## Mathematic Bohemica

Ivan Chajda<br>Ring-like structures derived from $\lambda$-lattices with antitone involutions

Mathematica Bohemica, Vol. 132 (2007), No. 1, 87-96

Persistent URL: http: //dml.cz/dmlcz/133992

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# RING-LIKE STRUCTURES DERIVED FROM $\lambda$-LATTICES WITH ANTITONE INVOLUTIONS 

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(Received December 12, 2005)


#### Abstract

Using the concept of the $\lambda$-lattice introduced recently by V. Snášel we define $\lambda$-lattices with antitone involutions. For them we establish a correspondence to ring-like structures similarly as it was done for ortholattices and pseudorings, for Boolean algebras and Boolean rings or for lattices with an antitone involution and the so-called Boolean quasirings.


Keywords: $\lambda$-lattice, $\lambda$-semilattice, ortholattice, $\lambda$-ortholattice, antitone involution, Boolean quasiring

MSC 2000: 06C15, 16Y99, 06B99, 81P10, 06A12

The well-known correspondence between Boolean algebras and Boolean rings (see e.g. [1]) was extended to orthomodular lattices by H. Länger [12] and to ortholattices by the author in [3]. A general setting was described by G. Eigenthaler, H. Länger and the author in [4] and [5]. It was generalized to lattices with antitone involution by D. Dorninger, H. Länger and M. Mączyński [6], [7] and to generalized orthomodular lattices, see [5]. It was motivated by the use of these ring-like structures in certain logics of quantum mechanics, see e.g. [6], [7] for the description. However, in quantum mechanics it can happen that we cannot distinguish between two possibilities. It leads us to study more general structures than lattices where still ring-like structures can be induced.

Suitable tools for our investigation thus can be the so-called $\lambda$-lattices introduced by V. Snášel [13] and treated in [11], [13] and the so-called $\lambda$-semilattices known also under the name commutative directoids in [9]. On the other hand, this level of generality can make problems in computing formulas because some well-known

This work is supported by the Czech Goverment via the research project MSM 6198959214.
results for lattices fail for $\lambda$-lattices. For example, if $x \mapsto x^{\prime}$ is an antitone involution on a lattice $\mathcal{L}=(L ; \vee, \wedge)$ then the De Morgan laws

$$
\begin{equation*}
(x \vee y)^{\prime}=x^{\prime} \wedge y^{\prime} \quad \text { and } \quad(x \wedge y)^{\prime}=x^{\prime} \vee y^{\prime} \tag{DM}
\end{equation*}
$$

hold in $\mathcal{L}$, see e.g. [1], [8]. However, this is not the case for $\lambda$-lattices. Hence we must often require some additional properties.

## 1. $\lambda$-LATtices with antitone involutions

First, we recall some well-known concepts. Let $(A ; \leqslant)$ be an ordered set. For $a, b \in A$ denote

$$
\begin{aligned}
U(a, b) & =\{x \in A ; a \leqslant x \text { and } b \leqslant x\}, \\
L(a, b) & =\{x \in A ; x \leqslant a \text { and } x \leqslant b\} .
\end{aligned}
$$

An ordered set $(A ; \leqslant)$ is called up-directed (down-directed) if $U(a, b) \neq \emptyset$ (or $L(a, b) \neq \emptyset$, respectively) for each $a, b \in A$. Further, $(A ; \leqslant)$ is directed if it is both an up- and down-directed set.

Let $(A ; \leqslant)$ be a down-directed set. Denote by $\operatorname{Exp} A$ the power set of $A$. Let $\lambda$ be a mapping $\lambda: \operatorname{Exp} A \rightarrow A$ such that
(i) $\lambda(L(a, b)) \in L(a, b)$,
(ii) if $a \leqslant b$ then $\lambda(L(a, b))=a$.

Define a binary operation $\wedge$ on $A$ as follows:

$$
a \wedge b=\lambda(L(a, b))
$$

The groupoid $(A ; \wedge)$ will be called a $\lambda$-semilattice (or a commutative directoid in [9]). It is easy to verify that $(A ; \wedge)$ satisfies the identities
(I) $x \wedge x=x$ (idempotency),
(C) $x \wedge y=y \wedge x$ (commutativity),
(SA) $x \wedge((x \wedge y) \wedge z)=(x \wedge y) \wedge z$ (skew associativity).
Also conversely, if $(A ; \wedge)$ is a groupoid satisfying (I), (C), (SA) and $\leqslant$ is defined by the rule
$a \leqslant b \quad$ if and only if $\quad a \wedge b=a$
then $(A ; \leqslant)$ is a down-directed set and $\lambda(L(a, b))=a \wedge b$ satisfies (i) and (ii) mentioned above.

Moreover, if $(A ; \leqslant)$ is a directed set and $\lambda: \operatorname{Exp} A \rightarrow A$ satisfies also
(iii) $\lambda(U(a, b)) \in U(a, b)$,
(iv) if $a \leqslant b$ then $\lambda(U(a, b))=b$
then we can introduce another operation $\vee$ by setting

$$
a \vee b=\lambda(U(a, b))
$$

and easily verify the following identities:
( $\left.\mathrm{I}^{\prime}\right) x \vee x=x$,
$\left(\mathrm{C}^{\prime}\right) x \vee y=y \vee x$,
$\left(\mathrm{SA}^{\prime}\right) x \vee((x \vee y) \vee z)=(x \vee y) \vee z$,
(Ab) $x \vee(x \wedge y)=x, x \wedge(x \vee y)=x$ (absorption).
Then the algebra $(A ; \vee, \wedge)$ is called a $\lambda$-lattice (see [11], [13]). Also conversely, if $(A ; \vee, \wedge)$ is an algebra of type $(2,2)$ satisfying the identities $(\mathrm{I}),\left(\mathrm{I}^{\prime}\right),(\mathrm{C}),\left(\mathrm{C}^{\prime}\right),(\mathrm{SA})$, $\left(\mathrm{SA}^{\prime}\right),(\mathrm{Ab})$ then the relation defined by
$a \leqslant b \quad$ if and only if $\quad a \vee b=b$
coincides with the already introduced induced order on $(A ; \wedge)$. If we put

$$
\lambda(U(a, b))=a \vee b
$$

then (iii) and (iv) are satisfied.
Hence, $\lambda$-semilattices and $\lambda$-lattices can be viewed either as algebras satisfying certain identities or as ordered sets with constrains on upper and lower bounds. Contrary to the case of lattices, the choice of $\lambda(L(a, b))$ or $\lambda(U(a, b))$ need not be unique. Consider e.g. the directed sets drawn in Fig. 1 and Fig. 2.


Fig. 1


Fig. 2

In the first case, we have three choices for $a \vee b$, namely $c$ or $d$ or 1 . Analogously, $c \wedge d$ can be $a$ or $b$ or 0 (for comparable elements the choice is unique due to the conditions (ii) and (iv)). In all 9 possible cases, the resulting algebra will be a $\lambda$ lattice. Analogously, in the second case (which is a lattice) we can pick up e.g. $a \vee b=c$ or $a \vee b=1$. In the latter case, the resulting algebra will be a $\lambda$-lattice which is not a lattice.

Let $(A ; \leqslant)$ be an ordered set. A mapping $x \mapsto x^{\prime}$ on $A$ is called an antitone involution if $x^{\prime \prime}=x$ and $x \leqslant y$ implies $y^{\prime} \leqslant x^{\prime}$. It is well-known that if $(L ; \vee, \wedge)$ is
a lattice and $x \mapsto x^{\prime}$ is an antitone involution on $(L ; \leqslant)$ where $\leqslant$ is the induced order then the De Morgan laws (DM) hold.

Consider our $\lambda$-lattice $\mathcal{L}=(A ; \vee, \wedge)$ visualized in Fig. 1. Define $x \mapsto x^{\prime}$ on $A=$ $\{0, a, b, c, d, 1\}$ by the table

| $x$ | 0 | $a$ | $b$ | $c$ | $d$ | 1 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $x^{\prime}$ | 1 | $d$ | $c$ | $b$ | $a$ | 0 |.

Then clearly it is an antitone involution on $A$. Let us pick up e.g. $a \vee b=c$ (for other elements the operation $\vee$ is determined as supremum). Then $(A ; \vee)$ is a $\lambda$ semilattice (w.r.t. $\vee$ ). Now, if e.g. $c \wedge d=a$ (for other elements it is determined as infimum) then $\mathcal{L}=(A ; \vee, \wedge)$ is a $\lambda$-lattice with an antitone involution but

$$
b^{\prime} \wedge a^{\prime}=c \wedge d=a \neq b=c^{\prime}=(b \vee a)^{\prime}
$$

thus in $\mathcal{L}$ the De Morgan laws do not hold. On the contrary, when choosing $c \wedge d=b$, De Morgan laws hold in $\mathcal{L}$. Hence, validity of (DM) depends not only on the induced order but also on our choice of operations.

## 2. $\lambda$-Boolean quasirings

By a Boolean quasiring (see [4], [6], [7]) we mean an algebra $\mathcal{R}=(R ;+, \cdot, 0,1)$ of type $(2,2,0,0)$ satisfying the identities

$$
\begin{align*}
& x+y=y+x,  \tag{R1}\\
& x+0=x  \tag{R2}\\
& (x \cdot y) \cdot z=x \cdot(y \cdot z),  \tag{R3}\\
& x \cdot y=y \cdot x  \tag{R4}\\
& x \cdot x=x  \tag{R5}\\
& x \cdot 0=0  \tag{R6}\\
& x \cdot 1=x  \tag{R7}\\
& 1+(1+x \cdot y) \cdot(1+y)=y . \tag{R8}
\end{align*}
$$

For our purposes, we modify this definition as follows. An algebra $\mathcal{R}=$ ( $R ;+, \cdot, 0,1$ ) of type $(2,2,0,0)$ is called a $\lambda$-Boolean quasiring if it satisfies the identities (R1), (R2), (R4)-(R8) and

$$
\begin{equation*}
x \cdot((x \cdot y) \cdot z)=(x \cdot y) \cdot z \tag{*}
\end{equation*}
$$

One can immediately verify that every Boolean quasiring is a $\lambda$-Boolean quasiring since (R3*) follows easily by (R3), (R4) and (R5). The mutual correspondence between Boolean quasirings and lattices with an antitone involution was established in [6]. We are going to extend this correspondence to $\lambda$-Boolean quasirings.

Theorem 1. Let $\mathcal{L}=\left(L ; \vee, \wedge,{ }^{\prime}, 0,1\right)$ be a bounded $\lambda$-lattice with an antitone involution. Define

$$
x+y=(x \vee y) \wedge(x \wedge y)^{\prime} \quad \text { and } \quad x \cdot y=x \wedge y .
$$

Then $\mathcal{R}(L)=(L ;+, \cdot, 0,1)$ is a $\lambda$-Boolean quasiring. If, moreover, the De Morgan laws hold in $\mathcal{L}$ then $\mathcal{R}(L)$ satisfies the correspondence identity (Cor) $(1+(1+x) \cdot(1+y)) \cdot(1+x \cdot y)=x+y$.

Proof. Since $(L ; \wedge)$ is a bounded $\lambda$-semilattice, the identities (R3*), (R4)-(R7) are immediate consequences of (I), (C), (SA) and the properties of the induced order. The identity (R1) is a trivial consequence of the definition and (R2) is evident. It remains to prove (R8) and (Cor).

For (R8) we use ( Ab ) to compute

$$
1+(1+x \cdot y) \cdot(1+y)=\left((x \wedge y)^{\prime} \wedge y^{\prime}\right)^{\prime}=\left(y^{\prime}\right)^{\prime}=y^{\prime \prime}=y
$$

For (Cor), we apply De Morgan laws to derive

$$
(1+(1+x) \cdot(1+y)) \cdot(1+x \cdot y)=\left(x^{\prime} \wedge y^{\prime}\right)^{\prime} \wedge(x \wedge y)^{\prime}=(x \vee y) \wedge(x \wedge y)^{\prime}=x+y
$$

We can prove also the converse.

Theorem 2. Let $\mathcal{R}=(R ;+, \cdot, 0,1)$ be a $\lambda$-Boolean quasiring. Define

$$
x \wedge y=x \cdot y, \quad x^{\prime}=1+x \quad \text { and } \quad x \vee y=1+(1+x) \cdot(1+y)
$$

Then $\mathcal{L}(R)=\left(R ; \vee, \wedge,{ }^{\prime}, 0,1\right)$ is a bounded $\lambda$-lattice with an antitone involution in which the De Morgan laws hold.

Proof. By $\left(\mathrm{R} 3^{*}\right),(\mathrm{R} 4)-(\mathrm{R} 7),(R ; \wedge)$ is a bounded $\lambda$-semilattice. If we put $y=x$ in (R8) and apply (R5), we obtain the identity

$$
\begin{equation*}
1+(1+x)=x \tag{*}
\end{equation*}
$$

proving that the unary operation $x^{\prime}=1+x$ is an involution on $R$. Suppose $x \leqslant y$. Then $x \wedge y=x \cdot y=x$ and, by (R8),

$$
1+x^{\prime} \cdot y^{\prime}=1+(1+x) \cdot(1+y)=1+(1+x \cdot y) \cdot(1+y)=y
$$

Thus, applying (*), we arrive at

$$
x^{\prime} \wedge y^{\prime}=x^{\prime} \cdot y^{\prime}=1+\left(1+x^{\prime} \cdot y^{\prime}\right)=1+y=y^{\prime}
$$

whence $y^{\prime} \leqslant x^{\prime}$, i.e. this operation is an antitone involution on $R$.
Further, using ( $*$ ), we obtain

$$
x^{\prime} \vee y^{\prime}=1+\left(1+x^{\prime}\right) \cdot\left(1+y^{\prime}\right)=1+x \cdot y=(x \wedge y)^{\prime}
$$

and

$$
x^{\prime} \wedge y^{\prime}=(1+x) \cdot(1+y)=1+(1+(1+x) \cdot(1+y))=(x \vee y)^{\prime}
$$

thus the De Morgan laws hold in $\mathcal{L}(R)$. Due to this fact, $(R ; \vee)$ is also a $\lambda$-semilattice and, by (R8),

$$
(x \wedge y) \vee y=\left((x \cdot y)^{\prime} \cdot y^{\prime}\right)^{\prime}=1+(1+x \cdot y) \cdot(1+y)=y
$$

The dual absorption law can be established by the De Morgan laws and involutorness. Hence, $\mathcal{L}(R)=\left(R ; \vee, \wedge,^{\prime}, 0,1\right)$ is a bounded $\lambda$-lattice with an antitone involution.

Theorem 3. Let $\mathcal{R}=(R ;+, \cdot, 0,1)$ be a $\lambda$-Boolean quasiring satisfying the correspondence indentity (Cor). Then $\mathcal{R}(\mathcal{L}(R))=\mathcal{R}$.

Let $\mathcal{L}=\left(L ; \vee, \wedge,{ }^{\prime}, 0,1\right)$ be a bounded $\lambda$-lattice with an antitone involution in which De Morgan laws hold. Then $\mathcal{L}(\mathcal{R}(L))=\mathcal{L}$.

Proof. Evidently, the multiplicative operations coincide in $\mathcal{R}(\mathcal{L}(R))$ and $\mathcal{R}$. To prove $\mathcal{R}(\mathcal{L}(R))=\mathcal{R}$ we need only to show that also $\oplus=+$ where $\oplus$ is the additive operation in $\mathcal{R}(\mathcal{L}(R))$. Applying (Cor) we compute

$$
x \oplus y=(x \vee y) \wedge(x \wedge y)^{\prime}=(1+(1+x) \cdot(1+y)) \cdot(1+x \cdot y)=x+y
$$

Analogously, the operation meet clearly coincides in $\mathcal{L}(\mathcal{R}(L))$ and $\mathcal{L}$. Hence, it remains to prove $\sqcup=\vee$ and $x^{*}=x^{\prime}$ where $\sqcup$ is the join and ${ }^{*}$ is the antitone involution in $\mathcal{L}(\mathcal{R}(L))$. We have

$$
x^{*}=1+x=(1 \vee x) \wedge(1 \wedge x)^{\prime}=1 \wedge x^{\prime}=x^{\prime}
$$

and

$$
x \sqcup y=1+(1+x) \cdot(1+y)=\left(x^{\prime} \wedge y^{\prime}\right)^{\prime}=x \vee y
$$

due to the De Morgan laws.

## 3. $\lambda$-ORTHOLATTICES

By an ortholattice (see e.g. [2], [10]) we mean a bounded lattice $\mathcal{L}=\left(L ; \vee, \wedge,{ }^{\prime}, 0,1\right)$ with an antitone involution which is a complementation, i.e. $x \wedge x^{\prime}=0$ and $x \vee x^{\prime}=1$ for each $x \in L$. Of course, De Morgan laws hold in $\mathcal{L}$ and hence $x \wedge x^{\prime}=0$ is equivalent to $x \vee x^{\prime}=1$.

We can extend this concept to $\lambda$-lattices.
By a $\lambda$-ortholattice we mean a bounded $\lambda$-lattice $\mathcal{L}=\left(L ; \vee, \wedge,{ }^{\prime}, 0,1\right)$ with an antitone involution such that $x \wedge x^{\prime}=0$ and $x \vee x^{\prime}=1$.

Formally, the definition is the same as for ortholattices, but we must be careful. For example, the $\lambda$-lattice depicted in Fig. 3 is an $\lambda$-ortholattice.


Fig. 3

Since it is a $\lambda$-lattice, we must specify joins and meets of non-comparable elements. If e.g. $a \vee b=c$ and $a^{\prime} \wedge b^{\prime}=d^{\prime}$ then $a^{\prime} \wedge b^{\prime}=d^{\prime} \neq c^{\prime}=(a \vee b)^{\prime}$ and hence the De Morgan laws do not hold. It means that for $a \vee b=c$ we must set $a^{\prime} \wedge b^{\prime}=c^{\prime}$ etc.

The mutual correspondence between ortholattices and the so-called pseudosemirings was settled in [3]. However, it is easy to establish such a correspondence for $\lambda$-ortholattices and certain $\lambda$-Boolean quasirings.

Theorem 4. Let $\mathcal{L}=\left(L ; \vee, \wedge,{ }^{\prime}, 0,1\right)$ be a $\lambda$-ortholattice in which the De Morgan laws hold. Define

$$
x+y=(x \vee y) \wedge(x \wedge y)^{\prime} \quad \text { and } \quad x \cdot y=x \wedge y
$$

Then $\mathcal{R}(L)=(L ;+, \cdot, 0,1)$ is a $\lambda$-Boolean quasiring of characteristic 2 (i.e. satisfying the identity $x+x=0$ ) satisfying the correspondence identity (Cor).

Let $\mathcal{R}=(R ;+, \cdot, 0,1)$ be a $\lambda$-Boolean quasiring of characteristic 2. Define

$$
x \vee y=1+(1+x) \cdot(1+y), \quad x \wedge y=x \cdot y, \quad x^{\prime}=1+x
$$

Then $\mathcal{L}(R)=\left(R ; \vee, \wedge,^{\prime}, 0,1\right)$ is a $\lambda$-ortholattice in which the De Morgan laws hold.

Moreover, $\mathcal{L}(\mathcal{R}(L))=\mathcal{L}$ and if $\mathcal{R}$ satisfies (Cor) then also $\mathcal{R}(\mathcal{L}(R))=\mathcal{R}$.
Proof. Of course, if $x \wedge x^{\prime}=0$ then $x+x=x \wedge x^{\prime}=0$, thus the induced $\lambda$-Boolean quasiring $\mathcal{R}(L)$ is of characteristic 2 for each $\lambda$-ortholattice $\mathcal{L}$. The rest of the proof follows by Theorem 1 .

Conversely, if $\mathcal{R}$ satisfies $x+x=0$ then $x \wedge x^{\prime}=(x \vee x) \wedge(x \wedge x)^{\prime}=x+x=0$ and, due to the De Morgan laws, also $x \vee x^{\prime}=1$ and thus $\mathcal{L}(R)$ is really a $\lambda$-ortholattice.

The remaining statements follow directly by Theorems 2 and 3 .
In what follows we concentrate on a special case of a $\lambda$-ortholattice, the so-called orthomodular $\lambda$-lattice which satisfies the orthomodular law

$$
\begin{equation*}
x \leqslant y \Rightarrow x \vee\left(y \wedge x^{\prime}\right)=y \tag{OML}
\end{equation*}
$$

which is equivalent to the identity

$$
\begin{equation*}
(x \wedge y) \vee\left(x \wedge(x \wedge y)^{\prime}\right)=x \tag{**}
\end{equation*}
$$

An example of an orthomodular $\lambda$-lattice can be constructed from the lattice drawn in Fig. 4.


Fig. 4

One can easily verify that $\mathcal{L}$ is an orthomodular lattice. To change it into a $\lambda$ lattice which is not a lattice, we define $a \vee b=1$ and $a^{\prime} \wedge b^{\prime}=0$ and leave all other joins and meets as they were in $\mathcal{L}$. Of course, the new algebra is a $\lambda$-lattice but not a lattice, but the orthomodular law remains valid.

We are wonder what a $\lambda$-Boolean quasiring corresponds to an orthomodular $\lambda$ lattice.

Theorem 5. Let $\mathcal{L}=\left(L ; \vee, \wedge,{ }^{\prime}, 0,1\right)$ an orthomodular $\lambda$-lattice in which the $D e$ Morgan laws hold. Define

$$
x+y=(x \vee y) \wedge(x \wedge y)^{\prime} \quad \text { and } \quad x \cdot y=x \wedge y
$$

Then $\mathcal{R}(L)$ is a $\lambda$-Boolean quasiring of characteristic 2 satisfying the identity

$$
\begin{equation*}
(1+x \cdot y)+(x+x \cdot y)=1+x \tag{OM}
\end{equation*}
$$

Let $\mathcal{R}=(R ;+, \cdot, 0,1)$ be a $\lambda$-Boolean quasiring of characteristic 2 satisfying the identity (OM). Define

$$
x \vee y=1+(1+x) \cdot(1+y), \quad x \wedge y=x \cdot y \quad \text { and } \quad x^{\prime}=1+x
$$

Then $\mathcal{L}(R)=\left(R ; \vee, \wedge,{ }^{\prime}, 0,1\right)$ is an orthomodular $\lambda$-lattice in which the De Morgan laws hold.

Moreover, $\mathcal{L}(\mathcal{R}(L))=\mathcal{L}$ and if $\mathcal{R}$ satisfies also (Cor) then $\mathcal{R}(\mathcal{L}(R))=\mathcal{R}$.
Proof. In virtue of Theorem 4, we need only to verify the orthomodular law for $\mathcal{L}(R)$ and the identity $(\mathrm{OM})$ for $\mathcal{R}(L)$.

Let $\mathcal{L}$ be an orthomodular $\lambda$-lattice in which the De Morgan laws hold. We can easily derive $x^{\prime} \vee\left(x^{\prime} \vee y^{\prime}\right)=x^{\prime} \vee y^{\prime}$ since $x^{\prime} \leqslant x^{\prime} \vee y^{\prime}$ and thus

$$
\begin{aligned}
x \cdot(1+x \cdot y) & =x \wedge(x \wedge y)^{\prime}=x \wedge\left(x^{\prime} \vee y^{\prime}\right)=x \wedge\left(x^{\prime} \vee\left(x^{\prime} \vee y^{\prime}\right)\right) \\
& =x \wedge(x \wedge(x \wedge y))^{\prime}=(x \vee(x \wedge y)) \wedge(x \wedge(x \wedge y))^{\prime}=x+x \cdot y
\end{aligned}
$$

Since $1+x \cdot y=(x \wedge y)^{\prime} \geqslant x \wedge(x \wedge y)^{\prime}=x \cdot(1+x \cdot y)$ and for $a \geqslant b$ we have

$$
a+b=(a \vee b) \wedge(a \wedge b)^{\prime}=a \wedge b^{\prime}=a \cdot b^{\prime}
$$

then also $(1+x \cdot y)+x \cdot(1+x \cdot y)=(1+x \cdot y) \cdot(1+x \cdot(1+x \cdot y))$. Now we apply $(* *)$ to compute

$$
\begin{aligned}
(1+x \cdot y)+(x+x \cdot y) & =(1+x \cdot y)+x \cdot(1+x \cdot y)=(1+x \cdot y) \cdot(1+x \cdot(1+x \cdot y)) \\
& =\left[(x \wedge y) \vee\left(x \wedge(x \wedge y)^{\prime}\right]^{\prime}=x^{\prime}=1+x\right.
\end{aligned}
$$

Thus $\mathcal{R}(L)$ satisfies (OM).
Conversely, let $\mathcal{R}$ be a $\lambda$-Boolean quasiring of characteristic 2 satisfying (OM). Let $x \leqslant y$ in $\mathcal{L}(R)$. Then

$$
\begin{aligned}
x \vee\left(y \wedge x^{\prime}\right) & =1+(1+x) \cdot(1+y \cdot(1+x)) \\
& =1+((1+x)+y \cdot(1+x))=1+((1+x \cdot y)+y \cdot(1+x \cdot y)) \\
& =1+((1+x \cdot y)+(y+x \cdot y))=1+(1+y)=y
\end{aligned}
$$

(by $(*)$ of the proof of Theorem 2). Hence, $\mathcal{L}(R)$ is an orthomodular $\lambda$-lattice.

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