\ &+ \beta_{13} c_{31}' \big) + c_{33} \big(\gamma_{11} c_{11}' + \gamma_{12} c_{21}' + \gamma_{13} c_{31}' \big) \,, \\ t_{12} &= c_{31} \big(\alpha_{11} c_{12}' + \alpha_{12} c_{22}' + \alpha_{13} c_{32}' \big) + c_{32} \big(\beta_{11} c_{12}' + \beta_{12} c_{22}' + \\ &+ \beta_{13} c_{32}' \big) + c_{33} \big(\gamma_{11} c_{12}' + \gamma_{12} c_{22}' + \gamma_{13} c_{32}' \big) + c_{21} \big(\alpha_{11} c_{11}' + \\ &+ \alpha_{12} c_{21}' + \alpha_{13} c_{31}' \big) + c_{22} \big(\beta_{11} c_{11}' + \beta_{12} c_{21}' + \beta_{13} c_{31}' \big) \,, \\ t_{13} &= \big(c_{31} \alpha_{13} + c_{32} \beta_{13} + c_{33} \gamma_{13} \big) \, c_{33}' + c_{21} \big(\alpha_{11} c_{12}' + \alpha_{12} c_{22}' + \\ &+ \alpha_{13} c_{32}' \big) + c_{22} \big(\beta_{11} c_{12}' + \beta_{12} c_{22}' + \beta_{13} c_{32}' \big) + c_{11} \big(\alpha_{11} c_{11}' + \\ &+ \alpha_{12} c_{21}' + \alpha_{13} c_{31}' \big) + c_{12} \big(\beta_{11} c_{11}' + \beta_{12} c_{21}' + \beta_{13} c_{31}' \big) \,, \end{split}$$

where $\alpha_{11} = c'_{11}p_{11} + c'_{21}p_{12} + c'_{31}p_{13}$, $\alpha_{12} = c'_{11}p_{12} + c'_{21}p_{22} + c'_{31}p_{23}$, $\alpha_{13} = c'_{11}p_{13} + c'_{21}p_{23} + c'_{31}p_{33}$ (the same for β_{1k} and γ_{1k} , k = 1, 2, 3, where we have q_{ij} and r_{ij} , respectively, instead of p_{ij}), $NAN^{-1} = S$, S is as above, $N = (c_{ij})$, $N^{-1} = (c_{ij})$. The elements c_{ij} are functions of the elements of the matrix A expressed as above (we assume $A_3 \neq 0$).

Lemma 4. The sets T_1 , T_2 , T_3 are smooth submanifolds of J_3^2 of codimension 7.

Proof. (1) for T_2 : Let $F = (F_1, F_2, F_3, F_4)$: $R^{30} \rightarrow R^4$, where F_i , i = 1, 2, 3, 4, are the functions from the definition of the sets T_k . It suffices to show that rank DF = 4.

$$F_{ij} = \det \begin{bmatrix} \frac{\partial F_1}{\partial a_{11}} & \frac{\partial F_1}{\partial a_{31}} & \frac{\partial F_1}{\partial a_{21}} & \frac{\partial F_1}{\partial r_{ij}} \\ \cdots \cdots \cdots \cdots \cdots \cdots \\ \frac{\partial F_4}{\partial a_{11}} & \frac{\partial F_4}{\partial a_{31}} & \frac{\partial F_4}{\partial a_{21}} & \frac{\partial F_4}{\partial r_{ij}} \end{bmatrix} = \frac{\partial F_4}{\partial r_{ij}} A_{23},$$

where $A_{23} = a_{12}A_3 - a_{13}A_2$. It suffices to show that $F_{11}^2 + F_{12}^2 \neq 0$. By the above formulae for c_{ij} we have that $c_{11} = k_1A_{23}, c_{12} = k_2A_{23}, c_{13} = 0$ and hence $A_{23} \neq 0$.

Therefore it suffices to show that

$$\Delta = \left(\frac{\partial F_4}{\partial r_{11}}\right)^2 + \left(\frac{\partial F_4}{\partial r_{12}}\right)^2 \neq 0.$$

If we express c'_{ij} as functions of c_{ij} , then the formula for t_{12} yields $t_{12} = -d^2 c^3_{33} c_{12} c_{22} r_{11} + d^2 c^3_{33} (c_{12} c_{21} + c_{22} c_{11}) r_{12} + (\text{terms independent of } r_{11} \text{ and } r_{12}), d^{-1} = \det N$ and therefore $\Delta = d^4 c^6_{33} (c^2_{12} c^2_{22} + (c_{12} c_{21} + c_{22} c_{11})^2)$. Since $c_{33} = A_3 \neq 0, c_{22} = kA_{23} \neq 0$, we obtain that $\Delta = 0$ if and only if $c_{11} = c_{12} = 0$. However, $c_{13} = 0$, N is regular and therefore $\Delta \neq 0$.

(2) for $T_1: F_4 = t_{11} = d c_{33}^3 c_{22}^2 r_{11} + (\text{terms independent of } r_{11}), c_{22} = kA_{23} = 0$ $= 0, c_{33} = A_3 = 0$ and therefore $\partial F_4 / \partial r_{11} = 0$, i.e. rank DF = 4.

(3) for T_3 : $F_4 = t_{13} = d^2 c_{32}^2 c_{22} (c_{11}c_{22} - c_{12}c_{21}) q_{12}$ + (terms independent of q_{12}). This implies that $\partial F_4 / \partial q_{12} \neq 0$, i.e. rank DF = 4 and the proof of the lemma is complete.

Denote by D^{∞} the set of all smooth mapping from $R^3 \times R^3$ into R^3 and for any $f \in D^{\infty}$ define the mapping $\Psi_f \colon R^3 \times R^3 \to J_3^2$, $\Psi_f(x, \varepsilon) = j^2 f_{\varepsilon}(x)$, $(x, \varepsilon) \in R^3 \times R^3$ $(f_{\varepsilon}(x) = f(x, \varepsilon))$. As a consequence of Lemma 4 and Thom's transversality theorem (see e.g. [13, Theorem 3.1]) we obtain

Lemma 5. (1) There exists a residual subset D_1^{∞} of D^{∞} such that if $f \in D_1^{\infty}$, then $\Psi_f(R^3 \times R^3) \cap (T_1 \cup T_2 \cup T_3) = \emptyset$.

(2) If $X \subset \mathbb{R}^3 \times \mathbb{R}^3$ is a compact set, then there exists an open dense subset D_X of D^{∞} such that if $f \in D_X$, then $\Psi_f(X) \cap (T_1 \cup T_2 \cup T_3) = \emptyset$.

Let $\Sigma = \{(a, A) \in J_3^1 : a = 0, \text{ det } A = 0, \text{ Sp } A = 0, c_1 = 0, A \text{ is a nonzero matrix}\}, A = (a_{ij}), c_1 \text{ are as above. From the above computations it follows that } \Sigma \text{ is a smooth submanifold of } J_3^1 \text{ of codimension 6.}$

Definition. The family (1.1) is called nondegenrate, if $t_{11} \cdot t_{12} \cdot t_{13} \neq 0$ and

 $(\Phi_f \text{ transversally intersects } \Sigma \text{ at } (0,0)), \text{ where } \Phi_f(x,\varepsilon) = j^1 f_{\varepsilon}(x).$

Denote by H^{∞} the set of all families of vector fields of the form (1.1). As a consequence of Lemma 5 and Thom's transversality theorem we obtain the following lemma.

Lemma 6. The set of all nondegenerate families of vector fields $H_1^{\infty} \subset H^{\infty}$ is open dense in H^{∞} .

Let $f \in H_1^{\infty}$ and suppose that it is already in the form (2.5). Define the mapping $\sigma_f: \mathbb{R}^6 \to \mathbb{R}^6$, $\sigma_f(y, \varepsilon) = (f(y, \varepsilon), \operatorname{Sp} D_y f_{\varepsilon}(y), \operatorname{det} D_y f_{\varepsilon}(y), H_{\varepsilon}(y))$, where $D_y f_{\varepsilon}(y) = (a_{ij}(y, \varepsilon))$ is the differential of the mapping f at y, $\operatorname{Sp} D_y f_{\varepsilon}(y) = a_{11} + a_{22} + a_{33}$, $H_{\varepsilon}(y) = -a_{22}a_{33} + a_{32}a_{23} + a_{11}a_{33} - a_{13}a_{31} - a_{11}a_{22} + a_{12}a_{21}$. The form of the mapping σ_f and the forms of the functions defining the set imply that the trans-

versality condition (2.6) is equivalent to the regularity of the mapping σ_f at the origin, i.e. to the condition det $D\sigma_f(0, 0) \neq 0$.

If $f = (f_1, f_2, f_3)$, then Sp $D_y f_{\varepsilon}(y) = \partial f_3 / \partial y_3$, det $D_y f_{\varepsilon}(y) = \partial f_3 / \partial y_1$, $H_{\varepsilon}(y) = \partial f_3 / \partial y_2$. Using the form of the family (2.5) one can show that det $D\sigma_f(0, 0) = -\omega_2$ det $D\varphi(0)$, $\varphi = (\varphi_1, \varphi_2, \varphi_3)$. Since $\omega_2 \neq 0$ and det $\sigma_f(0, 0) \neq 0$ for $f \in H_1^{\infty}$, we obtain that det $D\varphi(0) \neq 0$. This enables us to introduce new coordinates in the parameter space: $\mu_i = \varphi_i(\varepsilon)$, i = 1, 2, 3, and we obtain a family of the form (2.5) with μ_i , $Q_i(y_1, \varphi^{-1}(\mu))$ (i = 1, 2), $Q_3(y_3, \varphi^{-1}(\mu))$, $\Phi_j(y, \varphi^{-1}(\mu)) = (\mu_1, \mu_2, \mu_3)$. Dividing the right-hand side of the resulting family by the function $\Theta(y_1, \varphi^{-1}(\mu))$ the family becomes

(2.7)
$$\dot{z}_1 = z_2 \tilde{\Theta}(z_1, \mu), \quad \dot{z}_2 = z_3 \tilde{\Theta}(z_1, \mu), \quad \dot{z}_3 = R(z, \mu),$$

where $\tilde{\Theta}$, $R \in C^{\infty}$, $\tilde{\Theta}(0, 0) = 1$ and R has the same form as the right-hand side of (2.5) with $\Theta \equiv 1$. This family is C^{∞} -equivalent to (2.5). Now, if we put $u_1 = z_1$, $u_2 = z_2 \tilde{\Theta}(z_1, \mu)$, $u_3 = z_3$, the family becomes

$$\dot{u}_1 = u_2, \quad \dot{u}_2 = u_3 \hat{\Theta}(u_1, \mu), \quad u_3 = \hat{R}(u, \mu),$$

where $\hat{\Theta}$, $\hat{R} \in C^{\infty}$, $\hat{\Theta}(0, 0) = 1$, $\hat{R}(u, \mu)$ has the same form as R. Finally, introducing new coordinates $y_1 = u_1$, $y_2 = u_2$, $y_3 = u_3 \hat{\Theta}(u_1, \mu)$, one obtains a family of the form

(2.8)
$$\dot{y}_1 = y_2,$$

 $\dot{y}_2 = y_3,$
 $\dot{y}_3 = \mu_1 + \mu_2 y_1 + y_1^2 + \mu_3 y_2 + y_1 y_2 Q_1(y_1, \mu) +$
 $+ y_1 y_3 Q_2(y_1, \mu) + y_2 Q_3(y_3, \mu) + y_2^2 \Phi_1(y, \mu) + y_3^2 \Phi_2(y, \mu),$

where $Q_1, Q_2, Q_3, \Phi_1, \Phi_2$ are smooth functions, $Q_1(0, 0) = \omega_1, Q_2(0, 0) = \omega_2$. We have proved the following theorem

We have proved the following theorem.

Theorem. There exists an open dense subset H_1^{∞} of the set H^{∞} of all three-parameter families of vector fields of the form (1.1) such that if $f \in H_1^{\infty}$, then f is nondegenerate, and it is possible to transform this family by a smooth regular transformation of coordinates in a sufficiently small neighbourhood of the origin in $R^3 \times R^3$ to the form (2.8), where ω_1, ω_2 are invariants of the germ, represented by the family f.

3. **BIFURCATION DIAGRAM**

Let $f \in H_1^{\infty}$ be a family of the form (2.8). All critical points of this family have the form $(y_1, 0, 0)$, where y_1 is a real root of the equation

$$y^2 + \mu_2 y + \mu_1 = 0 \, .$$

Let U be a neighbourhood of the origin in the parameter space and let S_k (k = = 0, 1, 2) be the set of all $\mu \in U$ for which (2.8) has k critical points.

Lemma 7. There exists a smooth function $\mu_1 = S(\mu_2)$ such that $S_1 = \{\mu = (\mu_1, \mu_2, \mu_3) \in U: \mu_1 = S(\mu_2), S(0) = S'(0) = 0, S''(0) > 0\}$, i.e. S_1 is a fold dividing U into two components, one of which is S_0 and the other is S_2 .

If $\mu \in S_2$, then the vector field (2.8) has two critical points $F = (\xi_1, 0, 0)$, $G = (\xi_2, 0, 0)$, where $\xi_1 = -\frac{1}{2}(\mu_2 - \nu)$, $\xi_2 = \frac{1}{2}(-\mu_2 - \nu)$, $\nu = (\mu_2^2 - 4\mu_1)^{1/2}$.

Let $S_{12} = S_1 \cup S_2$ and let $K = (\xi, 0, 0)$ be a critical point of (2.8). Denote by L(K) the matrix of the linear part of (2.8), computed at K. The characteristic equation of L(K) is

(3.1)
$$\lambda^3 - a_2 \lambda^2 - a_1 \lambda - a_0 = 0,$$

where $|a_0| = |v|$, $a_1 = (\mu_3 + \xi) Q_1(\xi, \mu)$, $a_2 = \xi Q_2(\xi, \mu)$. If $\mu \in S_1$, then the vector field (2.8) has one critical point $K = (\xi, 0, 0)$, where $\xi = -\frac{1}{2}\mu_2$. If $\mu \in S_2$, then the matrix L(F) (L(G)) has the characteristic equation of the form (3.1), where $\xi = -\frac{1}{2}i_1(\xi = \xi_2)$ and $a_0 = v > 0$ ($a_0 = -v < 0$).

First assume $\mu \in S_1$. Then the matrix L(K) has zero as an eigenvalue. Obviously, it is of multiplicity 2 if and only if, in addition to the identity $\mu_2^2 - 4\mu_1 = 0$ defining the set S_1 , the following holds:

$$H(\mu_2, \mu_3) = (\mu_3 - \frac{1}{2}\mu_2) Q_1(-\frac{1}{2}\mu_2, \mu) = 0, \quad \mu_2 \neq 0.$$

Since $Q_1(0, 0) = \omega_1 \neq 0$, the last identity is satisfied in a sufficiently small neighbourhood of the origin only if $\mu_3 = \eta(\mu_2) = \frac{1}{2}\mu_2$. If $\chi(t) = (\frac{1}{4}t^2, t, \eta(t))$ and W is a neighbourhood of the origin in R^1 , then $\chi(W)$ is a one-dimensional smooth submanifold of S_1 . For $\mu \in Z_2(K) = \chi(W) \setminus \{0\}$, the matrix L(K) has zero as an eigenvalue of multiplicity 2 and the third eigenvalue is $\lambda_3 = -\frac{1}{2}\mu_2Q_1(-\frac{1}{2}\mu_2, \mu)$. The matrix L(K) has zero as an eigenvalue of multiplicity 1 and a couple of pure imaginary eigenvalues if and only if $a_0 = 0$, $a_2 = 0$, $a_1 < 0$, i.e. $\mu_1 = 0$, $\mu_2 = 0$, $\mu_3\omega_1 < 0$. Denote $Z_{1c} = \{\mu: \mu_1 = 0, \mu_2 = 0, \mu_3\omega_1 < 0\}$.

We have proved the following lemma.

Lemma 8. There exist one-dimensional smooth submanifolds Z_2 and Z_{1c} of S_1 such that the following holds:

- (1) Z_2 is the set of all $\mu \in U$ (U is a neighbourhood of the origin) for which the matrix L(K) has eigenvalues: $\lambda_1 = \lambda_2 = 0, \lambda_3 \neq 0$, where sign $\lambda_3 = -\text{sign } \mu_2 \omega_1$;
- (2) Z_{1c} is the set of all $\mu \in U$ for which the matrix has one zero eigenvalue and a couple of pure imaginary eigenvalues

(3)
$$\overline{Z}_2 \setminus Z_2 = \{0\}, \quad \overline{Z}_{1c} \setminus Z_{1c} = \{0\}.$$

Now assume $\mu \in S_2$. By means of the substitution $z + \frac{1}{3}a_2$ for λ in the characteristic equation (3.1) of the matrix L(K) we obtain

$$(3.2) z^3 + 3pz + 2q = 0,$$

1	1
T	T
_	_

where

(3.3)
$$p = -\frac{1}{3}(a_1 + \frac{1}{3}a_2^2), \quad q = -\frac{1}{2}(a_0 + \frac{1}{3}a_1a_2 + \frac{2}{27}a_2^3).$$

Fig. 1.

The discriminant of the equation (3.2) is $D = D(\mu) = q^2 + p^3$. Let us introduce new coordinates on S_{12} via the mapping

(3.4)

$$v_{1} = a_{0} = (\mu_{2}^{2} - 4\mu_{1})^{1/2},$$

$$\varrho_{F}: v_{2} = a_{2} = -\frac{1}{2}(\mu_{2} - (\mu_{2}^{2} - 4\mu_{1})^{1/2}) Q_{2}(\xi_{1}, \mu),$$

$$v_{3} = a_{1} = [\mu_{3} - \frac{1}{2}(\mu_{2} - (\mu_{2}^{2} - 4\mu_{1})^{1/2})] Q_{1}(\xi_{1}, \mu)$$

Obviously, ρ_F is a smooth diffeomorphism on S_2 , but it is not C^1 on S_1 and

(3.5)
$$\widetilde{S}_1 = \varrho_F(S_1) = \partial \varrho_F(S_{12}) = \{ \nu = (\nu_1, \nu_2, \nu_3) : \nu_1 = 0 \}.$$

In these coordinates the characteristic equation of L(F) is

$$(3.6) \qquad \qquad \lambda^3 - \nu_2 \lambda^2 - \nu_3 \lambda - \nu_1 = 0$$

The discriminant of this equation is $D_F = D_F(v) = p^3 + q^2$, where

(3.7)
$$p = -\frac{1}{3}(v_3 + \frac{1}{3}v_2^2), \quad q = -\frac{1}{2}(v_1 + \frac{1}{3}v_2v_3 + \frac{2}{27}v_2^3).$$

Denote $\mathscr{D}_F = \{v: D_F(v) = 0\}, \ \mathscr{D}_F^+ = \{v: D_F(v) > 0\}, \ \mathscr{D}_F^- = \{v: D_F(v) < 0\}.$

Lemma 9. If $v \in \mathcal{D}_F^-(\mathcal{D}_F; \mathcal{D}_F^+)$, then the equation (3.1) has three distinct real roots (two distinct real roots; one real and a couple of complex roots). $\mathcal{D}_F = \mathcal{F}^+ \cup \mathcal{F}^-$, where $\mathcal{F}^\pm = \{v: v_1 = F^\pm(v_2, v_3), v_3 + \frac{1}{3}v_2^2 \ge 0\}$,

(3.8)
$$F^{\pm}(v_2, v_3) = -\frac{1}{3}(v_2v_3 + \frac{2}{9}v_2^3) \pm \frac{2}{\sqrt{27}}(v_3 + \frac{1}{3}v_2^2)^{3/2}$$

The functions F^+ , F^- are smooth on $P_F^+ = \{v: v_3 + \frac{1}{3}v_2^2 > 0\}$, but only C^1 on $P_F^0 = \{v: v_3 + \frac{1}{3}v_2^2 = 0\}$. $F^{\pm}(v_2, 0) = \frac{2}{27}(-v_2^3 \pm |v_2|^2)$ and therefore $F^+(v_2, 0) = 0$ for $v_2 \ge 0$, $F^+(v_2, 0) = -\frac{4}{27}v_2^3 > 0$ for $v_2 < 0$, $F^-(v_2, 0) = -\frac{4}{27}v_2^3 < 0$ for $v_2 > 0$. $F^-(v_2, 0) = 0$ for $v_2 \le 0$. Since

$$\frac{\partial F^{\pm}(\nu_2, \nu_3)}{\partial \nu_3} = -\frac{1}{3}\nu_2 \pm \frac{1}{\sqrt{3}}(\nu_3 + \frac{1}{3}\nu_2^2)^{1/2}$$

we have

$$\frac{\partial F^{\pm}(v_2,0)}{\partial v_3} = \frac{1}{3}(-v_2 \pm |v_2|)$$



Fig. 2.

and therefore

$$\frac{\partial F^+(v_2,0)}{\partial v_3} = 0 \quad \text{for} \quad v_2 \ge 0 , \quad \frac{\partial F^-(v_2,0)}{\partial v_3} = 0 \quad \text{for} \quad v_2 \le 0 ,$$
$$\frac{\partial^2 F^{\pm}(v_2,0)}{\partial v_3^2} = \pm \frac{1}{2|v_2|} .$$

Moreover, it is obvious that $F^+(v_2, v_3) = F^-(v_2, v_3)$ if and only if $v_3 = -\frac{1}{3}v_2^2$. The above properties of the function F^+ , F^- enable us to sketch the picture of the set \mathcal{D}_F (Fig. 2). Since $\varrho_F(S_2) = \tilde{S}_2 = \{v: v_1 > 0\}$, we are interested in the restriction of \mathcal{D}_F , \mathcal{D}_F^- , \mathcal{D}_F^+ to this set only.

Obviously $\tilde{Z}_2 = \varrho_F(Z_2) = \{v: v_1 = 0, v_3 = 0\}$ and $\tilde{Z}_{1c} = \varrho_F(Z_{1c}) = \{v: v_1 = 0, v_3 < 0\}$.

Now we are interested in such $v \in \mathscr{D}_F^+$ for which the equation (3.6) has a couple of pure imaginary roots. For $v \in \mathscr{D}_F^+$ there is one real root $\lambda_1 = u + v + \frac{1}{3}v_2$ and a couple of complex ones $\lambda_{2,3} = \frac{1}{3}v_2 - \frac{1}{2}(u+v) \pm i(\sqrt{3}/2)(u-v)$, where u = $= (-q + (D_F)^{1/2})^{1/3}$, $v = (-q - (D_F)^{1/2})^{1/3}$, q, \mathscr{D}_F are as above. This implies that Re $\lambda_{2,3} = 0$ if and only if $v \in I_F = \{v \in \mathscr{D}_F^+ : H_F(v_1, v_2, v_3) = 0\}$, where $H_F(v_1, v_2, v_3) = 2v_2 - 3((-q + (D_F)^{1/2})^{1/3} + (-q - (D_F)^{1/2})^{1/3})$. For any $v_3^0 < 0$ we have $H_F(0, 0, v_3^0) = 0$. The function H_F is C^1 in a neighbourhood of the point $(0, 0, v_3^0)$ and $\partial H_F(0, 0, v_3^0)/\partial v_1 = 3/v_3^0$. Therefore there is a C^1 -function $v_1 =$ $= h(v_2, v_3)$ defined in a neighbourhood V of $(0, v_3^0)$ such that $h(0, v_3^0) = 0$ and $H_F(h(v_2, v_3), v_2, v_3) = 0$ in V. Moreover, $\partial h(0, v_3^0)/\partial v_2 = -v_3^0 > 0$ and hence the function $h(v_2, v_3^0)$ increases near the point $v_2 = 0$. We have

$$\frac{\partial H_F}{\partial v_1} = -(D_F)^{-1/2} \left(-\frac{\partial q}{\partial v_1} (D_F)^{1/2} + \frac{1}{2} \frac{\partial D_F}{\partial v_1} \right) ((D_F)^{1/2} - q)^{-3/2} - \left(\frac{\partial q}{\partial v_1} (D_F)^{1/2} + \frac{1}{2} \frac{\partial D_F}{\partial v_1} \right) ((D_F)^{1/2} + q)^{-3/2} = \\ = -\frac{1}{2} (D_F)^{-1/2} (((D_F)^{1/2} - q)^{1/3} + ((D_F)^{1/2} + q)^{1/3}) \neq 0 \quad \text{for} \quad v \in \mathcal{D}_F^+.$$

Therefore the set I_F is a two-dimensional C^1 -manifold defined not only locally near the set \tilde{Z}_{1c} . We can express the set $I_F \setminus \{v: v_1 = v_3 = 0, v_2 > 0\}$ as the graph of a C^1 -function $v_1 = h(v_2, v_3), v_3 < 0, v_2 \ge 0$. Since $H_F(0, v_2, 0) = 0, \partial H_F(0, v_2, 0)/\partial v_1 \ne 0$ for any $v_2 > 0$, the uniqueness of the implicit function implies that $\lim_{v_3 \to 0} h(v_2, v_3) = 0$. Defining $h(v_2, 0) = 0$, we obtain that I_F is the graph of a function $v_1 = h(v_2, v_3)$ defined for all $v_2 \ge 0$, $v_3 \le 0$, which is C^1 on $\{v: v_2 > 0, v_3 \le 0\}$. The boundary of the set I_F is $\{v: v_1 = 0, v_2 = 0, v_3 \le 0\} \cup \{v: v_1 = 0, v_3 = 0, v_2 \ge 0\}$.

Since for $v \in \mathcal{D}_F$ the equation (3.6) has one root of multiplicity two, it has no complex root and therefore the surface I_F does not intersect the surface \mathcal{D}_F .

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We prove that

$$\alpha = \lim_{v_3 \to 0} \frac{\partial h(v_2, v_3)}{\partial v_3} < 0 \quad \text{for any} \quad v_2 > 0$$

sufficiently small. We have

$$\frac{\partial h(v_2, v_3)}{\partial v_3} = -\left(\frac{\partial H_F(h')}{\partial v_3}\right) \left(\frac{\partial H_F(h')}{\partial v_1}\right)^{-1},$$

where $h' = (h(v_2, v_3), v_2, v_3)$ and $\partial H_F / \partial v_3 = -(D_F)^{-1/2} (-\partial q / \partial v_3 (((D_F)^{1/2} - q)^{1/3} + ((D_F)^{1/2} + q)^{1/3}) - \frac{1}{2} (((D_F)^{1/2} + q)^{2/3} - ((D_F)^{1/2} - q)^{2/3}))$. Using the above formulae for $\partial H_F / \partial v_1$ and $\partial H_F / \partial v_3$ we obtain

$$\alpha = -2 \lim_{v_3 \to 0} \left(- \frac{\partial q(h(v_2, v_3), v_2, v_3)}{\partial v_3} - \frac{1}{2} (((D_F)^{1/2} + q)^{1/3} - ((D_F)^{1/2} - q)^{1/3}) \right).$$

Since $h(v_2, 0) = 0$, $D_F(0, v_2, 0) = 0$,

$$q = \left(-\frac{v_2}{3^3}\right)^3, \quad \frac{\partial q(0, v_2, 0)}{\partial v_3} = -\frac{1}{6}v_2,$$

we obtain that $\alpha = -\frac{1}{3}\nu_2 < 0$ for $\nu_2 > 0$. This together with the fact that the set $I_F \cap \mathcal{D}_F$ is empty implies that I_F looks like in Fig. 2.

New let us consider the critical point $G = (\xi_2, 0, 0)$. Similarly to the case of the critical point F, we introduce new coordinates via the mapping

(3.9)
$$\begin{aligned} \varkappa_1 &= -(\mu_2^2 - 4\mu_1)^{1/2} ,\\ \varrho_G &: \varkappa_2 &= -\frac{1}{2}(\mu_2 + (\mu_2^2 - 4\mu_1)^{1/2} Q_2(\xi_2, \mu) ,\\ \varkappa_3 &= -\mu_3 - \frac{1}{2}(\mu_2 + (\mu_2^2 - 4\mu_1)^{1/2}) Q_1(\xi_2, \mu) . \end{aligned}$$

The mapping ρ_G is a smooth diffeomorphism on S_2 and

(3.10)
$$\hat{S}_{1} = \varrho_{G}(S_{1}) = \partial \varrho_{G}(S_{12}) = \{ \varkappa = (\varkappa_{1}, \varkappa_{2}, \varkappa_{3}) : \varkappa_{1} = 0 \},$$
$$\mu_{1} = \frac{1}{4}(\varphi_{2}^{2}(\varkappa) - \varkappa_{1}^{2}) = \varphi_{1}(\varkappa),$$
$$\varrho_{G}^{-1} : \mu_{2} = \varphi_{2}(\varkappa),$$
$$\mu_{3} = \varphi_{3}(\varkappa),$$

where the functions φ_2 , φ_3 satisfy the identities $\varkappa_2 = -\frac{1}{2}(\varphi_2(\varkappa) + \varkappa_1) Q_2(-\frac{1}{2}(\varphi_2(\varkappa) + \varkappa_1)) \varphi_1(-\frac{1}{2}(\varphi_2(\varkappa) + \varkappa_1)) Q_1(-\frac{1}{2}(\varphi_2(\varkappa) + \varkappa_1)) \varphi_1(-\frac{1}{2}(\varphi_2(\varkappa) + \varkappa_1)) \varphi_1(\varphi_2(\varkappa) + \varphi_2))$, φ_3 . Since $Q_1(0) = \omega_1 \neq 0$, $Q_2(0) = \omega_2 \neq 0$, the existence of the functions $\varphi_2, \varphi_3 \in C^\infty$ follows from the implicit function theorem. From these identities we obtain

$$\frac{\partial \varphi_2(0)}{\partial \varkappa_1} = -1 , \quad \frac{\partial \varphi_2(0)}{\partial \varkappa_2} = -\frac{2}{\omega_2} , \quad \frac{\partial \varphi_2(0)}{\partial \varkappa_3} = 0 ,$$
$$\frac{\partial \varphi_3(0)}{\partial \varkappa_1} = 0 , \quad \frac{\partial \varphi_3(0)}{\partial \varkappa_2} = -\frac{1}{\omega_2} , \quad \frac{\partial \varphi_3(0)}{\partial \varkappa_3} = \frac{1}{\omega_1} ,$$

and therefore

(3.11)
$$\mu_{1} = \frac{1}{4} \left(-\varkappa_{1} - \frac{2}{\omega_{2}} \varkappa_{2} - \varkappa_{1}^{2} + h_{2}(\varkappa) \right)$$
$$\varrho_{G}^{-1} \colon \mu_{2} = -\varkappa_{1} - \frac{2}{\omega_{2}} \varkappa_{2} + h_{2}(\varkappa) ,$$
$$\mu_{3} = -\frac{1}{\omega_{2}} \varkappa_{2} + \frac{1}{\omega_{1}} \varkappa_{3} + h_{3}(\varkappa) ,$$

where $h_2(\varkappa)$, $h_3(\varkappa) = o(\|\varkappa\|)$.

Hence we obtain

$$v_1 = -\varkappa_1,$$

$$H = \varrho_F \circ \varrho_G^{-1}: v_2 = \varkappa_2 + \tilde{h}_2(\varkappa),$$

$$v_3 = \varkappa_3 + \tilde{h}_3(\varkappa),$$

where \tilde{h}_2 , $\tilde{h}_3 = o(\|\varkappa\|)$. Since $H(0, \varkappa_2, \varkappa_3) = (0, \varkappa_2, \varkappa_3)$, we have $\tilde{h}_i(\varkappa) = \varkappa_1 \tilde{H}_i(\varkappa)$, i = 1, 2. The inverse mapping H^{-1} has the same form as H, i.e.

$$\begin{aligned} \varkappa_1 &= -\nu_1 , \\ H^{-1}: \varkappa_2 &= \nu_2 + \nu_1 H_2(\nu) , \\ \varkappa_3 &= \nu_3 + \nu_1 H_3(\nu) , \end{aligned}$$

where $H_i(v) = O(||v||)$, i = 1, 2. Therefore the characteristic equation of the matrix L(G) has the form

(3.12)
$$\lambda^{3} - (v_{2} + v_{1} H_{2}(v)) \lambda^{2} - (v_{3} + v_{1} H_{3}(v)) \lambda + v_{1} = 0.$$

The discriminant of this equation is $D_G = D_G(v) = \tilde{p}^3 + \tilde{q}^2$, where $\tilde{p} = \tilde{p}(v) = \hat{p}(H^{-1}(v))$, $\tilde{q} = \tilde{q}(v) = \hat{g}(H^{-1}(v))$, $\hat{p} = -\frac{1}{3}(\varkappa_3 + \frac{1}{3}\varkappa_2^2)$, $\hat{g} = -\frac{1}{2}(\varkappa_1 + \frac{1}{3}\varkappa_2\varkappa_3 + \frac{2}{27}\varkappa_3^2)$. Let $\mathcal{D}_G = \{v: D_G(v) = 0\}$, $\mathcal{D}_G^+ = \{v: D_G(v) > 0\}$, $\mathcal{D}_G^- = \{v: D_G(v) < 0\}$.

In the \varkappa -coordinates we have the same bifurcation diagram as we have obtained for the critical point F in the v-coordinates. In order to obtain the bifurcation diagram not only for F and G separately, but also for F and G as a couple, we need to sketch the bifurcation diagram for G also in the v-coordinates.

From the form of the mapping H it follows that H maps the \varkappa_3 -axis onto the ν_3 -axis onto the ν_3 -axis, the \varkappa_2 -axis onto the ν_2 -axis and the \varkappa_1 -axis is mapped by H onto a curve, which has its tangent at the origin close to the ν_1 -axis.

The discriminant surface \mathscr{D}_G has the form $\mathscr{D}_G = H^+ \cup H^-$, with $H^{\pm} = \{v: v_1 = \tilde{F}^{\pm}(v_2, v_3), v_1 \ge 0\}$, where $\tilde{F}^{\pm}(v_2, v_3)$ is the solution of the implicit equation

$$v_1 + F^{\pm}(v_2 + v_1H_2(v_2, v_3), v_3 + v_1H_3(v_2, v_3)) = 0.$$

From this equation, the uniqueness of its solutions and from the properties of the functions F^+ , F^- mentioned above it follows that the functions \tilde{F}^+ , \tilde{F}^- have the following properties:

$$\begin{split} \tilde{F}^{+}(v_{2},0) &= 0 \text{ for } v_{2} \geq 0, \qquad \tilde{F}^{+}(v_{2},0) < 0 \text{ for } v_{2} < 0, \\ \tilde{F}^{-}(v_{2},0) > 0 \text{ for } v_{2} > 0, \qquad \tilde{F}^{-}(v_{2},0) = 0 \text{ for } v_{2} \leq 0, \\ \frac{\partial \tilde{F}^{+}(v_{2},0)}{\partial v_{3}} &= 0 \text{ for } v_{2} \geq 0, \qquad \frac{\partial \tilde{F}^{+}(v_{2},0)}{\partial v_{3}} < 0 \text{ for } v_{2} < 0, \\ \frac{\partial \tilde{F}^{-}(v_{2},0)}{\partial v_{3}} > 0 \text{ for } v_{2} > 0, \qquad \frac{\partial \tilde{F}^{-}(v_{2},0)}{\partial v_{3}} = 0 \text{ for } v_{2} \leq 0, \\ \frac{\partial^{2} \tilde{F}^{+}(v_{2},0)}{\partial v_{3}^{2}} < 0 \text{ for } v_{2} > 0, \qquad \frac{\partial \tilde{F}^{-}(v_{2},0)}{\partial v_{3}} = 0 \text{ for } v_{2} \leq 0, \end{split}$$

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for $v_2 < 0$. The properties of the functions \tilde{F}^+ , \tilde{F}^- are the same as for the functions $-F^+$ and $-F^-$, respectively. From these properties we obtain that the surface \mathcal{D}_G looks like in Fig. 3.



Fig. 3.

Now we are interested in such $v \in \mathcal{D}_{G}^{+}$ for which the characteristic equation of the matrix L(G) has a couple of pure imaginary eigenvalues. For $v \in \mathcal{D}_{G}^{+}$ the equation (3.12) has one real root $\beta_{1} = U + V + \frac{1}{3}\kappa_{2}$ and a couple of complex ones $\beta_{2,3} = \frac{1}{3}\kappa_{2} - \frac{1}{2}(U + V) \pm i(\sqrt{3}/2)(U - V)$, where $\kappa_{2} = v_{2} + v_{1}H_{2}(v)$, $U = (-\tilde{q} + (D_{G})^{1/2})^{1/3}$, $V = (\tilde{q} - (D_{G})^{1/2})^{1/3}$. This implies that Re $\beta_{2,3} = 0$ if and only if $v \in I_{G} = \{v \in \mathcal{D}_{G}^{+} : H_{G}(v_{1}, v_{2}, v_{3}) = 0\}$, where $H_{G}(v_{1}, v_{2}, v_{3}) = 2\kappa_{2} - 3((-\tilde{q} + (D_{G})^{1/2})^{1/3} + (-\tilde{q} + (D_{G})^{1/2})^{1/3})$. For any $v_{3}^{0} < 0$ we have $H_{G}(0, 0, v_{3}^{0}) = 0$. The function H_{G} is C^{1} in a neighbourhood of the point $(0, 0, v_{3}^{0})$, and $\partial H_{G}(0, 0, v_{3}^{0})/\partial v_{1} = -v_{3}^{0} \neq 0$. Therefore there is a C^{1} -function $v_{1} = k(v_{2}, v_{3})$ defined in a neigbourhood of the point $(0, v_{3}^{0})/\partial v_{2} = v_{3}^{0} < 0$ for $v_{3}^{0} < 0$. Similarly to the case of the set I_{F} , it is possible to extend the tunction $v_{1} = k(v_{2}, v_{3})$ to a function \tilde{k} defined on the set $\{v: v_{2} \leq 0, v_{3} \leq 0\}$ so that $\tilde{k} \in C^{1}$ on $\{v: v_{2} \leq 0, v_{3} < 0\}$, $\tilde{k}(v_{2}, 0) = 0$ for $v_{2} \leq 0$, $\tilde{k}(0, v_{3}) = 0$ for $v_{3} \leq 0$ and $I_{G} = graph k$. Moreover,

$$\lim_{v_3\to 0} \frac{\partial k(v_2, v_3)}{\partial v_3} < 0 \quad \text{for any} \quad v_2 < 0 \; .$$

Similarly to the case of the set I_F , it is possible to show that the surface I_G does not intersect the surface \mathcal{D}_G . We have shown that I_G looks like in Fig. 3.

For $v_1 \in \tilde{S}_1$ there is only one critical point K, for which the matrix L(K) has the eigenvalues $\lambda_1 = 0, \lambda_{2,3} = \frac{1}{2}(v_2 \pm (v_2^2 + 4v_3)^{1/2})$. The sets $\tilde{Z}_2, \tilde{Z}_{1c} \subset \tilde{S}_1$ (see Lemma 8)



Fig. 4.

and $R = \{v \in \tilde{S}_1: v_2^2 + 4v_3 = 0\}$ divide the set \tilde{S}_1 into the following components:

$$D_{1} = \{ v \in \tilde{S}_{1} : \Psi(v_{2}, v_{3}) = v_{2}^{2} + 4v_{3} < 0, v_{2} < 0 \},$$

$$D_{2} = \{ v \in \tilde{S}_{1} : \Psi > 0, v_{2} < 0, v_{3} < 0 \}, \quad D_{3} = \{ v \in \tilde{S} : v_{3} > 0 \},$$

$$D_{4} = \{ v \in \tilde{S}_{1} : \Psi > 0, v_{2} > 0, v_{3} < 0 \}, \quad D_{5} = \{ v \in \tilde{S}_{1} : \Psi < 0, v_{2} > 0 \}$$

(see Fig. 4).

We have the following list of signs of eigenvalues of the matrix L(K):

$$\begin{split} \lambda_1 &= 0 \text{ for all } \nu \in \tilde{S}_1 \text{ and} \\ D_1 \colon \operatorname{Re} \lambda_{2,3} < 0 \,, \quad D_2 \colon \lambda_2 < 0 \,, \quad \lambda_3 < 0 \,, \quad D_3 \colon \lambda_2 > 0 \,, \quad \lambda_3 < 0 \,, \\ D_4 \colon \lambda_2 > 0 \,, \quad \lambda_3 > 0 \,, \quad D_5 \colon \operatorname{Re} \lambda_{2,3} > 0 \,, \\ Z_2^- \colon \lambda_2 &= 0 \,, \quad \lambda_3 < 0 \,, \quad Z_2^+ \colon \lambda_3 = 0 \,, \quad \lambda_2 > 0 \,, \quad \tilde{Z}_{1c} \colon \lambda_{2,3} = \pm i\omega \,, \end{split}$$

 $\omega \neq 0$, where $\tilde{Z}_2 = Z_2^+ \cup Z_2^-$, $Z_2^+ = \{ \nu \in \tilde{Z}_2 : \nu_2 > 0 \}$, $Z_2^- = \{ \nu \in \tilde{Z}_2 : \nu_2 < 0 \}$.

Let us introduce the following notations: $\mathscr{D}_1 = \mathscr{D}_F^- \cap \mathscr{D}_G^-$, $\mathscr{D}_2 = \mathscr{D}_F^+ \cap \mathscr{D}_G^-$, $\mathscr{D}_3 = \mathscr{D}_F^- \cap \mathscr{D}_G^+$; $I_F^+(I_G^+)(I_F^-(I_G^-))$ is the set of all $v \in \mathscr{D}_F^+(\mathscr{D}_G^+)$ for which the matrix



Fig. 5.

L(F)(L(G)) has a couple of complex eigenvalues with positive (negative) real parts,

$$\begin{split} I_1 &= I_F^- \cap I_G^-, \quad I_2 = I_F^+ \cap I_G^+, \quad I_3 = I_F^- \cap I_G^+, \quad J_1 = I_1 \cap \mathcal{D}_2, \\ J_2 &= I_1 \cap \mathcal{D}_1, \quad J_3 = I_1 \cap \mathcal{D}_3, \quad K_1 = I_2 \cap \mathcal{D}_3, \quad K_2 = I_2 \cap \mathcal{D}_1, \\ K_3 &= I_1 \cap \mathcal{D}_2, \\ A_1 &= \mathcal{D}_F \cap \mathcal{D}_G \cap \{ v: v_2 < 0, v_3 < 0 \}, \\ A_2 &= \mathcal{D}_F \cap \mathcal{D}_G \cap \{ v: v_2 < 0, v_3 > 0 \}, \\ A_3 &= \mathcal{D}_F \cap \mathcal{D}_G \cap \{ v: v_2 < 0, v_3 > 0 \}, \\ B_1 &= \mathcal{D}_F \cap I_G, \quad B_2 = \mathcal{D}_G \cap I_F \quad (\text{see Fig. 5}). \end{split}$$

If the matrix L(F)(L(G)) has only real eigenvalues, then we denote them by $\lambda_1, \lambda_2, \lambda_3(\beta_1, \beta_2, \beta_3)$. If $v \in \mathcal{D}_F^+(\mathcal{D}_G^+)$, then the matrix L(F)(L(G)) has one real and a couple of complex eigenvalues. Let us denote the real eigenvalue by $\lambda_2(\beta_3)$ and the complex one by $\lambda(\beta)$. Then det $L(F) = \lambda_2 |\lambda|^2 = v_1 > 0$, det $L(G) = \beta_3 |\beta|^2 = -v_1 < 0$ and therefore $\lambda_2 > 0, \beta_3 < 0$. Since $\lambda_2 + 2$ Re $\lambda = v_2$ and $\beta_3 + 2$ Re $\beta = v_2 + v_1 H_2(v)$ we obtain that Re $\lambda < 0$ for $v_2 < 0$ and Re $\beta > 0$ for $v_2 > 0, v_1$ sufficiently small. These properties of the eigenvalues together with the list of signs of eigenvalues of the matrix L(K) for $v \in \tilde{S}_1$ enable us to deduce the following list of signs of eigenvalues for $v_1 > 0$:

$$\begin{split} I_1: & \operatorname{Re} \lambda < 0, \quad \lambda_2 > 0, \quad \operatorname{Re} \beta < 0, \quad \beta_3 < 0, \\ I_2: & \operatorname{Re} \lambda > 0, \quad \lambda_2 > 0, \quad \operatorname{Re} \beta > 0, \quad \beta_3 < 0, \\ \Im_1: & \lambda_1 < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \beta_1 > 0, \quad \beta_2 > 0, \quad \beta_3 < 0, \\ \varnothing_1: & \lambda_1 < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \beta_1 > 0, \quad \beta_2 > 0, \quad \beta_3 < 0, \\ \Im_2: & \operatorname{Re} \lambda < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \operatorname{Re} \beta > 0, \quad \beta_3 < 0, \\ \Im_3: & \lambda_1 < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \operatorname{Re} \beta > 0, \quad \beta_3 < 0, \\ J_1: & \lambda_1 < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \operatorname{Rr} \beta < 0, \quad \beta_3 < 0, \\ J_2: & \lambda_1 < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \operatorname{Rr} \beta < 0, \quad \beta_3 < 0, \\ J_2: & \lambda_1 < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \operatorname{Rr} \beta < 0, \quad \beta_3 < 0, \\ J_2: & \lambda_1 < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \operatorname{Rr} \beta < 0, \quad \beta_3 < 0, \\ J_3: & \operatorname{Re} \lambda < 0, \quad \lambda_2 > 0, \quad \lambda_3 < 0, \quad \operatorname{Re} \beta > 0, \quad \beta_3 < 0, \\ K_1: & \lambda_1 > 0, \quad \lambda_2 > 0, \quad \lambda_3 > 0, \quad \operatorname{Re} \beta > 0, \quad \beta_3 < 0, \\ K_2: & \lambda_1 > 0, \quad \lambda_2 > 0, \quad \lambda_3 > 0, \quad \beta_1 > 0, \quad \beta_2 > 0, \quad \beta_3 < 0, \\ K_3: & \operatorname{Re} \lambda > 0, \quad \lambda_2 > 0, \quad \beta_1 = \beta_2 < 0, \quad \beta_3 < 0, \\ A_1: & \lambda_1 = \lambda_3 < 0, \quad \lambda_2 > 0, \quad \beta_1 = \beta_2 < 0, \quad \beta_3 < 0, \\ A_2: & \lambda_1 = \lambda_3 < 0, \quad \lambda_2 > 0, \quad \beta_1 = \beta_2 > 0, \quad \beta_3 < 0, \\ A_3: & \lambda_1 = \lambda_3 < 0, \quad \lambda_2 > 0, \quad \beta_1 = \beta_2 > 0, \quad \beta_3 < 0, \\ B_1: & \lambda_1 = \lambda_3 < 0, \quad \lambda_2 > 0, \quad \beta_1 = \beta_2 > 0, \quad \beta_3 < 0, \\ B_1: & \lambda_1 = \lambda_3 < 0, \quad \lambda_2 > 0, \quad \beta_1 = \beta_2 > 0, \quad \beta_3 < 0, \\ B_2: & \lambda_{2,3} = \pm i\gamma, \quad \gamma \neq 0, \quad \lambda_2 > 0, \quad \beta_1 = \beta_2 > 0, \quad \beta_3 < 0. \\ \end{array}$$

4. **BIFURCATIONS**

In this section we study the bifurcations of the family (2.8). Although we have obtained a relatively simple bifurcation diagram for the critical points, the bifurcation diagram for the corresponding eigenvalues indicates that the bifurcations of the phase portraits are complicated.

For $\mu^0 \in Z_2$ we have $(\mu_2^0)^2 - 4\mu_1^0 = 0$, $2\mu_3^0 - \mu_2^0 = 0$. The point $K = (-\frac{1}{2}\mu_2^0, 0, 0)$ is the unique critical point of the vector field v_{μ^0} (we denote by v_{μ} the vector field corresponding to the parameter μ). Let $\xi_i = \xi_i(\mu)$, i = 1, 2, be the roots of the equation $y^2 + \mu_2 y + \mu_1 = 0$ such that $\xi_i(\mu^0) = -\frac{1}{2}\mu_2^0$ (we assume $\mu \in S_1 \cup S_2$). If $y_1 - \xi_1 = x_1$, $y_2 = x_2$, $y_3 = x_3$, then the family (2.8) becomes

(4.1)
$$\dot{x}_{1} = x_{2},$$

$$\dot{x}_{2} = x_{3},$$

$$\dot{x}_{3} = x_{1}(x_{1} + \xi_{1} - \xi_{2}) + \mu_{3}x_{2}\tilde{Q}_{1}(x_{1}, \mu) + \xi_{1}x_{2}\tilde{Q}_{1}(x_{1}, \mu) +$$

$$+ x_{1}x_{2}\tilde{Q}_{1}(x_{1}, \mu) + \xi_{1}x_{3}\tilde{Q}_{2}(x_{1}, \mu) + x_{1}x_{3}\tilde{Q}_{2}(x_{1}, \mu) +$$

$$+ x_{2}x_{3}\tilde{Q}_{3}(x_{3}, \mu) + x_{2}^{2}\tilde{\Phi}_{1}(x, \mu) + x_{3}^{2}\tilde{\Phi}_{2}(x, \mu),$$

where the functions \tilde{Q}_i , $\tilde{\Phi}_j$ have the same properties as the functions Q_i , Φ_j from (2.8). The family (4.1) has two critical points $K_1 = (0, 0, 0)$ and $K_2 = (\xi, 0, 0)$, where $\xi = \xi_2 - \xi_1$. The matrix of the linearization at K_1 is $L(K_1) = A(\mu) = (a_{ij})$, where $a_{12} = a_{23} = 1$, $a_{31} = \xi_1 - \xi_2$, $a_{32} = (\mu_3 + \xi_1) \tilde{Q}_1(0, \mu)$, $a_{33} = \xi_1 \tilde{Q}_2(0, \mu)$ and the other entries are equal to zero. For $\mu^0 \in Z_2$ also $a_{31} = a_{32} = 0$ and $a_{33} = \gamma = -\frac{1}{2}\mu_2^0 \tilde{Q}_2(0, 0, \mu_2^0, 0)$. If

$$C = \begin{pmatrix} -\gamma & 1 - \gamma & 1 \\ 0 & -\gamma & 1 \\ 0 & 0 & 1 \end{pmatrix}, \text{ then } C^{-1} = \begin{pmatrix} -\gamma^{-1} & \gamma^{-1} - \gamma^{-2} & \gamma^{-2} \\ 0 & -\gamma^{-1} & \gamma^{-1} \\ 0 & 0 & 1 \end{pmatrix}$$

and using the change of coordinates u = Cx we obtain

(4.2)
$$\begin{pmatrix} \dot{u}_1 \\ \dot{u}_2 \\ \dot{u}_3 \end{pmatrix} = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \gamma \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \\ u_3 \end{pmatrix} + B_0(\mu) \begin{pmatrix} u_1 \\ u_2 \\ u_3 \end{pmatrix} + \begin{pmatrix} F_0(u, \mu) \\ F_0(u, \mu) \\ F_0(u, \mu) \end{pmatrix},$$

where $F_0(u, \mu) = f(C^{-1}u, \mu)$, f is the nonlinear part of the right-hand side of the third equation of (4.1), $B_0(\mu^0) = 0$, $F_0(u, \mu^0) = A_{200}u_1^2 + A_{020}u_2^2 + A_{002}u_3^2 + A_{110}u_1u_2 + A_{101}u_1u_3 + A_{011}u_2u_3 + o(||u||^2)$. By [6, Theorem 2.2] (see also [4], [11] the parametrized central manifold can be expressed as the graph of a function $u_3 = h(u_1, u_2, \mu)$ defined locally, in a neighbourhood of the point $(0, 0, \mu^0)$ for which

$$h(0, 0, \mu^{0}) = \frac{\partial h(0, 0, \mu^{0})}{\partial u_{1}} = \frac{\partial h(0, 0, \mu^{0})}{\partial u_{2}} = 0$$

Therefore the reduction of the family (4.2) to the central manifold has the form

(4.3)
$$\begin{pmatrix} \dot{u}_1 \\ \dot{u}_2 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} + B(\mu) \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} + \begin{pmatrix} F(u, \mu) \\ F(u, \mu) \end{pmatrix},$$

where $B(\mu^0) = 0$, $F(u, \mu^0) = A_{200}u_1^2 + A_{110}u_1u_2 + A_{020}u_2^2 + o(||u||^2)$, $A_{200} = \gamma^{-2}$, $A_{110} = -2\gamma^{-1}(\gamma^{-1} - \gamma^{-2}) + \gamma^{-2}Q_1(0, \mu^0)$.

Let us restrict the set of parameters to a neighbourhood $U(\mu^0) \subset P(\mu^0)$ of the point μ^0 , where $P(\mu^0)$ is a two-dimensional surface crossing the set Z_2 transversally at μ^0 . Using Bogdanov's method (see [3]) it is possible to rewrite the family (4.3) in suitable coordinates on $U(\mu^0)$, $\varepsilon = \varrho(\mu)$, $v = \delta(u)$, $\varrho(\mu^0) = 0$, $\delta(0) = 0$ to the form

(4.4)
$$\dot{v}_1 = v_2$$
,
 $\dot{v}_2 = \varepsilon_1 + \varepsilon_2 v_1 + g(v, \varepsilon)$

where $g(v, 0) = (Qv, v) + o(||v||^2)$, $Q = (q_{ij})$ is a symmetric matrix with $q_{11} \neq 0$. By [3, Lemma 2] $q_{12} = \tilde{q}_{12} \cdot \hat{g}_{11}^{-1}$, where $\tilde{q}_{11} = \gamma^{-2}$, $\tilde{q}_{12} = -2\gamma^{-1}(\gamma^{-1} - \gamma^{-2}) + \gamma^{-2}Q_1(0, \mu^0)$. Therefore sign $q_{12} = -\text{sign } \mu_2^0 \omega_2$ for μ_2^0 sufficiently small.

Denote by $v_{\varepsilon}^+(v_{\varepsilon}^-)$ the family (4.4) with $q_{12} > 0$ ($q_{12} < 0$). We remark that it is possible to transform the family v_{ε}^- to the same form with $q_{12} > 0$ by using the change of coordinates $x_2 \to -x_2$, $t \to -t$. The complete bifurcation diagram for the family v_{ε}^+ is described in [1, 3].

Now it is convenient to use the v-coordinales (see (3.4)). Since $v_2 = -\frac{1}{2}(\mu_2 - \nu_1)$. $Q_2(\xi_1, \mu)$, we have that $q_{12} > 0$ ($q_{12} < 0$) for $v^0 = (0, v_2^0, 0) \in Z_2^+(Z_2^-)$. This means that the bifurcations near $v^0 \in Z_2^+(Z_2^-)$ correspond to the bifurcations of the family $v_{\epsilon}^+(v_{\epsilon}^-)$.



Assume $v^0 \in \mathbb{Z}_2^-$. For the family v_{ε}^- there exists a curve R (see [3]), on which a stable focus bifurcates into a stable closed orbit and the focus becomes unstable. By the bifurcation diagram shown in Fig. 5, this Hopf bifurcation may occur only near



the point G. For $v \in I_1(I_3)$ the matrix L(G) has one real eigenvalue $\beta_3 < 0$ and a couple of complex eigenvalues β , $\overline{\beta}$ with Re $\beta < 0$ (Re $\beta > 0$). This means that if the parameter goes in the direction $I_1 \rightarrow I_3$, crossing the surface I_G transversally, then the stable focus G bifurcates into a stable closed orbit and the focus becomes unstable. This determines the orientation of Bogdanov's bifurcation cycle. By [3] there must be a curve P in $U(\mu^0) \cap I_3$ with the end-point at ν^0 such that if the parameter v approaches this curve, the period of the closed orbit tends to infinity, i.e. the closed orbit bifurcates into a homoclinic orbit. This implies that for the family (2.8) (in the v-coordinates) there is a surface $S_G \subset I_3 \cap \mathcal{D}_3$ such that if the parameter v approaches this surface, the period of the closed orbit, arising on I_G , tends to infinity. Since for a parameter from the set Z_{1c} the corresponding central manifold is three-dimensional, the two-dimensional central manifold corresponding to a parameter from the set Z_2^- is destroyed if the parameter passes out of a neighbourhood of the set Z_2^- . Therefore the global properties of the surface S_G cannot be found by the methods of plane vector fields and so it is difficult to find them. We know the form of S_G near Z_2^- .

If $v^0 \in \mathbb{Z}_2^+$, then by the bifurcation diagram, the Hopf bifurcation may occur near the point F only. For the family v_{ε}^+ there exists a curve, on which an unstable focus



Fig. 9.

bifurcates into an unstable closed orbit. For $v \in K_3(I_3)$, the matrix L(F) has one real eigenvalue $\lambda_2 > 0$ and a couple of complex eigenvalues λ , λ with Re $\lambda > 0$ (Re $\lambda < 0$). This means that if the parameter v goes in the direction $K_3 \rightarrow I_3$, crossing the surface I_F transversally, an unstable focus bifurcates into an unstable closed orbit and this determines the orientation of Bogdanov's cycle. Similarly as above, Bogdanov's results imply that there must be a surface $S_F \subset I_3 \cap \mathcal{D}_2$ on which the closed orbit arising on I_F bifurcates into a homoclinic orbit. The problem of global properties of S_F remains open.

From the above considerations we conclude that in a neighbourhood of $v^0 \in \tilde{Z}_2$ the bifurcation diagram and the bifurcations look like in Figures 6-9.

We have described the bifurcations near the set \tilde{Z}_2 . For the results bifurcations near the set \tilde{Z}_{1c} we refer to the papers [5], [7–10]. The problem how the phase portraits appearing for the parameter from a neighbourhood of \tilde{Z}_2 may bifurcate into different phase portraits corresponding to the values of the parameters from a neighbourhood of the set \tilde{Z}_{1c} remains open.

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