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A *p*-Laplacian system with resonance and nonlinear boundary conditions on an unbounded domain

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Abstract. We study a nonlinear elliptic system with resonance part and nonlinear boundary conditions on an unbounded domain. Our approach is variational and is based on the well known Landesman-Laser type conditions.

Keywords: quasilinear problem, $p\mbox{-} Laplacian system, Landesman-Laser condition, resonance$

Classification: 35D05, 35J45, 35J50

1. Introduction and statement of results

Let Ω be an unbounded domain in \mathbb{R}^N , $N \geq 3$, with a noncompact and smooth boundary $\partial \Omega$. In this paper we consider the following quasilinear elliptic system

(1)
$$\begin{cases} -\Delta_p u = \lambda_1 a(x) |u|^{p-2} u + \lambda_1 b(x) |u|^{\alpha} |v|^{\beta} v + g_1(x, u) - h_1(x), \ x \in \Omega \\ -\Delta_p v = \lambda_1 d(x) |v|^{p-2} v + \lambda_1 b(x) |u|^{\alpha} |v|^{\beta} u + g_2(x, u) - h_2(x), \ x \in \Omega \end{cases}$$

subject to the nonlinear boundary conditions

(2)
$$\begin{cases} |\nabla u|^{p-2}\nabla u \cdot \eta + c_1(x)|u|^{p-2}u = 0, & x \in \partial\Omega\\ |\nabla v|^{p-2}\nabla v \cdot \eta + c_2(x)|v|^{p-2}v = 0, & x \in \partial\Omega \end{cases}$$

where $\Delta_p u = \operatorname{div}(|\nabla u|^{p-2}\nabla u)$ and η is the unit outward normal vector on $\partial\Omega$. On a single equation level with Ω bounded and Dirichlet boundary conditions, the problem has been studied by Arcoya and Orsina [1] taking into consideration the well known Landesman-Laser type conditions for the resonance part. The extension to the case of a system, again with Ω bounded and Dirichlet boundary conditions, was first considered by Zographopoulos in [7].

In order to confront with our problem we need a suitable space setting which we describe next.

For $\xi \in \mathbb{R}$, we set $w_{\xi}(x) := \frac{1}{(1+|x|)^{\xi}}$, and assume that the space $L^r(w_{\xi}, \Omega) := \{u : \int_{\Omega} w_{\xi}(x) | u|^r < +\infty\}, r \ge 1$, is supplied with the norm

$$||u||_{w_{\xi},r} = \left(\int_{\Omega} w_{\xi}(x)|u|^{r}\right)^{1/r}.$$

. .

Let $C^{\infty}_{\delta}(\Omega)$ be the space of $C^{\infty}_{0}(\mathbb{R}^{N})$ -functions restricted on Ω . For $p \in (1, +\infty)$, the weighted Sobolev space E_{p} is the completion of $C^{\infty}_{\delta}(\Omega)$ in the norm

$$|||u|||_p = \left(\int_{\Omega} |\nabla u|^p + \int_{\Omega} w_p(x) |u|^p\right)^{1/p}.$$

By Lemma 2 in [5], we see that if $c(\cdot)$ is a positive continuous function defined on \mathbb{R}^N such that

$$kw_{p-1}(x) \le c(x) \le Kw_{p-1}(x),$$

for some positive constants k and K, then the norm

$$\|u\|_{1,p} = \left(\int_{\Omega} |\nabla u|^p + \int_{\partial\Omega} c(x) |u|^m\right)^{1/p}$$

is equivalent to $\|\|\cdot\|\|_p$.

We will consider our system on the space $E = E_p \times E_p$, supplied with the norm

$$||(u, v)|| = ||u||_{1,p} + ||v||_{1,p}.$$

The following lemma is useful for our compactness arguments.

Lemma 1. (i) If

$$p \leq r \leq \frac{pN}{N-p}$$
 and $N > \alpha \geq N - r\frac{N-p}{p}$,

then the embedding $E \subseteq L^r(w_\alpha, \Omega)$ is continuous. If the upper bound for r in the first inequality and the upper bound for α in the second are strict, then the embedding is compact.

(ii) If

$$p \le m \le \frac{p(N-1)}{N-p}$$
 and $N > \beta \ge N - 1 - m \frac{N-p}{p}$,

then the trace operator $T: E \to L^m(w_\beta, \partial\Omega)$ is continuous. If the upper bound for *m* in the first inequality and the lower bound for β are strict, then the trace operator is compact.

(iii) If

$$1 \le q \frac{p}{q},$$

then the embedding $L^p(w_{\alpha_1}, \Omega) \subseteq L^q(w_{\alpha_2}, \Omega)$ is continuous.

PROOF: The first and second part of the lemma is Theorem 1 in [5], while the third is a consequence of the following inequality

$$\int_{\Omega} \frac{1}{(1+|x|)^{\alpha_2}} |u|^q \, dx \le \left(\int_{\Omega} \frac{1}{(1+|x|)^d} \, dx \right)^{\frac{p-q}{p}} \left(\int_{\Omega} \frac{1}{(1+|x|)^{\alpha_1}} |u|^p \, dx \right)^{\frac{q}{p}},$$

where $d = \frac{\alpha_2 p - \alpha_1 q}{p - q}$. Note that the integral $\int_{\Omega} \frac{1}{(1 + |x|)^d} dx$ converges since d > N.

We study (1)-(2) in connection with the eigenvalue problem

(3)
$$\begin{cases} -\Delta_p u = \lambda_1 a(x) |u|^{p-2} u + \lambda_1 b(x) |u|^{\alpha} |v|^{\beta} v, \\ -\Delta_p v = \lambda_1 d(x) |v|^{p-2} v + \lambda_1 b(x) |u|^{\alpha} |v|^{\beta} u, \end{cases}$$

subject to the boundary conditions (2), which was considered in [4] under the following set of assumptions, also needed for the present problem:

(H1)
$$2 , $\alpha, \beta \ge 0$ with $\alpha + \beta = p - 2$ and $\alpha + 1$, $\beta + 1 \le \frac{pp^*}{N}$, where $p^* = \frac{Np}{N-p}$.$$

(H2) (i) There exist positive constants α_1 , A with $\alpha_1 \in \left(p + \frac{(\beta+1)(N-p)}{p^*}, N\right)$ such that $0 < a(x) \le Aw_{\alpha_1}(x)$ a.e. in Ω .

(ii) There exist positive constants α_2 , D with $\alpha_2 \in \left(p + \frac{(\alpha+1)(N-p)}{p^*}, N\right)$ such that

$$0 < d(x) \le Dw_{\alpha_2}(x)$$
 a.e. in Ω .

(iii) $m\{x \in \Omega : b(x) > 0\} > 0$ and

$$0 \le b(x) \le Bw_s(x)$$
 a.e. in Ω ,

where B > 0 and $s \in (p, N)$.

(H3) $c_1(\cdot)$ and $c_2(\cdot)$ are positive continuous functions defined on \mathbb{R}^N with

$$kw_{p-1}(x) \le c_1(x), c_2(x) \le Kw_{p-1}(x),$$

for some positive constants k and K. Let

$$I(u,v) = \frac{\alpha+1}{p} \int_{\Omega} |\nabla u|^p + \frac{\alpha+1}{p} \int_{\partial \Omega} c_1(x) |u|^p + \frac{\beta+1}{p} \int_{\Omega} |\nabla v|^p + \frac{\beta+1}{p} \int_{\partial \Omega} c_2(x) |v|^p$$

and

$$J(u,v) = \frac{\alpha+1}{p} \int_{\Omega} a(x)|u|^p + \frac{\beta+1}{p} \int_{\Omega} d(x)|v|^p + \int_{\Omega} b(x)|u|^{\alpha}|v|^{\beta}uv.$$

Theorem 2 ([4]). Let Ω be an unbounded domain in \mathbb{R}^N , $N \geq 2$, with a noncompact and smooth boundary $\partial \Omega$. Assume that hypotheses (H1), (H2) and (H3) hold. Then

(a) the system (3) admits a positive principal eigenvalue λ_1 given by

$$\lambda_1 = \inf\{I(u, v) : J(u, v) = 1\}.$$

Each component of the associated normalized eigenfunction (u_1, v_1) is positive on $\overline{\Omega}$ and of class $C^{1,\delta}_{\text{loc}}(\Omega)$ for some $\delta \in (0,1)$.

(b) the set of eigenfunctions corresponding to λ_1 forms a one dimensional manifold $X \subseteq E$ defined by

$$X = \{c(u_1, v_1); \ c \in \mathbb{R} \setminus \{0\}\}.$$

(c) λ_1 is isolated, in the sense that there exists $\eta > 0$ such that the interval $(0, \lambda_1 + \eta)$ does not contain any other eigenvalue than λ_1 .

We make the following assumptions concerning the resonance part:

(H4) (i) $g_1(\cdot, \cdot), g_2(\cdot, \cdot)$ are Caratheodory functions such that $|g_1(x,s)| \leq \frac{C_1}{(1+|x|)^{\alpha_3}}$ and $|g_2(x,s)| \leq \frac{C_2}{(1+|x|)^{\alpha_4}}$, where $\alpha_3 > N - \frac{N-\alpha_1}{p}, \alpha_4 > N - \frac{N-\alpha_2}{p}, C_1, C_2$ are positive constants, and the limits

$$\lim_{s \to \pm \infty} g_i(x,s) = g_i^{\pm}(x), \qquad i = 1, 2,$$

 $s \to \pm \infty^{s \to (+, +)}$ exist for almost every $x \in \Omega$.

(ii)
$$|h_1(x)| \le \frac{H_1}{(1+|x|)^{\alpha_3}}$$
 and $|h_2(x)| \le \frac{H_2}{(1+|x|)^{\alpha_4}}$ for some positive constants H_1, H_2 .

Furthermore, we will need the following inequalities

(4)
$$L^+ < (\alpha + 1) \int_{\Omega} h_1(x) u_1 + (\beta + 1) \int_{\Omega} h_2(x) v_1 < L^-,$$

(5)
$$L^{-} < (\alpha + 1) \int_{\Omega} h_1(x) u_1 + (\beta + 1) \int_{\Omega} h_2(x) v_1 < L^+,$$

where (u_1, v_1) is the normalized eigenfunction of (3)–(2) with positive components and

$$L^{+} = (\alpha + 1) \int_{\Omega} g_{1}^{+}(x)u_{1} + (\beta + 1) \int_{\Omega} g_{2}^{+}(x)v_{1},$$

$$L^{-} = (\alpha + 1) \int_{\Omega} g_{1}^{-}(x)u_{1} + (\beta + 1) \int_{\Omega} g_{2}^{-}(x)v_{1}.$$

Inequalities (4) and (5) are the adaptation to the case of systems of the Landesman-Laser type conditions for scalar equations.

The energy functional of the problem (1)-(2) is

$$\begin{split} \Phi(u,v) &= \frac{\alpha+1}{p} \int_{\Omega} |\nabla u|^p + \frac{\alpha+1}{p} \int_{\partial\Omega} c_1(x)|u|^p - \lambda_1 \frac{\alpha+1}{p} \int_{\Omega} a(x)|u|^p \\ &- (\alpha+1) \int_{\Omega} G_1(x,u) + (\alpha+1) \int_{\Omega} h_1(x)u \\ &+ \frac{\beta+1}{p} \int_{\Omega} |\nabla v|^p + \frac{\beta+1}{p} \int_{\partial\Omega} c_2(x)|v|^p - \lambda_1 \frac{\beta+1}{p} \int_{\Omega} d(x)|v|^p \\ &- (\beta+1) \int_{\Omega} G_2(x,v) + (\beta+1) \int_{\Omega} h_2(x)v - \lambda_1 \int_{\Omega} b(x)|u|^\alpha |v|^\beta uv, \end{split}$$

where

$$G_i(x,s) = \int_0^s g_i(x,t) dt, \quad i = 1, 2.$$

In view of (H1)–(H3), the functional Φ is well defined and continuously differentiable on *E*. By a *weak solution* of (1)–(2) we mean an element of *E* which is a critical point of Φ .

The main result of this work is the following theorem:

Theorem 3. (i) Assume that hypotheses (H1)-(H3) and inequality (4) or (5) hold. Then the system (1)-(2) admits a weak solution.

2. The main result

In view of Theorem 2(a), it is clear that $\lambda_1 \leq \min\{\lambda_u, \lambda_v\}$, where λ_u, λ_v are the first eigenvalues of the problems $-\Delta_p u = \lambda a(x)|u|^{p-2}u$ and $-\Delta_p v = \lambda d(x)|v|^{p-2}v$, with the boundary conditions (2), respectively. The following lemma shows that this inequality is actually strict.

Lemma 4. $\lambda_1 < \min\{\lambda_u, \lambda_v\}.$

PROOF: Let $u_0 > 0$ be an eigenfunction corresponding to λ_u and $v_0 > 0$ an eigenfunction corresponding to λ_v . If $\lambda_u = \lambda_v$, then

$$\lambda_1 \le \frac{I(u_0, v_0)}{J(u_0, v_0)} < \lambda_u,$$

so without loss of generality we may assume that $\lambda_u < \lambda_v$. Let t > 0 be such that

(6)
$$\frac{\beta+1}{p} \int_{\Omega} d(x) |v_0|^p < \frac{\lambda_u}{\lambda_v - \lambda_u} \int_{\Omega} b(x) |tu_0|^{\alpha} |v_0|^{\beta} t u_0 v_0.$$

Then, in view of (6),

$$\lambda_1 \le \frac{I(tu_0, v_0)}{J(tu_0, v_0)} < \lambda_u = \min\{\lambda_u, \lambda_v\}.$$

Note that due to assumptions H(1)–H(4), the operators $A,\,N,\,B,\,C:E\to E^*$ given by

$$\begin{split} \langle A(u,v),(\varphi,\psi)\rangle &:= \int_{\Omega} |\nabla u|^{p-2} u \nabla \varphi + \int_{\Omega} |\nabla v|^{p-2} \nabla v \nabla \psi, \\ \langle N(u,v),(\varphi,\psi)\rangle &:= \int_{\Omega} a(x) |u|^{p-2} u \varphi - \int_{\partial\Omega} c_1(x) |u|^{p-2} u \varphi \\ &+ \int_{\Omega} d(x) |v|^{p-2} v \psi - \int_{\partial\Omega} c_2(x) |v|^{p-2} v \psi, \\ \langle B(u,v),(\varphi,\psi)\rangle &:= \int_{\Omega} b(x) |u|^{\alpha} |v|^{\beta} v \varphi + \int_{\Omega} b(x) |u|^{\alpha} |v|^{\beta} u \psi, \\ \langle C(u,v),(\varphi,\psi)\rangle &:= \int_{\Omega} (g_1(x,u) - h_1(x)) \varphi + \int_{\Omega} (g_2(x,v) - h_2(x)) \psi, \end{split}$$

are well defined. Following standard arguments based on the embeddings given in Lemma 1, we have:

Lemma 5. The operators A, N, B and C are continuous. Moreover, N, B and C are compact.

We can now proceed with the proof of the main result:

PROOF OF THEOREM 3: We assume first that (4) holds. We claim that Φ satisfies the PS-condition. Indeed, let $\{(u_n, v_n)\}_{n \in N}$ be a PS-sequence in E. Then

(7)
$$-c \le \Phi(u_n, v_n) \le c,$$

for some c > 0, and there exists a sequence $\{\varepsilon_n\}_{n \in \mathbb{N}}$ converging to 0^+ , such that

(8)
$$-\varepsilon_n \|(u,v)\| \le \Phi'(u_n,v_n)(u,v) \le \varepsilon_n \|(u,v)\|$$
 for every $(u,v) \in E$.

We will show that the sequence $\{(u_n, v_n)\}_{n \in N}$ is bounded in E. Assume the contrary, that is $||(u_n, v_n)|| \to +\infty$. Let

(9)
$$\widehat{u}_n := \frac{u_n}{\|(u_n, v_n)\|}, \quad \widehat{v}_n := \frac{v_n}{\|(u_n, v_n)\|}.$$

Since $\|\hat{u}_n\|_{E_p} \leq 1$ and $\|\hat{v}_n\|_{E_p} \leq 1$, by passing to subsequences if necessary, we may assume that $\hat{u}_n \to \hat{u}$ and $\hat{v}_n \to \hat{v}$ weakly in E_p . Due to our hypotheses on h_1 and g_1 we obtain

(10)
$$\lim_{n \to +\infty} \int_{\Omega} \frac{G_1(x, u_n)}{\|(u_n, v_n)\|^p} = \lim_{n \to +\infty} \int_{\Omega} \frac{h_1 u_n}{\|(u_n, v_n)\|^p} = 0$$

and similarly for $G_2(\cdot, \cdot)$ and $h_2(\cdot)$. Dividing (7) by $||(u_n, v_n)||^p$ and using (10), we arrive at

$$\begin{split} &\lim_{n \to +\infty} \sup_{n \to +\infty} \left[\frac{\alpha + 1}{p} \left\{ \int_{\Omega} |\nabla \widehat{u}_{n}|^{p} + \int_{\partial \Omega} c_{1}(x) |\widehat{u}_{n}|^{p} - \lambda_{1} \int_{\Omega} a(x) |\widehat{u}_{n}|^{p} \right\} \\ &+ \frac{\beta + 1}{p} \left\{ \int_{\Omega} |\nabla \widehat{v}_{n}|^{p} + \int_{\partial \Omega} c_{2}(x) |\widehat{v}_{n}|^{p} - \lambda_{1} \int_{\Omega} d(x) |\widehat{v}_{n}|^{p} \right\} \\ &- \lambda_{1} \int_{\Omega} b(x) |\widehat{u}_{n}|^{\alpha} |\widehat{v}_{n}|^{\beta} \widehat{u}_{n} \widehat{v}_{n} \right] \leq 0, \end{split}$$

and Lemma 1 gives

$$\begin{split} &\limsup_{n \to +\infty} \left[\frac{\alpha + 1}{p} \left\{ \int_{\Omega} |\nabla \widehat{u}_n|^p + \int_{\partial \Omega} c_1(x) |\widehat{u}_n|^p \right\} \\ &+ \frac{\beta + 1}{p} \left\{ \int_{\Omega} |\nabla \widehat{v}_n|^p + \int_{\partial \Omega} c_2(x) |\widehat{v}_n|^p \right\} \right] \\ &\leq \lambda_1 \left(\frac{\alpha + 1}{p} \int_{\Omega} a(x) |\widehat{u}|^p + \frac{\beta + 1}{p} \int_{\Omega} d(x) |\widehat{v}|^p + \int_{\Omega} b(x) |\widehat{u}|^{\alpha} |\widehat{v}|^{\beta} \widehat{u} \widehat{v} \right). \end{split}$$

The reverse inequality (with the limsup replaced by liminf) also holds due to the lower semicontinuity of the norms. Thus (\hat{u}, \hat{v}) is a nonzero solution of (3) with $\|(\hat{u}, \hat{v})\| = 1$. In view of Lemma 4, $\hat{u} \neq 0$ and $\hat{v} \neq 0$. By Theorem 2, \hat{u} and \hat{v} have the same sign. Suppose that both \hat{u} and \hat{v} are positive, the other case can be treated similarly. Thus $\hat{u} = u_1$ and $\hat{v} = v_1$. If we replace (u, v) by (u_n, v_n) in (8), write the relation for $-\Phi'$, multiply the members of (7) by p, add memberwise the resulting inequalities, and divide by $\|(u_n, v_n)\|$, we obtain

$$\begin{aligned} \left| (\alpha+1)(p-1) \int_{\Omega} h_1(x) \widehat{u}_n + (\beta+1)(p-1) \int_{\Omega} h_2(x) \widehat{v}_n \right. \\ \left. - (\alpha+1)p \int_{\Omega} \widehat{g}_1(x,u_n) \widehat{u}_n + (\alpha+1) \int_{\Omega} g_1(x,u_n) \widehat{u}_n - (\beta+1)p \int_{\Omega} \widehat{g}_2(x,v_n) \widehat{v}_n \right. \\ \left. + (\beta+1) \int_{\Omega} g_2(x,v_n) \widehat{v}_n \right| &\leq \frac{c}{\|(u_n,v_n)\|} + \varepsilon_n, \end{aligned}$$

where

$$\widehat{g}_i(x,s) := \begin{cases} \frac{G_i(x,s)}{s} & \text{if } s \neq 0, \\ g_i(x,0) & \text{if } s = 0, \end{cases} \quad i = 1, 2.$$

By letting $n \to +\infty$, we get

(11)

$$\lim_{n \to +\infty} \left\{ (\alpha + 1) \int_{\Omega} [g_1(x, u_n) \widehat{u}_n - p \widehat{g}_1(x, u_n) \widehat{u}_n] + (\beta + 1) \int_{\Omega} [g_2(x, v_n) \widehat{v}_n - p \widehat{g}_2(x, v_n) \widehat{v}_n] \right\}$$

$$= (\alpha + 1)(1 - p) \int_{\Omega} h_1(x) \widehat{u} + (\beta + 1)(1 - p) \int_{\Omega} h_2(x) \widehat{v}$$

By (9), $u_n(x)$ and $v_n(x)$ tend to $+\infty$, so

$$g_1(x, u_n) \to g_1^+(x)$$
 and $g_2(x, v_n) \to g_2^+(x)$ a.e. in Ω .

Therefore

(12)
$$\lim_{n \to +\infty} \int_{\Omega} \left[g_1(x, u_n) \widehat{u}_n - p \widehat{g}_1(x, u_n) \widehat{u}_n \right] = (1-p) \int_{\Omega} g_1^+(x) \widehat{u},$$

with a similar relation holding for $g_2(\cdot, \cdot)$ as well. In view of (11) and (12), we have

$$(\alpha+1)\int_{\Omega}g_{1}^{+}(x)u_{1} + (\beta+1)\int_{\Omega}g_{2}^{+}(x)v_{1} = (\alpha+1)\int_{\Omega}h_{1}(x)u_{1} + (\beta+1)\int_{\Omega}h_{2}(x)v_{1},$$

contradicting (4). Thus $\{(u_n, v_n)\}_{n \in N}$ is bounded. Therefore, up to subsequences, $u_n \to u_0$ and $v_n \to v_0$ weakly in E_p and strongly in $L^p(w_{\alpha_1}, \Omega)$ and $L^p(w_{\alpha_2}, \Omega)$, respectively. By taking $(u, v) = (u_n, v_n) - (u_0, v_0)$ in (8), and using Lemma 1, we derive that

$$\begin{aligned} &(\alpha+1)\left\{\int_{\Omega}\left(|\nabla u_{n}|^{p-2}\nabla u_{n}-|\nabla u_{0}|^{p-2}\nabla u_{0}\right)\left(\nabla u_{n}-\nabla u_{0}\right)\right.\\ &+\int_{\partial\Omega}c_{1}\left(|u_{n}|^{p-2}u_{n}-|u_{0}|^{p-2}u_{0}\right)\left(u_{n}-u_{0}\right)\right\}\\ &+\left(\beta+1\right)\left\{\int_{\Omega}\left(|\nabla v_{n}|^{p-2}\nabla v_{n}-|\nabla v_{0}|^{p-2}\nabla v_{0}\right)\left(\nabla v_{n}-\nabla v_{0}\right)\right.\\ &+\int_{\partial\Omega}c_{1}\left(|v_{n}|^{p-2}v_{n}-|v_{0}|^{p-2}v_{0}\right)\left(v_{n}-v_{0}\right)\right\}\rightarrow0,\end{aligned}$$

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which, in view of inequality 2.5 in [2], implies that $(u_n, v_n) \to (u_0, v_0)$ in E. We show next that Φ is coercive. Indeed, if this were not the case, there would exist a sequence $\{(u_n, v_n)\}_{n \in \mathbb{N}}$ with $||(u_n, v_n)|| \to +\infty$ and

(13)
$$|\Phi(u_n, v_n)| \le M, \text{ for some } M > 0.$$

Working as before, we get that

$$\begin{split} \lim_{n \to +\infty} & \left[\frac{\alpha + 1}{p} \left\{ \int_{\Omega} |\nabla \widehat{u}_n|^p + \int_{\partial \Omega} c_1(x) |\widehat{u}_n|^p - \lambda_1 \int_{\Omega} a(x) |\widehat{u}_n|^p \right\} \\ & + \frac{\beta + 1}{p} \left\{ \int_{\Omega} |\nabla \widehat{v}_n|^p + \int_{\partial \Omega} c_2(x) |\widehat{v}_n|^p - \lambda_1 \int_{\Omega} d(x) |\widehat{v}_n|^p \right\} \\ & - \lambda_1 \int_{\Omega} b(x) |\widehat{u}_n|^\alpha |\widehat{v}_n|^\beta \widehat{u}_n \widehat{v}_n \right] = 0, \end{split}$$

where \hat{u}_n and \hat{v}_n are defined in (9). Thus $(\hat{u}_n, \hat{v}_n) \to (u_1, v_1)$ or $(\hat{u}_n, \hat{v}_n) \to -(u_1, v_1)$ in *E*. If $(\hat{u}_n, \hat{v}_n) \to (u_1, v_1)$, by (13) and the variational characterization of λ_1 , we obtain

$$(\alpha+1)\int_{\Omega}g_{1}^{+}(x)u_{1} + (\beta+1)\int_{\Omega}g_{2}^{+}(x)v_{1} \ge (\alpha+1)\int_{\Omega}h_{1}(x)u_{1} + (\beta+1)\int_{\Omega}h_{2}(x)v_{1},$$

while if $(\widehat{u}_n, \widehat{v}_n) \to -(u_1, v_1)$, we get

$$(\alpha+1)\int_{\Omega}g_{1}^{-}(x)u_{1}+(\beta+1)\int_{\Omega}g_{2}^{-}(x)v_{1} \leq (\alpha+1)\int_{\Omega}h_{1}(x)u_{1}+(\beta+1)\int_{\Omega}h_{2}(x)v_{1},$$

contradicting (4). We can now use Theorem 4.7 in [3] to get a weak solution of (1)-(2).

Assume next that (5) holds. We split E as the direct sum of the eigenspace X and $Y = \{(u, v) \in E : \int_{\Omega} u u_1^{p-1} + \int_{\Omega} v v_1^{p-1} = 0\}$. Then Φ has a saddle point geometry, i.e.,

- (i) $\Phi(t(u_1, v_1)) \to -\infty$ if $|t| \to +\infty$, and
- (ii) Φ is bounded from below on Y.

Indeed, since

$$\begin{split} \Phi(t(u_1, v_1)) &= (\alpha + 1) \left[\int_{\Omega} h_1(x) t u_1 - \int_{\Omega} G_1(x, t u_1) \right] \\ &+ (\beta + 1) \left[\int_{\Omega} h_2(x) t v_1 - \int_{\Omega} G_2(x, t v_1) \right] \\ &= (\alpha + 1) t \left[\int_{\Omega} h_1(x) u_1 - \int_{\Omega} \frac{G_1(x, t u_1)}{t u_1} u_1 \right] \\ &+ (\beta + 1) t \left[\int_{\Omega} h_2(x) v_1 - \int_{\Omega} \frac{G_2(x, t v_1)}{t v_1} v_1 \right], \end{split}$$

by taking the limit as $|t| \to \infty$ and working as in the first part of the proof, we can use (5) to get (i). To prove (ii) we exploit the isolation of λ_1 , see Theorem 2, to derive that there exists $\hat{\lambda} > \lambda_1$ such that

$$\widehat{\lambda} < \frac{I(u,v)}{J(u,v)}$$

for every $(u, v) \in Y$. If $(u, v) \in Y$, in view of Lemma 1,

$$\begin{split} \Phi(u,v) &= I(u,v) - \lambda_1 J(u,v) + (\alpha+1) \left[\int_{\Omega} h_1(x) u - \int_{\Omega} G_1(x,u) \right] \\ &+ (\beta+1) \left[\int_{\Omega} h_2(x) v - \int_{\Omega} G_2(x,v) \right] \\ &> \left(1 - \frac{\lambda_1}{\widehat{\lambda}} \right) I(u,v) - (\alpha+1) c_1 \|u\|_{1,p} - (\beta+1) c_2 \|v\|_{1,p}, \end{split}$$

for some $c_1, c_2 > 0$. Consequently, Φ is bounded from below on Y. An application of the saddle point theorem, see [6], provides a weak solution of (1)–(2).

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