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# ASYMPTOTIC BEHAVIOUR OF THE SOLUTIONS OF A CERTAIN TYPE OF THE THIRD ORDER DIFFERENTIAL EQUATIONS

# JANA FEŤKOVÁ

In [1] M. Greguš among other things investigates the asymptotic properties of solutions without zero-points and also asymptotic properties of oscillatory solutions of the third order linear differential equation in the normal form

(A<sub>0</sub>) 
$$y'''(\tau) + 2A(\tau)y'(\tau) + [A'(\tau) + b(\tau)]y(\tau) = 0.$$

The functions A, A',  $b \in C^0(\langle \tau_0, \beta \rangle)$ ,  $-\infty \leq \tau_0 < \beta \leq +\infty$ . He simultaneously investigates the adjoint differential equation to the equation (A<sub>0</sub>) in the form

(B<sub>0</sub>) 
$$z'''(\tau) + 2A(\tau)z'(\tau) + [A'(\tau) - b(\tau)]z(\tau) = 0.$$

The function  $b(\tau)$  (Laguerre invariant) is supposed to have the property (V):  $b(\tau) \ge 0$  for all  $\tau \in \langle \tau_0, \beta \rangle$  and  $b(\tau) \equiv 0$  does not hold in any interval.

In this paper we investigate asymptotic behaviour of solutions without zero-points and the asymptotic behaviour of the oscillatory solutions of a more general third order linear ordinary differential equation of the form

(A<sub>1</sub>) 
$$[r(t)[r(t)x'(t)]']' + 2a_1(t)x'(t) + [a_1'(t) + b_1(t)]x(t) = 0, t \ge t_0,$$

where r(t) > 0,  $rx' \in C^{0}(\langle t_{0}, \infty \rangle)$ ,  $[rx']' \in C^{0}(\langle t_{0}, \infty \rangle)$ ,  $[r[rx']'] \in C^{0}(\langle t_{0}, \infty \rangle)$ ;  $a_{1} \in C^{1}(\langle t_{0}, \infty \rangle)$ ,  $b_{1} \in C^{0}(\langle t_{0}, \infty \rangle)$ ,  $-\infty < t_{0} < +\infty$ .

The adjoint equation to  $(A_1)$  is of the form

(B<sub>1</sub>) 
$$[r(t)[r(t)\xi'(t)]']' + 2a_1(t)\xi'(t) + [a_1'(t) - b_1(t)]\xi(t) = 0.$$

Two cases will be considered:

I. (1)  $\int_{-\infty}^{\infty} dt/r(t) = \infty \text{ and}$ II. (2)  $\int_{-\infty}^{\infty} dt/r(t) < \infty.$  In both cases the solutions of  $(A_1)$  will be transformed into solutions of equations which will have the form  $(A_0)$ , and conversely. On the basis of these transformations we shall apply the results from [1] and [2] for the solutions of  $(A_1)$ . The idea to apply the transformations stems from the paper [3] by Philos.

I. If 
$$\int_{-\infty}^{\infty} dt/r(t) = \infty$$
, then we denote

(3) 
$$R(t) = \int_{t_0}^t \mathrm{d}s/r(s) \quad \text{for } t \ge t_0.$$

The function  $R \in C^4(\langle t_0, \infty \rangle)$ , is increasing and maps the interval  $\langle t_0, \infty \rangle$  onto the interval  $\langle 0, \infty \rangle$ . Its inverse function  $R^{-1}(\tau)$  is increasing on  $\langle 0, \infty \rangle$ , and the latter will be mapped onto the interval  $\langle t_0, \infty \rangle$ .

**Lemma I.** Suppose that (1) holds. Let x(t) be a solution of  $(A_1)$  in the interval  $\langle t_0, \infty \rangle$ . Then the function

(4) 
$$y(\tau) = x[R^{-1}(\tau)] \quad \text{for all } \tau \in \langle 0, \infty \rangle$$

is a solution of

$$(A_{01}) y''' + 2a_1[R^{-1}(\tau)]y' + [a_1'[R^{-1}(\tau)] + b_1[R^{-1}(\tau)]]r[R^{-1}(\tau)] y = 0$$

in  $\langle 0, \infty \rangle$ .

Conversely, if  $y(\tau)$  is a solution of  $(A_{01})$  in  $\langle 0, \infty \rangle$ , and the function x(t) is determined by relation (4), i.e.

(4,1) 
$$x(t) = y[R(t)], \quad t \in \langle t_0, \infty \rangle,$$

then the function (4,1) satisfies the equation (A<sub>1</sub>) in  $\langle t_0, \infty \rangle$ .

Proof. Differentiating the relation (4), and considering R'(t) = 1/r(t) for all  $t \in \langle t_0, \infty \rangle$ , we obtain the equalities

(5) 
$$y'(\tau) = r(t) x'(t)$$

(6)  $y''(\tau) = r(t) [r(t) x'(t)]'$ 

(7) 
$$y'''(\tau) = r(t)[r(t)[r(t)x'(t)]']', \quad \tau = R(t), \ \tau \in \langle 0, \infty \rangle,$$

for the function,  $y(\tau)$  given by (4). From the equalities (5)—(7) and from  $t = R^{-1}(\tau)$  it follows on the basis of (A<sub>1</sub>) that

$$y'''(\tau) = -2a_1(t)r(t)x'(t) - [a_1'(t) + b_1(t)]r(t)x(t) =$$
  
= -2a\_1[R^{-1}(\tau)]y' - [a\_1'[R^{-1}(\tau)] + b\_1[R^{-1}(\tau)]]r[R^{-1}(\tau)]y,

and so  $y(\tau)$  is a solution of  $(A_{01})$  from the corresponding interval.

The relation (4) means the bijective mapping of the space of solutions of  $(A_1)$ .

onto the space of solutions of  $(A_{01})$ . Indeed, for arbitrary values  $y_0^{(i)}$ , i = 0, 1, 2 the system of conditions

$$y_0 = x(t_0)$$
  

$$y'_0 = r(t_0) x'(t_0)$$
  

$$y''_0 = r(t_0) [r(t_0) x''(t_0) + r'(t_0) x'(t_0)]$$

has the only solution in the variables  $x(t_0)$ ,  $x'(t_0)$ ,  $x''(t_0)$ . From that the second part of the assertion of Lemma I follows.

Remark 1. Between the solutions of the adjoint equation

(B<sub>1</sub>) 
$$[r(t)[r(t)\xi'(t)]']' + 2a_1(t)\xi'(t) + [a_1'(t) - b_1(t)]\xi(t) = 0$$

and those of the equation

(B<sub>01</sub>) 
$$z''' + 2a_1[R^{-1}(\tau)]z' + [a'_1[R^{-1}(\tau)] - b_1[R^{-1}(\tau)]]r[R^{-1}(\tau)]z = 0$$

there is a similar relation as for the equations (A<sub>1</sub>) and (A<sub>01</sub>). If  $\xi(t)$  is the solution of (B<sub>1</sub>) in  $\langle t_0, \infty \rangle$ , then the function

(4,2) 
$$z(\tau) = \xi[R^{-1}(\tau)] \text{ for all } \tau \in \langle 0, \infty \rangle$$

is the solution of  $(B_{01})$ .

Remark 2. If we put

(8) 
$$A_{01}(\tau) = a_1[R^{-1}(\tau)]; \quad b_{01}(\tau) = b_1[R^{-1}(\tau)]r[R^{-1}(\tau)]$$

for all  $\tau \in \langle 0, \infty \rangle$ , so on the basis of the equality

$$[a_1[R^{-1}(\tau)]]' = a_1'[R^{-1}(\tau)]r[R^{-1}(\tau)]$$

the equations  $(A_{01})$  and  $(B_{01})$  can be written in the form

$$y''' + 2A_{01}(\tau)y' + [A'_{01}(\tau) + b_{01}(\tau)]y = 0,$$
  
$$z''' + 2A_{01}(\tau)z' + [A'_{01}(\tau) - b_{01}(\tau)]z = 0.$$

II. If  $\int_{-\infty}^{\infty} dt/r(t) < \infty$ , then we define the function

(9) 
$$\varrho(t) = \int_{t_0}^{\infty} \mathrm{d}s/r(s), \quad t \ge t_0$$

which belongs to the class  $C^4(\langle t_0, \infty \rangle)$ , is decreasing and maps the interval  $\langle t_0, \infty \rangle$  onto the interval  $(0, \varrho(t_0))$ . Denote  $\varrho(t_0) = \varrho_0$ . Its inverse function  $\varrho^{-1}$  is decreasing in the interval  $(0, \varrho_0)$ , and it maps this interval onto the interval  $\langle t_0, \infty \rangle$ . Then a composite function  $\varrho^{-1}(1/\tau)$  and  $\varrho^{-1}(e^{-\tau})$  belongs to the class  $C^4(\langle \tau_0, \infty \rangle)$ , where  $\tau_0 = 1/\varrho_0$  and  $\tau_0 = -\ln \varrho_0$  respectively, is increasing in this interval, and maps the interval  $\langle \tau_0, \infty \rangle$  onto  $\langle t_0, \infty \rangle$ .

**Lemma II.** Assume that (2) holds. Let x(t) be the solution of  $(A_1)$  in the interval  $\langle t_0, \infty \rangle$ . Then the function

(10) 
$$y(\tau) = \tau^2 x [\varrho^{-1}(1/\tau)] \quad \text{for all } \tau \in \langle \tau_0, \infty \rangle$$

is the solution of  $(A_{02})$  in the interval  $\langle t_0, \infty \rangle$ , where

(A<sub>02</sub>) 
$$y''' + 2\tau^{-4}a_1[\varrho^{-1}(1/\tau)]y' + \tau^{-6}\{r[\varrho^{-1}(1/\tau)][a_1'[\varrho^{-1}(1/\tau)] + b_1[\varrho^{-1}(1/\tau)]] - 4\tau a_1[\varrho^{-1}(1/\tau)]\}y = 0.$$

Conversely, if  $y(\tau)$  is the solution of equation  $(A_{02})$  in  $\langle \tau_0, \infty \rangle$ , then the function x(t) determined by the relation (10), i.e.

(10,1) 
$$x(t) = \varrho^2(t) y[\varrho^{-1}(t)], \quad t \in \langle t_0, \infty \rangle$$

satisfies (A<sub>1</sub>) in  $\langle t_0, \infty \rangle$ .

The proof can be obtained in a similar way as in the case of Lemma I. Remark 3. By relation

(10,2) 
$$z(\tau) = \tau^2 \xi[\varrho^{-1}(1/\tau)]$$

the space of the solutions of  $(\mathbf{B}_1)$  is transformed bijectively onto the space of the solutions of

(B<sub>02</sub>) 
$$z'''(\tau) + 2\tau^{-4}a_1[\varrho^{-1}(1/\tau)]z' + \tau^{-6}\{r[\varrho^{-1}(1/\tau)][a'_1[\varrho^{-1}(1/\tau)] - b_1[\varrho^{-1}(1/\tau)] - 4\tau a_1[\varrho^{-1}(1/\tau)]\}z = 0.$$

Remark 4. If we denote

(11) 
$$A_{02}(\tau) = \tau^{-4} a_1[\varrho^{-1}(1/\tau)]; \quad b_{02}(\tau) = \tau^{-6} b_1[\varrho^{-1}(1/\tau)] r[\varrho^{-1}(1/\tau)],$$

so on the basis of equality

$$A'_{02}(\tau) = -4\tau^{-5}a_1[\varrho^{-1}(1/\tau)] + \tau^{-6}a'_1[\varrho^{-1}(1/\tau)]r[\varrho^{-1}(1/\tau)]$$

the equations  $(A_{02})$  and  $(B_{02})$  can be written in the form

$$y''' + 2A_{02}(\tau)y' + [A'_{02}(\tau) + b_{02}(\tau)]y = 0,$$
  
$$z''' + 2A_{02}(\tau)z' + [A'_{02}(\tau) - b_{02}(\tau)]z = 0.$$

**Lemma III.** Suppose that (2) is satisfied. Let x(t) be the solution of  $(A_1)$  in  $\langle t_0, \infty \rangle$ . Then the function

(12) 
$$y(\tau) = e^{\tau} x[\varrho^{-1}(e^{-\tau})] \quad for \ all \ \tau \ge \tau_0 = -\ln \varrho(t_0)$$

is the solution of  $(A_{03})$  in  $\langle \tau_0, \infty \rangle$ , where

(A<sub>03</sub>) 
$$y''' + 2[e^{-2r}a_1[\varrho^{-1}(e^{-r})] - 1/2]y' + e^{-3r}\{r[\varrho^{-1}(e^{-r})][a'_1[\varrho^{-1}(e^{-r})] + b_1[\varrho^{-1}(e^{-r})] - 2e^{-2r}a_1[\varrho^{-1}(e^{-r})]\}y = 0.$$

Conversely, if  $y(\tau)$  is the solution of  $(A_{03})$  in  $\langle \tau_0, \infty \rangle$ , and x(t) is determined by (12), i.e.

(12,1) 
$$x(t) = \varrho(t) y[-\ln \varrho(t)] \quad t \in \langle t_0, \infty \rangle,$$

then the function (12,1) satisfies an equation (A<sub>1</sub>) in  $\langle t_0, \infty \rangle$ .

The proof can be obtained in a similar way as in Lemma I.

Remark 5. By relation

(12,2) 
$$z(\tau) = e^{-\tau} \xi[\varrho^{-1}(e^{-\tau})]$$

the space of solutions of  $(B_1)$  is transformed bijectively onto the space of solutions of

(B<sub>03</sub>) 
$$z''' + 2[e^{-2r}a_1[\varrho^{-1}(e^{-r})] - 1/2]z' + e^{-3r}\{[a'_1[\varrho^{-1}(e^{-r})] - b_1[\varrho^{-1}(e^{-r})]]r[\varrho^{-1}(e^{-r})] - 2e^ra_1[\varrho^{-1}(e^{-r})]\}z = 0.$$

Remark 6. By designating

$$A_{03}(\tau) = e^{-2\tau} a_1[\varrho^{-1}(e^{-\tau})] - 1/2;$$
  

$$b_{03}(\tau) = e^{-3\tau} b_1[\varrho^{-1}(e^{-\tau})] r[\varrho^{-1}(e^{-\tau})]$$

and from the equality

$$A'_{03}(\tau) = e^{-3\tau} r[\varrho^{-1}(e^{-\tau})] a'_1[\varrho^{-1}(e^{-\tau})] - 2e^{-2\tau} a_1[\varrho^{-1}(e^{-\tau})]$$

the equations  $(A_{03})$  and  $(B_{03})$  can be written in the form

$$y''' + 2A_{03}(\tau) y' + [A'_{03}(\tau) + b_{03}(\tau)] y = 0$$
  
$$z''' + 2A_{03}(\tau) z' + [A'_{03}(\tau) - b_{03}(\tau)] z = 0.$$

Next we shall present the results, which by using Lemma I and some lemmas and theorems proved by M. Greguš, for the solutions of  $(A_1)$ .

**Lemma 1.1.** Let (1) hold. Let  $a_1(t) \leq 0$ ,  $a'_1(t) + b_1(t) \geq 0$  in  $\langle t_0, \infty \rangle$  and let the function  $b_1(t)$  have the property

(V<sub>1</sub>):  $b_1(t) \ge 0$  for all  $t \in \langle t_0, \infty \rangle$  while  $b_1(t) \ne 0$ in any subinterval.

Let  $0 < r(t) \leq M$  in  $\langle t_0, \infty \rangle$ . Let  $\xi$  be the solution of the differential equation (B<sub>1</sub>) with  $\xi(\alpha_1) = \xi'(\alpha_1) = 0$ ,  $\xi''(\alpha_1) > 0$  for  $\alpha_1 \in \langle t_0, \infty \rangle$ . Then  $\xi(t) \to \infty$ ,  $\xi'(t) \to \infty$ as  $t \to \infty$ .

Proof. If we transform the equation  $(A_1)$  into  $(A_{01})$  by using Lemma I., we can see that  $(A_{01})$  satisfies all assumptions of Lemma 3.1 in [2, p. 120] and moreover, in this transformation [relation (4,1)] there corresponds to the solution  $\xi(t)$  of  $(B_1)$  with a double zero-point  $\alpha_1$  the solution  $z(\tau)$  of  $(B_{01})$  having a double zero-point in  $\tau_1 = R[\alpha_1]$ . According to Lemma 3.1 in [2]  $z(\tau) \to \infty$ ,  $z'(\tau) \to \infty$  as  $\tau \to \infty$ . From (4,1), (1) and (4,2) we have:  $\xi(t) = z[R(t)] \to \infty$ ,  $\xi'(t) = r^{-1}(t) z'[R(t)] \to \infty$  as  $t \to \infty$ , because  $1/r(t) \ge 1/M > 0$  for all  $t \ge t_0$ . Lemma 1.2. Let the hypotheses of Lemma 1.1 be satisfied and, moreover, let

(13) 
$$\int_{t_0}^{\infty} R^2(t) [a_1'(t) + b_1(t)] dt = \infty.$$

Then the solution  $\xi$  mentioned in Lemma 1.1 satisfies also

(14) 
$$r(t)[r(t)\xi'(t)]' + 2a_1(t)\xi(t) \to \infty \quad \text{as } t \to \infty.$$

Proof. With respect to both Lemma 1.1, from which we take the notation, and Lemma 3.2 in [2, p. 120] it is sufficient to prove that the condition

$$\int_0^\infty \tau^2 [a_1'[R^{-1}(\tau)] + b_1[R^{-1}(\tau)]] r[R^{-1}(\tau)] d\tau = \infty$$

is equivalent to (13) and the property of the solution z of  $(\mathbf{B}_{01}) \lim_{\tau \to \infty} [z''(\tau) + 2A_{01}(\tau)z(\tau)] = \infty$  means the relation (14). However,

$$\int_0^\infty \tau^2 [a_1'[R^{-1}(\tau)] + b_1[R^{-1}(\tau)]] r[R^{-1}(\tau)] d\tau = \int_{t_0}^\infty R^2(t) [a_1'(t) + b_1(t)] dt$$

and on the basis of (4,1), (4,2) we get the statement of Lemma 1.2.

By using (3), (4), (5), (6) and the monotonicity of R(t) we obtain from Theorem 4 in [1] the following:

**Theorem 1.1.** Let the assumptions of Lemma 1.2 be satisfied in  $\langle t_0, \infty \rangle$ . Then there exists exactly one solution x of the differential equation (A<sub>1</sub>) [up to linear dependence] with the following properties:  $x(t) \neq 0$ , sgn  $x(t) = \text{sgn}[r(t)x'(t)]' \neq$  $\neq \text{sgn } x'(t)$  for  $t \in \langle t_0, \infty \rangle$ ; x(t), r(t)x'(t), r(t)[r(t)x'(t)]' are monotonic functions of  $t \in \langle t_0, \infty \rangle$  and  $x(t) \to 0$ ,  $r(t)x'(t) \to 0$ ,  $r(t)[r(t)x'(t)]' \to 0$  as  $t \to \infty$ .

From Theorem 5 in [1] there follows

**Theorem 1.2.** Let  $a_1(t) \leq 0$ ,  $a'_1(t) + b_1(t) \geq 0$  and  $b_1(t)$  have the property  $(V_1)$ for  $t \in \langle t_0, \infty \rangle$ . If the differential equation  $(A_1)$  has an oscillatory solution on  $\langle t_0, \infty \rangle$ , then all solutions of  $(A_1)$  are oscillatory on  $\langle t_0, \infty \rangle$  with one exception of the solution x [up to linear dependence] with the following properties:  $x(t) \neq 0$ ,  $\operatorname{sgn} x(t) = \operatorname{sgn} [r(t)x'(t)]' \neq \operatorname{sgn} x'(t)$  for  $t \in \langle t_0, \infty \rangle$ ; x(t), r(t)x'(t), r(t)[r(t)x'(t)]'are monotonic functions for  $t \in \langle t_0, \infty \rangle$  and  $r(t)x'(t) \to 0$ ,  $r(t)[r(t)x'(t)]' \to 0$  as  $t \to \infty$ .

Example. Consider the equation

(15)

$$[t^{x}[t^{x}x'(t)]']' + \kappa t^{2(\kappa-1)}x'(t) + [\kappa(\kappa-1) + (\kappa-2)(\kappa-3)]t^{2\kappa-3}x(t) = 0$$

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for  $t \ge t_0 > 0$ .  $\varkappa$  is a non-positive number. As easily verified, all the hypotheses of Theorem 1.2 are fulfilled, and the solution x(t) = 1/t of (15) has all the properties satisfying the conclusions of Theorem 1.2.

Since on the basis of (3) and (8),

$$\int_0^\infty b_{01}(\tau) \, \mathrm{d}\tau = \int_{t_0}^\infty b_1(t) \, \mathrm{d}t,$$

from Theorem 3.4 in [2, p. 122] it follows

**Theorem 1.3.** Let  $b_1(t)$  have the property  $(V_1)$  in  $\langle t_0, \infty \rangle$  and let

$$\int_{t_0}^{\infty} b_1(t) \,\mathrm{d}t = \infty.$$

Then the equation (A<sub>1</sub>) has at least one solution x with no zeros in  $\langle t_0, \infty \rangle$  and satisfying  $\liminf x(t) = 0$ .

By using Lemma II we can obtain the following results for the equation  $(A_1)$ .

**Lemma 2.1.** Let (2) hold. Let  $a_1(t) \leq 0$ ,  $\varrho(t)r(t)(a'_1(t) + b_1(t)) \geq 4a_1(t)$  for  $t \in \langle t_0, \infty \rangle$ . Let  $b_1(t)$  have the property  $(V_1)$ . Let  $\xi$  be the solution of the equation (B<sub>1</sub>) with the properties:  $\xi(\alpha_2) = \xi'(\alpha_2) = 0$ ,  $\xi''(\alpha_2) > 0$ ,  $\alpha_2 \in \langle t_0, \infty \rangle$ . Then  $\varrho^{-2}(t)\xi(t) \to \infty$ ,  $[r(t)\xi'(t) + 2\varrho^{-1}(t)\xi(t)] \to \infty$  as  $t \to \infty$ , where  $\varrho(t)$  is determined by (9).

Lemma 2.2. Let the assumptions of Lemma 2.1 be fulfilled and, moreover,

$$\int_{t_0}^{\infty} \varrho^2(t) [a_1'(t) + b_1(t) - 4 \varrho^{-1}(t) r^{-1}(t) a_1(t)] dt = \infty.$$

Then the solution  $\xi$  mentioned in Lemma 2.1 satisfies also  $[\varrho^2(t) r(t) [r(t) \xi'(t)]' + 2\varrho(t) r(t) \xi'(t) + 2(1 + \varrho^2(t) a_1(t)) \xi(t)] \rightarrow \infty$  as  $t \rightarrow \infty$ .

Further, we can prove a theorem where we use the formula (9) and derivatives of the relation (10).

**Theorem 2.1.** Let the assumptions of Lemma 2.2 be satisfied in  $\langle t_0, \infty \rangle$ . Then there exists exactly one solution x of  $(A_1)$  [up to linear dependence] with the following properties:  $x(t) \neq 0$ ,  $\operatorname{sgn} x(t) = \operatorname{sgn} [\varrho^2(t)r(t)[r(t)x'(t)]' + 2\varrho(t)r(t) \cdot x'(t) + 2x(t)] \neq \operatorname{sgn} [r(t)x'(t) + 2\varrho^{-1}(t)x(t)]$  for  $t \in \langle t_0, \infty \rangle$ ;  $\varrho^{-2}(t)x(t), r(t) \cdot x'(t) + 2\varrho^{-1}(t)x(t), \varrho^2(t)r(t)[r(t)x'(t)]' + 2\varrho(t)r(t)x'(t) + 2x(t)$  are monotonic functions and  $\varrho^{-2}(t)x(t) \to 0$ ,  $[r(t)x'(t) + 2\varrho^{-1}(t)x(t)] \to 0$ ,  $[\varrho^2(t)r(t) \cdot [r(t)x'(t)]' + 2\varrho(t)x'(t)r(t) + 2x(t)] \to 0$  as  $t \to \infty$ .

**Theorem 2.2.** Let (2) hold. Let  $a_1(t) \leq 0$ ,  $\varrho(t)r(t)[a'_1(t) + b_1(t)] \geq 4a_1(t)$  for  $t \in \langle t_0, \infty \rangle$  and  $b_1(t)$  have the property (V<sub>1</sub>). If the differential equation (A<sub>1</sub>) has an oscillatory solution on  $\langle t_0, \infty \rangle$ , then all solutions of (A<sub>1</sub>) are oscillatory with one

exception of the solution x [up to linear dependence] with the following properties:  $x(t) \neq 0$ ,  $\operatorname{sgn} x(t) = \operatorname{sgn} [\varrho^2(t) r(t) [r(t) x'(t)]' + 2\varrho(t) r(t) x'(t) + 2x(t)] \neq \operatorname{sgn} \cdot [r(t) x'(t) + 2\varrho^{-1}(t) x(t)]$  for  $t \in \langle t_0, \infty \rangle$ ;  $\varrho^{-2}(t) x(t)$ ,  $r(t) x'(t) + 2\varrho^{-1}(t) x(t)$ ,  $\varrho^2(t) r(t) [r(t) x'(t)]' + 2\varrho(t) r(t) x'(t) + 2x(t)$  are monotonic functions on  $\langle t_0, \infty \rangle$ and  $[r(t) x'(t) + 2\varrho^{-1}(t) x(t)] \rightarrow 0$ ,  $[\varrho^2(t) r(t) [r(t) x'(t)]' + 2\varrho(t) r(t) x'(t) + 2x(t)] \rightarrow 0$  as  $t \rightarrow \infty$ .

**Theorem 2.3.** Let  $b_1(t)$  have the property  $(V_1)$  in  $\langle t_0, \infty \rangle$  and let

$$\int_{t_0}^{\infty} \varrho^4(t) \, b_1(t) \, r^{-1}(t) \, \mathrm{d}t = \infty.$$

Then the equation (A<sub>1</sub>) has at least one solution x with no zeros in  $\langle t_0, \infty \rangle$  and satisfying  $\liminf_{t \to \infty} \varrho^{-2}(t) x(t) = 0$ .

Further assertions can be derived by using Lemma III.

**Lemma 3.1.** Suppose that (2) holds and  $\varrho(t)$  is determined by (9). Let  $\varrho^2(t) a_1(t) \leq 1/2$ ,  $\varrho(t) r(t) [a'_1(t) + b_1(t)] \geq 2a_1(t)$  and let  $b_1(t)$  have the property  $(V_1)$  in  $\langle t_0, \infty \rangle$ . Let  $\xi$  be such a solution of  $(B_1)$  that  $\xi(\alpha_3) = \xi'(\alpha_3) = 0$ ,  $\xi''(\alpha_3) > 0$ ,  $\alpha_3 \in \langle t_0, \infty \rangle$ . Then  $\varrho^{-1}(t) \xi(t) \to \infty$ ,  $[r(t) \xi'(t) + \varrho^{-1}(t) \xi(t)] \to \infty$  as  $t \to \infty$ .

Lemma 3.2. Let the assumptions of Lemma 3.1 be satisfied and, moreover, let

$$\int_{t_0}^{\infty} \left[ \ln \varrho(t) \right]^2 \left[ \varrho^2(t) \left[ a_1'(t) + b_1(t) \right] - 2\varrho(t) r^{-1}(t) a_1(t) \right] dt$$

diverge. Then the solution  $\xi$  mentioned in Lemma 3.1 satisfies also  $[\varrho(t)r(t) \cdot [r(t)\xi'(t)]' + r(t)\xi(t) + 2\varrho(t)a_1(t)\xi(t)] \rightarrow \infty$  as  $t \rightarrow \infty$ .

**Theorem 3.1.** Let the hypotheses of Lemma 3.2 be fulfilled in  $\langle t_0, \infty \rangle$ . Then there exists exactly one solution x of (A<sub>1</sub>) [up to linear dependence] with the following properties:  $x(t) \neq 0$  for  $t \in \langle t_0, \infty \rangle$ ,  $\varrho^{-1}(t) x(t), r(t) x'(t) + \varrho^{-1}(t) x(t)$ ,  $\varrho(t) r(t) [r(t) x'(t)]' + r(t) x'(t) + \varrho^{-1}(t) x(t)$  are monotonic functions in  $\langle t_0, \infty \rangle$ ;  $\operatorname{sgn} x(t) = \operatorname{sgn} [\varrho(t) r(t) [r(t) x'(t)]' + r(t) x'(t) + \varrho^{-1}(t) x(t)] \neq \operatorname{sgn} [r(t) x'(t) +$  $<math>+ \varrho^{-1}(t) x(t)]$  for  $t \in \langle t_0, \infty \rangle$  and  $\varrho^{-1}(t) x(t) \to 0$ ,  $[r(t) x'(t) + \varrho^{-1}(t) x(t)] \to 0$ ,  $\varrho(t) r(t) [r(t) x'(t)]' \to 0$  as  $t \to \infty$ .

**Theorem 3.2.** Let (2) hold and let  $\varrho^2(t) a_1(t) \leq 1/2$ ,  $\varrho(t)r(t)[a_1'(t) + b_1(t)] \geq 2a_1(t)$  and let  $b_1(t)$  have the property  $(V_1)$  in  $\langle t_0, \infty \rangle$ . If the differential equation  $(A_1)$  has an oscillatory solution in  $\langle t_0, \infty \rangle$ , then all solutions of  $(A_1)$  are oscillatory with one exception of the solution x [up to the linear dependence] with the following properties:  $x(t) \neq 0$ ,  $\operatorname{sgn} x(t) = \operatorname{sgn} [\varrho(t)r(t)[r(t)x'(t)]' + r(t)x'(t)' + \varrho^{-1}(t) \cdot x(t)] \neq \operatorname{sgn} [r(t)x'(t) + \varrho^{-1}(t)x(t)]$  for  $t \in \langle t_0, \infty \rangle$ ;  $x(t) \varrho^{-1}(t), r(t)x'(t) + \varrho^{-1}(t)x(t), \varrho(t)r(t)[r(t)x'(t)]' + r(t)x'(t) + \varrho^{-1}(t)x(t)$  are monotonic functions and  $[r(t)x'(t) + \varrho^{-1}(t)x(t)] \rightarrow 0, \varrho(t)r(t)[r(t)x'(t)]' \rightarrow 0$  as  $t \to \infty$ .

**Theorem 3.3.** Let  $b_1(t)$  have the property  $(V_1)$  and let

$$\int_{t_0}^{\infty} \varrho^2(t) b_1(t) \,\mathrm{d}t = \infty.$$

Then the differential equation (A<sub>1</sub>) has at least one solution x with no zeros in  $\langle t_0, \infty \rangle$  and satisfying  $\liminf \rho^{-1}(t) x(t) = 0$ .

Finally the following theorem which generalizes Theorem 1.15 in [2, p. 20] can be proved by using the Lemmas I—III.

**Theorem 4.** If the function  $b_1(t) \ge 0$  for  $t \in \langle t_0, \infty \rangle$ , then there exists a solution of the equation (A<sub>1</sub>) without zeros in  $\langle t_0, \infty \rangle$ .

Example. In the equation

$$[t[tx'(t)]']' + (1 - 3\varkappa^2)\ln^{-2}(t)x'(t) + (t^{-1}\ln^{-3}(t))(3\varkappa^2 + 2\varkappa^3 - 1)x(t) = 0,$$

where  $\varkappa \ge 0$  is a constant, the Laguerre invariant  $b_1(t) = 2(t^{-1}\ln^{-3}(t))\varkappa^3$  is nonnegative on an interval  $\langle t_0, \infty \rangle$ ,  $t_0 > e$ . This equation has by Theorem 4 a solution without zeros. Such solutions, for example, are  $x_1(t) = \ln^{1+\varkappa}(t)$ ,  $x_2(t) = \ln^{1+\varkappa}(t) \cdot \ln(\ln t)$ .

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## АСИМПТОТИЧЕСКИЕ СВОЙСТВА РЕШЕНИЙ ОПРЕДЕЛЕННОГО ТИПА ДИФФЕРЕНЦИАЛЬНЫХ УРАВНЕНИЙ ТРЕТЬЕГО ПОРЯДКА

### Jana Feťková

#### Резюме

В работах М. Грегуша [1] и [2] приведены асимптотические свойства решений без нулевых точек и также асимптотические свойства осциллирующих решений дифференциального уравнения третьего порядка в форме (A<sub>0</sub>). В этой статье исследованы выше приведенные свойства для дифференциального уравнения в форме (A<sub>1</sub>).

Предполагается, что условия (1) или (2) выполены. В обоих случаях преобразуются решения уравнения ( $A_1$ ) в решения уравнения ( $A_{0i}$ ) i = 1, 2, 3 (которые в форме ( $A_0$ )) и наоборот. На основании преобразования применяются результаты работ [1] и [2] к решению уравнения ( $A_1$ ).

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