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## On Certain Groups of Holomorphic Maps

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O. Consider the space $\mathbf{C}^{2}$ with the complex coordinates $(x, y)$. By $\Gamma_{s}$ denote the pseudogroup of local holomorphic diffeomorphisms of $\mathbf{C}^{2} \tilde{x}=\tilde{x}(x, y), \tilde{y}=\tilde{y}(x, y)$ satisfying $\partial(\tilde{x}, \tilde{y})!\partial(x, y)=1$. We are going to prove the following

Theorem. Let $G \subset \Gamma_{s}$ be a Lie group such that $\operatorname{dim} G=3$ and the orbits of $G$ are real hypersurfaces $M^{3} \subset \mathbf{R}^{4} \equiv \mathbf{C}^{2}$ with non-trivial Levi form. Then $G$ is locally $\Gamma_{\delta}$-equivalent to one of the following groups:
(I) $\quad \tilde{x}=x-\frac{a}{\alpha} y-\frac{1}{2} i B a^{2}+c, \quad \tilde{y}=y+\alpha b+i \alpha B a ; \quad a, b, c \in \mathbf{R}$;
(II) $\tilde{x}=a x-\frac{1}{\alpha} b y+c, \tilde{y}=-\alpha B b x+a y+\alpha d ; a, b, c, d \in \mathbf{R} ; a^{2}-B b^{2}=1$;
(III) $\tilde{x}=\frac{(a x+b)(1-b c-a c x) y^{2}-\alpha^{2} a^{2} c}{(1-b c-a c x)^{2} y^{2}+\alpha^{2} a^{2} c^{2}}$,

$$
\tilde{y}=\frac{(1-b c-a c x)^{2} y^{2}+\alpha^{2} a^{2} c^{2}}{a y} ; a, b, c \in \mathbf{R} ;
$$

(IV) consider (III) with $a \in i \mathbf{R}, b \in \mathbf{C}, c=\bar{b}$.

Here, $0 \neq \alpha \in \mathbf{C}$ and $\mathbf{0} \neq B \in \mathbf{R}$ are parameters. The corresponding orbits $M^{3}$ are
(I') $\left(\frac{y}{\alpha}-\frac{\bar{y}}{\bar{\alpha}}\right)^{2}+4 i B(x-\bar{x})=r$,
(II') $\left(\frac{y}{\alpha}-\frac{\bar{y}}{\bar{\alpha}}\right)^{2}-B(x-\bar{x})^{2}=r$,
(III') $\quad(x-\bar{x})^{2} y^{2} \bar{y}^{2}+(\alpha \bar{y}+\bar{\alpha} y)^{2}+4 r y \bar{y}=0$
with $\quad r \in \mathbf{R}$.
The groups (III) and (IV) will be studied elsewhere.
In the second part of this paper, I solve the equivalence problem for hypersurfaces of $\mathbf{C}^{2}$ with respect to the pseudogroup of all local biholomorphic mappings.

It is well known that two real hypersurfaces in $\mathbf{C}^{2}$ are not generally holomorphically equivalent. The problem of the construction of invariants of $M^{3} \subset \mathbf{C}^{2}$ with respect to the pseudogroup of holomorphic mappings has been treated by E. Cartan (Annali di Mat., t. 11. 1932, 17-90); unfortunately, his treatment is very confused.

Let $V^{3}$ be a differentiable manifold together with a structure consisting of a choice of two tangent directions at each of its points. In what follows, I shall construct (in the general case) an $\{e\}$-structure on $V^{3}$ invariantly associated to the given structure; by means of this $\{e\}$-structure, the equivalence problem of the structures of the just described type will be solved. Further, I will show that the construction of an invariant $\{e\}$-structure on $M^{3} \subset \mathbf{C}^{2}$ is equivalent to the preceding construction. The special cases will be treated in a forthcoming paper.

Parts of this paper have been written during my stays at the universities at Berlin (GDR) and Riga (USSR).

1. Consider the space $\mathbf{C}^{2}, \mathbf{C}$ being the complex numbers, with the complex coordinates $x=x^{1}+i y^{1}, y=x^{2}+i y^{2}$. Its real form is the space $\mathbf{R}^{4}$ ( $\mathbf{R}$ being reals) with the coordinates ( $x^{1}, y^{1}, x^{2}, y^{2}$ ) together with the endomorphism $I$ : $\mathbf{R}^{4} \rightarrow \mathbf{R}^{4}, I^{2}=-i d$, defined by ( $i=1,2$ )

$$
\begin{equation*}
I \frac{\partial}{\partial x^{i}}=\frac{\partial}{\partial y^{i}}, \quad I \frac{\partial}{\partial y^{i}}=-\frac{\partial}{\partial x^{i}} . \tag{1.1}
\end{equation*}
$$

In general, on any complex vector space $V$, scalar multiplication by real numbers is, of course, defined. Relative to addition, and scalar multiplication by real numbers only, the elements of $V$ clearly form a real vector space, which will be denoted by $V_{0}$ and called the real vector space underlying the complex vector space $V$. If $V_{0}$ is the underlying real vector space of a complex space $V$, then there is an automorphism $I_{0}$ of $V_{0}$ satisfying $I_{0}^{2}=-i d$, induced by the automorphism $I$ of $V$ given by $I A=i A, A \in V$. Further, $\operatorname{dim}_{R} V_{0}=2 \operatorname{dim}_{c} V$. Let $V$ be a finite dimensional complex vector space and $A_{1}, \ldots ., A_{n}$ its basis, then $A_{1}, I_{0} A_{1}, \ldots ., A_{n}, I_{0} A_{n}$ give a basis for $V_{0}$. Let $W_{0}$ be a real vector space (of finite dimension). We say that a complex structure is given on $W_{0}$ if there is given an endomorphism $I_{0}$ of $W_{0}$ satisfying $I_{0}^{2}=-i d$; this endomorphism is an automorphism, since $I_{0}{ }^{-1}$ exists and is given by $-I_{0}$. Let $W_{0}$ be a real vector space with a complex structure defined by $I_{0}$. Then: (i) There exists a basis for $W_{0}$ of the form $A_{1}, I_{0} A_{1}, \ldots, A_{n}, I_{0} A_{n}$; in particular, $\operatorname{dim}_{R} W_{0}$ is even; (ii) there exists a complex space $W$ such that $W_{0}$ is the underlying real vector space of $W$ and $I_{0}$ is induced by the complex structure of $W$. Let us prove this last proposition. Since $\operatorname{dim}_{R} W_{0}>0$, there exists a vector $A_{1} \neq 0$ in $W_{0}$. Then $A_{1}$ and $I_{0} A_{1}$ are independent. In fact, if there exist real numbers $a, b$ such that $a A_{1}+b I_{0} A_{1}=0$, then $a I A_{1}-b A_{1}=0$ and $\left(a^{2}+b^{2}\right) A_{1}=0$. This implies $a=b=0$. We proceed by induction, and assume that an independent set $A_{1}$, $I_{0} A_{1}, \ldots, A_{k}, I_{0} A_{k}$ of vectors in $W_{0}$ has been found $(k \geq 1)$. If $\operatorname{dim}_{R} W_{0}=2 k$, there is nothing further to prove. If $\operatorname{dim}_{R} W_{0}>2 k$, then there is a non-zero vector $A_{k+1} \in W_{0}$ which is independent of the vectors $A_{1}, \ldots, I_{0} A_{k}$. The vectors $A_{1}$,
$I_{0} A_{1}, \ldots, A_{k+1}, I_{0} A_{k+1}$ form an independent set. In fact, if $a_{1}, \ldots, a_{k+1}, b_{1}, \ldots, b_{k+1}$ are real numbers such that

$$
\begin{equation*}
\sum_{=1}^{k+1} a_{j} A_{j}+\sum_{j=1}^{k+1} b_{j} I_{0} A_{j}=0 \tag{1.2}
\end{equation*}
$$

then

$$
\sum_{j=1}^{k+1} a_{j} I_{0} A_{j}-\sum_{j=1}^{k+1} b_{j} A_{j}=0
$$

From these, we obtain
$\sum_{j=1}^{k}\left(a_{j} a_{k+1}+b_{j} b_{k+1}\right) A_{j}+\sum_{j=1}^{k+1}\left(b_{j} a_{k+1}-a_{j} b_{k+1}\right) I_{0} A_{j}+\left(a_{k+1}^{2}+b_{k+1}^{2}\right) A_{k+1}=0$.
All coefficients being zero, we have $a_{k+1}=b_{k+1}=0$, and (1.2) implies $a_{1}=$ $=\ldots .=a_{k}=b_{1}=\ldots=b_{k}=0$. The complex vector space $W$ is constructed from the elements of $W_{0}$ by defining the operation of scalar multiplication by a complex number $c=a+i b$ as $c A=a A+b I_{0} A$.

Let $\Gamma$ be the pseudogroup of all local holomorphic diffeomorphisms of $\mathbf{C}^{2}$. Each $\gamma \in \Gamma$ induces a diffeomorphism of $\mathbf{R}^{4}$ denoted by $\gamma$, too. The local diffeomorphism $\gamma$ of $\mathbf{R}^{4}$ given by

$$
\begin{equation*}
\tilde{x}^{i}=f^{i}\left(x^{j}, y^{j}\right), \quad \tilde{y}^{i}=g^{i}\left(x^{j}, y^{j}\right) ; \quad i=1,2 ; \tag{1.3}
\end{equation*}
$$

is an element of $\Gamma$ if and only if the functions $f^{i}, g^{i}$ satisfy the Cauchy-Riemann equations

$$
\begin{equation*}
\frac{\partial f^{i}}{\partial x^{j}}=\frac{\partial g^{i}}{\partial y^{j}}, \quad \frac{\partial f^{i}}{\partial y^{j}}=-\frac{\partial g^{i}}{\partial x^{j}} ; \quad i, j=1,2 . \tag{1.4}
\end{equation*}
$$

Let $\Gamma_{s} \subset \Gamma$ be the pseudogroup of diffeomorphisms $\tilde{x}=\tilde{x}(x, y), \tilde{y}=\tilde{y}(x, y)$ of the space $\mathbf{C}^{2}$ or $\mathbf{R}^{4}$ resp. satisfying

$$
\frac{\partial(\tilde{x}, \tilde{y})}{\partial(x, y)} \equiv\left|\begin{array}{ll}
\frac{\partial \tilde{x}}{\partial x} & \frac{\partial \tilde{x}}{\partial y}  \tag{1.5}\\
\frac{\partial \tilde{y}}{\partial x} & \frac{\partial \tilde{y}}{\partial y}
\end{array}\right|=1 .
$$

It is easy to see that $\gamma \in \Gamma$ is an element of $\Gamma_{s}$ if and only if $\gamma$ preserves the 2-form

$$
\begin{equation*}
\Phi=d x \wedge d y \tag{1.6}
\end{equation*}
$$

indeed,

$$
\widetilde{\Phi}=d \tilde{x} \wedge d \tilde{y}=\frac{\partial(\tilde{x}, \tilde{y})}{\partial(x, y)} \Phi
$$

Define

$$
\begin{equation*}
\varphi=d x^{1} \wedge d x^{2}-d y^{1} \wedge d y^{2}, \quad \psi=d x^{1} \wedge d y^{2}+d y^{1} \wedge d x^{2} \tag{1.7}
\end{equation*}
$$

obviously, $\Phi=\varphi+i \psi$. Of course, we may write

$$
\begin{equation*}
\varphi=\frac{1}{2}(d x \wedge d y+d \bar{x} \wedge d \bar{y}), \quad \psi=-\frac{1}{2} i(d x \wedge d y-d \bar{x} \wedge d \bar{y}) \tag{1.8}
\end{equation*}
$$

We have

$$
\begin{equation*}
\varphi(v, w)=-\varphi(I v, I w), \quad \psi(v, w)=-\varphi(v, I w) \text { for } v, w \in \mathbf{R}^{4} . \tag{1.9}
\end{equation*}
$$

Indeed, let

$$
\begin{align*}
& v=a^{1} \frac{\partial}{\partial x^{1}}+b^{1} \frac{\partial}{\partial y^{1}}+a^{2} \frac{\partial}{\partial x^{2}}+b^{2} \frac{\partial}{\partial y^{2}}, \\
& w=c^{1} \frac{\partial}{\partial x^{1}}+d^{1} \frac{\partial}{\partial y^{1}}+c^{2} \frac{\partial}{\partial x^{2}}+d^{2} \frac{\partial}{\partial y^{2}} . \tag{1.10}
\end{align*}
$$

Then

$$
\begin{aligned}
& I v=-b^{1} \frac{\partial}{\partial x^{1}}+a^{1} \frac{\partial}{\partial y^{1}}-b^{2} \frac{\partial}{\partial x^{2}}+a^{2} \frac{\partial}{\partial y^{2}}, \\
& I w=-d^{1} \frac{\partial}{\partial x^{1}}+c^{1} \frac{\partial}{\partial y^{1}}-d^{2} \frac{\partial}{\partial x^{2}}+c^{2} \frac{\partial}{\partial y^{2}}
\end{aligned}
$$

and

$$
\begin{align*}
& \varphi(v, w)=a^{1} c^{2}-a^{2} c^{1}-b^{1} d^{2}+b^{2} d^{1}=-\varphi(I v, I w), \\
& \psi(v, w)=a^{1} d^{2}-b^{2} c^{1}+b^{1} c^{2}-a^{2} d^{1}=-\varphi(v, I w) . \tag{1.11}
\end{align*}
$$

In $\mathbf{C}^{2}$, this may be rewritten as follows. Introduce the well known vector fields

$$
\frac{\partial}{\partial x}=\frac{1}{2}\left(\frac{\partial}{\partial x^{1}}-i \frac{\partial}{\partial y^{1}}\right), \quad \frac{\partial}{\partial \bar{x}}=\frac{1}{2}\left(\frac{\partial}{\partial x^{1}}+i \frac{\partial}{\partial y^{1}}\right), \ldots .
$$

Then

$$
\frac{\partial}{\partial x^{1}}=\frac{\partial}{\partial x}+\frac{\partial}{\partial \bar{x}}, \quad \frac{\partial}{\partial y^{1}}=i\left(\frac{\partial}{\partial x}-\frac{\partial}{\partial \bar{x}}\right), \ldots .
$$

and the vectors $v, w$ may be written as

$$
\begin{aligned}
& v=A^{1} \frac{\partial}{\partial x}+A^{2} \frac{\partial}{\partial y}+\bar{A}^{1} \frac{\partial}{\partial \bar{x}}+\bar{A}^{2} \frac{\partial}{\partial \bar{y}}, \\
& w=C^{1} \frac{\partial}{\partial x}+C^{2} \frac{\partial}{\partial y}+\bar{C}^{1} \frac{\partial}{\partial \bar{x}}+\bar{C}^{2} \frac{\partial}{\partial \bar{y}}
\end{aligned}
$$

with

$$
A^{i}=a^{i}+i b^{i}, \quad C^{i}=c^{i}+i d^{i} ; \quad i=1,2
$$

It is easy to check that

$$
\begin{aligned}
& I v=i A^{1} \frac{\partial}{\partial x}+i A^{2} \frac{\partial}{\partial y}-i \overline{A^{1}} \frac{\partial}{\partial \bar{x}}-i \overline{A^{2}} \frac{\partial}{\partial \bar{y}}, \\
& I w=i C^{1} \frac{\partial}{\partial x}+i C^{2} \frac{\partial}{\partial y}-i \overline{C^{1}} \frac{\partial}{\partial \bar{x}}-i \overline{C^{2}} \frac{\partial}{\partial \bar{y}}
\end{aligned}
$$

and

$$
\begin{aligned}
& \varphi(v, w)=\frac{1}{2}\left(A^{1} C^{2}-A^{2} C^{1}+\bar{A}^{1} \bar{C}^{2}-\bar{A}^{2} \bar{C}^{1}\right)=-\varphi(I v, I w) \\
& \psi(v, w)=-\frac{1}{2} i\left(A^{1} C^{2}-A^{2} C^{1}-\bar{A}^{1} \bar{C}^{2}+\bar{A}^{2} \bar{C}^{1}\right)=-\varphi(v, I w) .
\end{aligned}
$$

Let $X=X(x, y), Y=Y(x, y)$ be a local holomorphic diffeomorphism of $\mathbf{C}^{2}$.
Then

$$
d X \wedge d Y+d \bar{X} \wedge d \bar{Y}=\frac{\partial(X, Y)}{\partial(x, y)} d x \wedge d y+\frac{\overline{\partial(X, Y)}}{\partial(x, y)} d \bar{x} \wedge d \bar{y}
$$

Thus: Let $\gamma$ be a local diffeomorphism of $\mathbf{R}^{4}$ defined on $U \subset \mathbf{R}^{4}$. Then $\gamma \in \Gamma_{s}$ if and only if
for each

$$
\begin{gather*}
\left(d \gamma_{a} \cdot I\right)\left(v_{a}\right)=\left(I . d \gamma_{a}\right)\left(v_{a}\right),  \tag{1.12}\\
\varphi\left(v_{a}, w_{a}\right)=\varphi\left(d \gamma_{a}\left(v_{a}\right), d \gamma_{a}\left(w_{a}\right)\right) \\
a \in U ; \quad v_{a}, w_{a} \in T_{a}\left(\mathbf{R}^{4}\right) \equiv \mathbf{R}^{4} .
\end{gather*}
$$

From now on, consider the following situation: In $\mathbf{R}^{4}$ with the coordinates ( $x^{1}, y^{1}, x^{2}, y^{2}$ ) be given a complex structure $I(1.1)$ and the form (1.71); let $\Gamma_{s}$ be the pseudogroup of local diffeomorphisms of $\mathbf{R}^{4}$ satisfying (1.12).

Now, let $M^{3} \subset \mathbf{R}^{4}$ be a hypersurface. At each point $m \in M^{3}$, consider the space

$$
\begin{equation*}
\tau_{m}=T_{m}\left(M^{3}\right) \bigcap I T_{m}\left(M^{3}\right) \tag{1.13}
\end{equation*}
$$

Obviously, $\operatorname{dim} \tau_{m}=2$ and $I\left(\tau_{m}\right)=\tau_{m}$. The pseudogroup $\Gamma_{s}$ induces on $M^{3}$ the following structure: at each point $m \in M^{3}$, we have a tangent plane $\tau_{m}$ and its endomorphism $I_{m}: \tau_{m} \rightarrow \tau_{m}$ satisfying $I_{m}^{2}=-i d$; further, there is given a 2 -form $\varphi^{\star}$ (the restriction of $\varphi$ ) on $M^{3}$ such that

$$
\varphi^{\star}\left(v_{m}, w_{m}\right)=-\varphi^{\star}\left(I_{m} v_{m}, I_{m} w_{m}\right) \quad \text { for } v_{m}, w_{m} \in \tau_{m} .
$$

Of course, $\varphi^{\star}$ 三 $\mathbf{0}$.
2. Let us suppose that the field of planes $\tau_{m}$ is non-integrable. Let us investigate this supposition more carefully. Define a partial complex structure on a manifold $X$, $\operatorname{dim} X=p$, as an assignment of a tangent space $\tau_{x} \subset T_{x}(X)$ and an endomorphism $I_{x}: \tau_{x} \rightarrow \tau_{x}, I_{x}^{2}=-i d$, to each point $x \in X$; let $\operatorname{dim} \tau_{x}=2 q$. Consider a fixed point $x_{0} \in X$ and its neighbourhood $U$ such that there are tangent vector fields $v_{1}, \ldots, v_{q}, w_{1}, \ldots, w_{q}, u_{1}, \ldots, u_{p-2 q}$ in $U$ satisfying $v_{i}(x), w_{i}(x) \in \tau_{x}$ and $I_{x} v_{i}(x)=$ $=w_{i}(x)$ in $U$; write $i, j, \ldots=1, \ldots ., q ; \alpha, \beta, \ldots=1, \ldots, p-2 q$. Then

$$
\begin{align*}
& {\left[v_{i}, v_{j}\right]=a_{i j}^{k} v_{k}+b_{i j}^{k} w_{k}+c_{i}^{a} u_{a},}  \tag{2.1}\\
& {\left[v_{i}, w_{j}\right]=d_{i j}^{k} v_{k}+e_{i j}^{k} w_{k}+f_{i j}^{a} u_{a},} \\
& {\left[w_{i}, w_{j}\right]=g_{i j}^{k} v_{k}+h_{i j}^{k} w_{k}+k_{i j}^{a} u_{a} .}
\end{align*}
$$

Let $V_{0} \in \tau_{x}$, be a fixed vector. On $U$, consider an arbitrary vector field $V$ such that $V\left(x_{0}\right)=V_{0}$ and $V(x) \in \tau_{x}$ for each $x \in U$. Then there are functions $p^{i}$, $q^{4}$ (on $U$ ) such that

$$
\begin{equation*}
V=p^{i} v_{i}-q^{i} w_{i} . \tag{2.2}
\end{equation*}
$$

At each point $x \in U$, consider the vector $I V$; of course,

$$
\begin{equation*}
I V=q^{i} v_{i}+p^{i} w_{i} . \tag{2.3}
\end{equation*}
$$

We have

$$
\begin{align*}
{[V, I V]=} & {\left[p^{i} v_{i}-q^{i} w_{i}, q^{j} v_{j}+p^{j} w_{j}\right]=}  \tag{2.4}\\
= & \left(p^{i} \cdot v_{i} q^{k}-q^{i} \cdot w_{i} q^{k}-q^{i} \cdot v_{i} p^{k}-p^{i} \cdot w_{i} p^{k}+a_{i j}^{k} p^{i} q^{j}+\right. \\
+ & \left.d_{i j}^{k} p^{j} p^{j}+d_{i j}^{k} q^{i} q^{j}-g_{i j}^{k} p^{j} q^{i}\right) v_{k}+\left(p^{i} \cdot v_{i} p^{k}-q^{i} \cdot w_{i} p^{k}+\right. \\
& +q^{i} \cdot v_{i} q^{k}+p^{i} \cdot w_{i} q^{k}+b_{i j}^{k} p^{i} q^{j}+e_{i j}^{k} p^{i} p^{j}+e_{i j}^{k} q^{i} q^{j}- \\
& \left.-h_{i j}^{k} q^{i} p^{j}\right) w_{k}+\left(c_{i j}^{a} p^{i} q^{j}+f_{i j}^{a} p^{i} p^{j}+f_{i j}^{a} q^{i} q^{j}-k_{i j}^{a} p^{j} q^{i}\right) u_{a} .
\end{align*}
$$

Let $\pi_{x}: T_{x}(X) \rightarrow T_{x}(X) / \tau_{x}$ be the natural projection. We see from (2.4) that

$$
\begin{equation*}
L_{x_{0}}\left(V_{0}\right)=\pi_{x_{0}}\left([V, I V]\left(x_{0}\right)\right) \in T_{x_{0}}(X) / \tau_{x_{0}} \tag{2.5}
\end{equation*}
$$

does not depend on the choice of the field $V$ extending the vector $V_{0}$. Thus we get a well defined map

$$
\begin{equation*}
L_{x}: \tau_{x} \rightarrow T_{x}(X) / \tau_{x} \tag{2.6}
\end{equation*}
$$

which is called the Levi map of the given partial complex structure (at the point $x \in X$ ). If $v_{i}, w_{i}, u_{a} \in T_{x}(X)$ as above and $\tilde{u}_{a}=\pi\left(u_{a}\right) \in T_{x}(X) / \tau_{x}$, then

$$
\begin{gather*}
L_{x}(V) \equiv L_{x}\left(p^{i} v_{i}-q^{i} w_{i}\right)=  \tag{2.7}\\
=\left(c_{i j}^{a} p^{i} q^{j}+f_{i j}^{a} p^{i} p^{j}+f_{i i}^{a} q^{i} q^{j}-k_{i j}^{a} p^{j} q^{i}\right) \tilde{u}_{a} .
\end{gather*}
$$

From this and (2.1), we see that the field $\left\{\tau_{x}\right\}$ is integrable if and only if $L_{x}(V)=0$ for each $x \in X$ and each $V \in \tau_{x}$.

To compare our notion of the Levi map with the well established notion of the Levi map used in the literature, let us calculate the Levi map of a real hypersurface $X^{2 n-1} \subset \mathbf{C}^{n}$. Suppose that $X^{2 n-1}$ is given by the equation

$$
\begin{equation*}
F\left(z^{1}, \ldots, z^{n}, \bar{z}^{1}, \ldots, \bar{z}^{n}\right)=0 \tag{2.8}
\end{equation*}
$$

in the neighbourhood of the point $z^{1}=0, \ldots, z^{n}=0$. Of course,

$$
\begin{equation*}
F\left(z^{1}, \ldots ., z^{n}, \bar{z}^{1}, \ldots, \bar{z}^{n}\right)=\overline{F\left(z^{1}, \ldots ., z^{n}, \bar{z}^{1}, \ldots, \bar{z}^{n}\right)} \tag{2.9}
\end{equation*}
$$

$F\left(z^{i}, \bar{z}^{i}\right)$ being a real function. In a suitable small neighbourhood of the origin of $\mathbf{C}^{n}$, consider the one-paramateric set of hypersurfaces

$$
\begin{equation*}
F\left(z^{1}, \ldots ., z^{n}, \bar{z}^{1}, \ldots, \bar{z}^{n}\right)=\alpha, \quad \alpha \in(-\varepsilon, \varepsilon) \tag{2.10}
\end{equation*}
$$

Let $v$ be a real vector field around the origin of $\mathbf{C}^{n}$. Then

$$
\begin{equation*}
v=A^{i} \frac{\partial}{\partial z^{i}}+\bar{A}^{i} \frac{\partial}{\partial \bar{z}^{i}}, \tag{2.11}
\end{equation*}
$$

and the vector field $I v$ is given by

$$
\begin{equation*}
I v=i A^{i} \frac{\partial}{\partial z^{i}}-i \overline{A^{i}} \frac{\partial}{\partial \bar{z}^{i}} . \tag{2.12}
\end{equation*}
$$

Indeed, write $z^{i}=x^{i}+i y^{i}$, and (as usual)

$$
\frac{\partial}{\partial z^{i}}=\frac{1}{2}\left(\frac{\partial}{\partial x^{i}}-i \frac{\partial}{\partial y^{i}}\right), \frac{\partial}{\partial \bar{z}^{i}}=\frac{1}{2}\left(\frac{\partial}{\partial x^{i}}+i \frac{\partial}{\partial y^{i}}\right) .
$$

Then

$$
\frac{\partial}{\partial x^{i}}=\frac{\partial}{\partial z^{i}}+\frac{\partial}{\partial \bar{z}^{i}}, \frac{\partial}{\partial y^{i}}=i\left(\frac{\partial}{\partial z^{i}}-\frac{\partial}{\partial \bar{z}^{i}}\right)
$$

and

$$
I \frac{\partial}{\partial x^{i}}=-\frac{\partial}{\partial y^{i}}, I \frac{\partial}{\partial y^{i}}=-\frac{\partial}{\partial x^{i}} .
$$

Then

$$
\begin{aligned}
& v=a^{i} \frac{\partial}{\partial x^{i}}+b^{i} \frac{\partial}{\partial y^{i}}=\left(a^{i}+i b^{i}\right) \frac{\partial}{\partial z^{i}}+\left(a^{i}-i b^{i}\right) \frac{\partial}{\partial \bar{z}^{i}}, \\
& I v=-b^{i} \frac{\partial}{\partial x^{i}}+a^{i} \frac{\partial}{\partial y^{i}}=\left(-b^{i}+i a^{i}\right) \frac{\partial}{\partial z^{i}}+\left(-b^{i}-i a^{i}\right) \frac{\partial}{\partial \bar{z}^{i}}= \\
&= i\left(a^{i}+i b^{i}\right) \frac{\partial}{\partial z^{i}}-i\left(a^{i}-i b^{i}\right) \frac{\partial}{\partial \bar{z}^{i}} .
\end{aligned}
$$

We are looking now for the vector fields $v$ (2.11) which are tangent to the hypersurfaces (2.10), the vector fields Iv (2.12) having the same property. This yields

$$
A^{i} \frac{\partial F}{\partial z^{i}}+\bar{A}^{i} \frac{\partial F}{\partial \bar{z}^{i}}=0, \quad i A^{i} \frac{\partial F}{\partial z^{i}}-i \overline{A^{i}} \frac{\partial F}{\partial \bar{z}^{i}}=0
$$

i.e.,

$$
\begin{equation*}
A^{i} \frac{\partial F}{\partial z^{i}}=0, \quad \overline{A^{i}} \frac{\partial F}{\partial \bar{z}^{i}}=0 \tag{2.13}
\end{equation*}
$$

Because of $F=\bar{F}$, we have

$$
\frac{\partial F}{\partial \bar{z}^{i}}=\frac{\overline{\partial F}}{\partial z^{i}} ;
$$

indeed, write $F\left(z^{i}, \bar{z}^{i}\right)=f\left(x^{1}, \ldots, x^{n}, y^{1}, \ldots, y^{n}\right)$, then

$$
\frac{\partial F}{\partial z^{i}}=\frac{1}{2}\left(\frac{\partial f}{\partial x^{i}}-i \frac{\partial f}{\partial y^{i}}\right), \frac{\partial F}{\partial \bar{z}^{i}}=\frac{1}{2}\left(\frac{\partial f}{\partial x^{i}}+i \frac{\partial f}{\partial y^{i}}\right) .
$$

Thus the system (2.13) is equivalent to

$$
\begin{equation*}
A^{i} \frac{\partial F}{\partial z^{i}}=0 \tag{2.14}
\end{equation*}
$$

It is easy to see that the coordinates $\boldsymbol{z}^{\boldsymbol{i}}$ in $\mathbf{C}^{\boldsymbol{n}}$ may be chosen in such a way (by a linear change) that

$$
\begin{align*}
& F\left(z^{1}, \ldots, z^{n}, \bar{z}^{1}, \ldots, \bar{z}^{n}\right)=z^{n}+\bar{z}^{n}+G\left(z^{1}, \ldots, z^{n-1}, \bar{z}^{1}, \ldots, \bar{z}^{n-1}, \bar{z}^{n}-z^{n}\right) ; \\
& G(0, \ldots, 0)=0 ;  \tag{2.15}\\
& \quad \frac{\partial G(0, \ldots, 0)}{\partial z^{a}}=0, \quad \frac{\partial G(0, \ldots, 0)}{\partial \bar{z}^{a}}=0, \quad \frac{\partial G(0, \ldots, 0)}{\partial\left(\bar{z}^{n}-z^{n}\right)}=0 \\
& \text { for } \alpha=1, \ldots, n-1 .
\end{align*}
$$

The geometrical meaning is very simple: The tangent hyperplane $T_{0}\left(X^{2 n-1}\right)$ at the origin is given by $z^{n}+\bar{z}^{n}=0$, i.e., $x^{n}=0$.

Of course, $\partial F\left(z^{i}, \bar{z}^{i}\right) / \partial z^{n} \neq 0$ in a neighbourhood of the origin, and (2.14) may be written as

$$
\begin{equation*}
A^{a} \frac{\partial F}{\partial z^{a}}+A^{n} \frac{\partial F}{\partial z^{n}}=0 \quad(\alpha, \beta, \ldots=1, \ldots, n-1) \tag{2.16}
\end{equation*}
$$

Its general solution is given by

$$
A^{a}=B^{a} \frac{\partial F}{\partial z^{n}}, \quad A^{n}=-B^{a} \frac{\partial F}{\partial z^{a}},
$$

$B^{1}, \ldots, B^{n-1}$ being arbitrary complex-valued functions, and we get

$$
\begin{gather*}
v=B^{a} \frac{\partial F}{\partial z^{n}} \frac{\partial}{\partial z^{\alpha}}-B^{a} \frac{\partial F}{\partial z^{\alpha}} \frac{\partial}{\partial z^{n}}+\overline{B^{\alpha}} \frac{\partial F}{\partial \bar{z}^{n}} \frac{\partial}{\partial \bar{z}^{\alpha}}-\overline{B^{\alpha}} \frac{\partial F}{\partial \bar{z}^{\alpha}} \frac{\partial}{\partial \bar{z}^{n}},  \tag{2.17}\\
I v=i B^{\alpha} \frac{\partial F}{\partial z^{n}} \frac{\partial}{\partial z^{\beta}}-i B^{\beta} \frac{\partial F}{\partial z^{\beta}} \frac{\partial}{\partial z^{n}}-i \bar{B}^{\beta} \frac{\partial F}{\partial \bar{z}^{n}} \frac{\partial}{\partial \bar{z}^{\beta}}+i \bar{B}^{\beta} \frac{\partial F}{\partial \bar{z}^{\beta}} \frac{\partial}{\partial \bar{z}^{n}} .
\end{gather*}
$$

At the origin of $\mathbf{C}^{n}$, we have

$$
\frac{\partial F(0, \ldots, 0)}{\partial z^{n}}=1, \quad \frac{\partial F(0, \ldots, 0)}{\partial \bar{z}^{n}}=1, \quad \frac{\partial F(0, \ldots, 0)}{\partial z^{a}}=0, \quad \frac{\partial F(0, \ldots, 0)}{\partial \bar{z}^{a}}=0,
$$

and the vectors (2.171) are given by

$$
\begin{equation*}
v=B^{a} \frac{\partial}{\partial z^{a}}+\overline{B^{a}} \frac{\partial}{\partial \bar{z}^{a}} . \tag{2.18}
\end{equation*}
$$

Thus the space $\tau_{0} \subset T_{0}\left(X^{2 n-1}\right)$ is spanned by the vectors

$$
\frac{\partial}{\partial z^{a}}+\frac{\partial}{\partial \bar{z}^{a}}, \quad i \frac{\partial}{\partial z^{a}}-i \frac{\partial}{\partial \bar{z}^{a}} ; \quad \alpha=1, \ldots, n-1 .
$$

From (2.17), we get

$$
\begin{aligned}
& {[v, I v]_{0}=B^{a} . i B^{\beta} \frac{\partial^{2} F(0)}{\partial z^{a} \partial z^{n}} \frac{\partial}{\partial z^{\beta}}-B^{a} . i B^{\beta} \frac{\partial^{2} F(0)}{\partial z^{a} \partial z^{\beta}} \frac{\partial}{\partial z^{n}}-} \\
& -B^{a} \cdot \overline{i B^{\beta}} \frac{\partial^{2} F(0)}{\partial z^{a} \partial \bar{z}^{n}} \frac{\partial}{\partial \bar{z}^{\beta}}+B^{a} \cdot i \bar{B}^{\beta} \frac{\partial^{2} F(0)}{\partial z^{a} \partial \bar{z}^{\beta}} \frac{\partial}{\partial \bar{z}^{n}}+\bar{B}^{a} \cdot i B^{\beta} \frac{\partial^{2} F(0)}{\partial \bar{z}^{a} \partial z^{n}} \frac{\partial}{\partial z^{\beta}}- \\
& -\bar{B}^{a} \cdot i B^{\beta} \frac{\partial^{2} F(0)}{\partial \bar{z}^{a} \partial z^{\beta}} \frac{\partial}{\partial z^{n}}-\bar{B}^{a} \cdot i \bar{B}^{\beta} \frac{\partial^{2} F(0)}{\partial \bar{z}^{\alpha} \partial \bar{z}^{n}} \frac{\partial}{\partial \bar{z}^{\beta}}+\bar{B}^{a} \cdot i \overline{B^{\beta}} \frac{\partial^{2} F(0)}{\partial \bar{z}^{\alpha} \partial \bar{z}^{\beta}} \frac{\partial}{\partial \bar{z}^{n}}- \\
& -i B^{\beta} B^{a} \frac{\partial^{2} F(0)}{\partial z^{\beta} \partial z^{n}} \frac{\partial}{\partial z^{\alpha}}+i B^{\beta} B^{\alpha} \frac{\partial^{2} F(0)}{\partial z^{\alpha} \partial z^{\beta}}-\frac{\partial}{\partial z^{n}}-i B^{\beta} \bar{B}^{\alpha} \frac{\partial^{2} F(0)}{\partial z^{\beta} \partial \bar{z}^{n}} \frac{\partial}{\partial \bar{z}^{\alpha}}+ \\
& +i B^{\bar{\beta}} \bar{B}^{\alpha} \frac{\partial^{2} F(0)}{\partial z^{\beta} \partial \bar{z}^{\alpha}} \frac{\partial}{\partial \bar{z}^{n}}+i \bar{B}^{\beta} B^{\alpha} \frac{\partial^{2} F(0)}{\partial \bar{z}^{\beta} \partial z^{n}} \frac{\partial}{\partial z^{\alpha}}-i \bar{B}^{\beta} B^{\alpha} \frac{\partial^{2} F(0)}{\partial z^{a} \partial \bar{z}^{\beta}} \frac{\partial}{\partial z^{n}}+ \\
& +i \bar{B}^{\beta} \bar{B}^{\alpha} \frac{\partial^{2} F(0)}{\partial \bar{z}^{\beta} \overline{\bar{z}}^{n}} \frac{\partial}{\partial \bar{z}^{\alpha}}-i \bar{B}^{\beta} \bar{B}^{\alpha} \frac{\partial^{2} F(0)}{\partial \bar{z}^{a} \partial \bar{z}^{\beta}} \frac{\partial}{\partial \bar{z}^{n}}=
\end{aligned}
$$

$$
\begin{gathered}
=2 i B^{a} \overline{B^{\beta}} \frac{\partial^{2} F(0)}{\partial \bar{z}^{\beta} \partial z^{n}} \frac{\partial}{\partial z^{\alpha}}-2 i B^{\beta} \overline{B^{\alpha}} \frac{\partial^{2} F(0)}{\partial z^{\beta} \partial \bar{z}^{n}} \frac{\partial}{\partial \bar{z}^{a}}- \\
-2 i \bar{B}^{a} B^{\beta} \frac{\partial^{2} F(0)}{\partial \bar{z}^{a} \partial z^{\beta}} \frac{\partial}{\partial z^{n}}+2 i B^{a} \bar{B}^{\beta} \frac{\partial^{2} F(0)}{\partial z^{a} \partial \bar{z}^{\beta}} \frac{\partial}{\partial \bar{z}^{n}}= \\
=i\left(B^{a} \bar{B}^{\beta} \frac{\partial^{2} F(0)}{\partial \bar{z}^{\beta} \partial z^{n}}-B^{\beta} \overline{B^{a}} \frac{\partial^{2} F(0)}{\partial z^{\beta} \partial \bar{z}^{n}}\right)\left(\frac{\partial}{\partial z^{a}}+\frac{\partial}{\partial \bar{z} \bar{a}}\right)+ \\
+\left(B^{a} \overline{B^{\beta}} \frac{\partial^{2} F(0)}{\partial \bar{z}^{\beta} \partial z^{n}}+B^{\beta} \overline{B^{\alpha}} \cdot \frac{\partial^{2} F(0)}{\partial z^{\beta} \partial \bar{z}^{n}}\right) \cdot i\left(\frac{\partial}{\partial z^{\alpha}}-\frac{\partial}{\partial \bar{z}^{\alpha}}\right)- \\
-2 i B^{a} \bar{B}^{\beta} \frac{\partial^{2} F(0)}{\partial z^{a} \partial \bar{z}^{\beta}}\left(\frac{\partial}{\partial z^{n}}-\frac{\partial}{\partial \bar{z}^{n}}\right) .
\end{gathered}
$$

Of course,

$$
\frac{\partial}{\partial y^{n}}=i\left(\frac{\partial}{\partial z^{n}}-\frac{\partial}{\partial \bar{z}^{n}}\right) .
$$

Consider once again the natural projection $\pi_{0}: \tau_{0} \rightarrow T_{0}\left(X^{2 n-1}\right) / \tau_{0}$, and write

$$
\pi_{0}\left(\frac{\partial}{\partial y^{n}}\right)=u
$$

then

$$
L_{0}\left(B^{a} \frac{\partial}{\partial z^{a}}+\bar{B}^{a} \frac{\partial}{\partial \bar{z}^{a}}\right)=-2 B^{a} \bar{B}^{\beta} \frac{\partial^{2} F(0)}{\partial z^{a} \partial \bar{z}^{\beta}} u
$$

This is the classical formula for the Levi map. It is easy to prove that $L_{x} \equiv 0$ at each point $x \in X^{2 n-1}$ is equivalent to the condition that $X^{2 n-1}$ is locally holomorphically equivalent to a hyperplane of $\mathbf{C}^{n}$.
3. Let us consider a manifold $M^{3}$ with the structure described at the end of No 1. At each point $m \in M^{3}$, let us choose a frame ( $v_{1}, v_{2}, v_{3}$ ), $v_{i} \in T_{m}\left(M^{3}\right)$, such that $\tau_{m}$ is spanned by $v_{1}, v_{2}$ and $I_{m} v_{1}=v_{2}$. Each other frame of the same type is given by

$$
\begin{gather*}
w_{1}=\alpha v_{1}-\beta v_{2}, \quad w_{2}=\beta v_{1}+\alpha v_{2}  \tag{3.1}\\
w_{3}=\gamma v_{1}+\delta v_{2}+\varphi v_{3} ; \quad\left(\alpha^{2}+\beta^{2}\right) \varphi \neq 0 .
\end{gather*}
$$

Let $v, v^{\prime} \in T_{m}\left(M^{3}\right)$,

$$
\begin{equation*}
v=a v_{1}+b v_{2}+c v_{3}, \quad v^{\prime}=a^{\prime} v_{1}+b^{\prime} v_{2}+c^{\prime} v_{3} . \tag{3.2}
\end{equation*}
$$

Then

$$
\begin{equation*}
\varphi^{\star}\left(v, v^{\prime}\right)=A\left(a b^{\prime}-a^{\prime} b\right)+B\left(a c^{\prime}-a^{\prime} c\right)+C\left(b c^{\prime}-b^{\prime} c\right) \tag{3.3}
\end{equation*}
$$

where $A, B, C$ are reals. For $v, v^{\prime} \in \tau_{m}$, we have $c=c^{\prime}=0$ and

$$
\varphi^{\star}\left(v, v^{\prime}\right)=A\left(a b^{\prime}-a^{\prime} b\right), \quad \varphi^{\star}\left(I v, I v^{\prime}\right)=A\left(a b^{\prime}-a^{\prime} b\right)
$$

From the condition $\varphi^{\star}\left(v_{m}, w_{m}\right)=-\varphi^{\star}\left(I_{m} v_{m}, I_{m} v_{m}\right)$, we get $A=0$. Let

$$
v=\tilde{a} w_{1}+\tilde{b} w_{2}+\tilde{c} w_{2}, \quad v^{\prime}=\tilde{a}^{\prime} w_{1}+\tilde{b}^{\prime} w_{2}+\tilde{c}^{\prime} w_{3},
$$

$w_{1}, w_{2}, w_{3}$ being given by (3.1). Write

$$
\varphi^{\star}\left(v, v^{\prime}\right)=\tilde{B}\left(\tilde{a} \tilde{c}^{\prime}-\tilde{a}^{\prime} \tilde{c}\right)+\tilde{C}\left(\tilde{b} \tilde{c}^{\prime}-\tilde{b}^{\prime} \tilde{c}\right)
$$

Then

$$
a=\alpha \tilde{a}+\beta \tilde{b}, \quad b=-\beta \tilde{a}+\alpha \tilde{b}, \quad c=\varphi \tilde{c}
$$

and

$$
\begin{equation*}
\tilde{B}=\varphi(\alpha B-\beta C), \quad \tilde{C}=\varphi(\beta B+\alpha C) \tag{3.4}
\end{equation*}
$$

The case $B=C=0$ being excluded (otherwise $\varphi^{\star} \equiv 0$ ), there exist frames ( $w_{1}, w_{2}, w_{3}$ ) with $\tilde{B}=1, \tilde{C}=0$, and we have the following result: On $M^{3}$, the considered structure induces a $G$-structure $B_{G}\left(M^{3}\right)$ such that ( $\left.v_{1}, v_{2}, v_{3}\right)_{m} \in B_{G}\left(M^{3}\right)$ if and only if $v_{1}, v_{2} \in \tau_{m}, I_{m} v_{1}=v_{2}$ and $\varphi^{\star}\left(v, v^{\prime}\right)=a c^{\prime}-a^{\prime} c, v$ and $v^{\prime}$ being given by (3.2); if $\left\{w_{1}, w_{2}, w_{3}\right\}_{m} \in B_{G}\left(M^{3}\right)$, then

$$
\begin{equation*}
w_{1}=\alpha v_{1}, \quad w_{2}=\alpha v_{2}, \quad w_{3}=\gamma v_{1}+\delta v_{2}+\alpha^{-1} v_{3} ; \quad \alpha \neq 0 . \tag{3.5}
\end{equation*}
$$

The last assertion follows easily from (3.4); indeed, we should have $1=\alpha \varphi, 0=\beta \varphi \varphi_{x}$.
Consider a $G$-structure $B_{G}\left(M^{3}\right)$ of this type, i.e., $G$ is the group of the matrices

$$
\left(\begin{array}{lll}
\alpha & 0 & 0  \tag{3.6}\\
0 & \alpha & 0 \\
\gamma & \delta & \alpha^{-1}
\end{array}\right), \alpha \neq 0 .
$$

In a domain $V \subset M^{3}$, choose a section $\left(v_{1}, v_{2}, v_{3}\right)$ of $B_{G}\left(M^{3}\right)$; then

$$
\begin{align*}
& {\left[v_{1}, v_{2}\right]=a_{1} v_{1}+a_{2} v_{2}+a_{3} v_{3}}  \tag{3.7}\\
& {\left[v_{1}, v_{3}\right]=b_{1} v_{1}+b_{2} v_{2}+b_{3} v_{3},} \\
& {\left[v_{2}, v_{3}\right]=c_{1} v_{1}+c_{2} v_{2}+c_{3} v_{3}}
\end{align*}
$$

$a_{1}, \ldots, c_{3}$ being functions on $V$. In what follows, let us restrict ourselves to manifolds with non-integrable field of planes $\tau_{m}$; thus $a_{3} \neq 0$ on $V$. From the Jacobi identity

$$
\left[v_{1},\left[v_{2}, v_{3}\right]\right]+\left[v_{2},\left[v_{3}, v_{1}\right]\right]+\left[v_{3},\left[v_{1}, v_{2}\right]\right]=0
$$

we get

$$
\begin{align*}
& v_{1} c_{1}-v_{2} b_{1}+v_{3} a_{1}+a_{1} c_{2}+b_{1} c_{3}-b_{3} c_{1}-a_{2} c_{1}=0  \tag{3.8}\\
& v_{1} c_{2}-v_{2} b_{2}+v_{3} a_{2}+b_{2} c_{3}+a_{2} b_{1}-b_{3} c_{2}-a_{1} b_{2}=0 \\
& v_{1} c_{3}-v_{2} b_{3}+v_{3} a_{3}+a_{3} c_{2}+a_{3} b_{1}-a_{1} b_{3}-a_{2} c_{3}=0
\end{align*}
$$

Let ( $w_{1}, w_{2}, w_{3}$ ) be another section of $B_{G}\left(M^{3}\right)$, let us have (3.5) with $\alpha, \gamma, \delta$ real-valued functions on $V$. Then

$$
\begin{align*}
& {\left[w_{1}, w_{2}\right]=A_{1} w_{1}+A_{2} w_{2}+A_{3} w_{3}}  \tag{3.9}\\
& {\left[w_{1}, w_{3}\right]=B_{1} w_{1}+B_{2} w_{2}+B_{3} w_{3}} \\
& {\left[w_{2}, w_{3}\right]=C_{1} w_{1}+C_{2} w_{2}+C_{3} w_{3}}
\end{align*}
$$

We have

$$
\begin{gather*}
{\left[w_{1}, w_{2}\right]=\left[\alpha v_{1}, \alpha v_{2}\right]=\alpha \cdot v_{1} \alpha \cdot v_{2}-\alpha \cdot v_{2} \alpha \cdot v_{1}+\alpha^{2}\left(a_{1} v_{1}+a_{2} v_{2}+a_{3} v_{3}\right)=} \\
=A_{1} \alpha v_{1}+A_{2} \alpha v_{2}+A_{3}\left(\gamma v_{1}+\delta v_{2}+\alpha^{-1} v_{3}\right), \tag{3.10}
\end{gather*}
$$

i.e.
$-\alpha \cdot v_{2} \alpha+\alpha^{2} a_{1}=\alpha A_{1}+\gamma A_{3}, \alpha \cdot v_{1} \alpha+\alpha^{2} a_{2}=\alpha A_{2}+\delta A_{3}, \alpha^{2} a_{3}=\alpha^{-1} A_{3}$.

Thus there exists a section ( $w_{1}, w_{2}, w_{3}$ ) satisfying $A_{3}=1, A_{1}=A_{2}=0$, and we have the following result: There exists (locally) exactly one section ( $v_{1}, v_{2}, v_{3}$ ) of $B_{G}\left(M^{3}\right)$ satisfying

$$
\begin{align*}
& {\left[v_{1}, v_{2}\right]=}  \tag{3.11}\\
& {\left[v_{1}, v_{3}\right]=b_{1} v_{1}+b_{2} v_{2}+b_{3} v_{3}} \\
& {\left[v_{2}, v_{3}\right]=c_{1} v_{1}+c_{2} v_{2}+c_{3} v_{3}}
\end{align*}
$$

The integrability conditions (3.8) reduce to

$$
\begin{align*}
& v_{1} c_{1}-v_{2} b_{1}+b_{1} c_{3}-b_{3} c_{1}=0  \tag{3.12}\\
& v_{1} c_{2}-v_{2} b_{2}+b_{2} c_{3}-b_{3} c_{2}=0 \\
& v_{1} c_{3}-v_{2} b_{3}+c_{2}+b_{1}=0
\end{align*}
$$

Now, let $B_{G}\left(M^{3}\right)$ be transitive. Then $b_{1}, \ldots, c_{3}$ are constants, and the equations (3.12) reduce to

$$
\begin{equation*}
b_{1} c_{3}-b_{3} c_{1}=0, \quad b_{2} c_{3}-b_{3} c_{2}=0, \quad c_{2}+b_{1}=0 \tag{3.13}
\end{equation*}
$$

Let $b_{3} c_{3} \neq 0$. Then there are real numbers $A, B, C$ such that

$$
\begin{align*}
& {\left[v_{1}, v_{2}\right]=}  \tag{3.14}\\
& {\left[v_{1}, v_{3}\right]=A B C v_{1}-A B^{2} v_{2}+B v_{3},} \\
& {\left[v_{2}, v_{3}\right]=A C^{2} v_{1}-A B C v_{2}+C v_{3} ; \quad B C \neq 0}
\end{align*}
$$

Let $b_{3} \neq 0, c_{3}=0$. Then $c_{1}=c_{2}=b_{1}=0$ and (3.11) are of the form

$$
\begin{align*}
& {\left[v_{1}, v_{2}\right]=}  \tag{3.15}\\
& {\left[v_{1}, v_{3}\right]=A v_{2}+B v_{3},} \\
& {\left[v_{2}, v_{3}\right]=0 ;}
\end{align*} \quad B \neq 0 ;
$$

the case $b_{3}=0, c_{3} \neq 0$ is symmetric. For $b_{3}=c_{3}=0$, we get

$$
\begin{align*}
& {\left[v_{1}, v_{2}\right]=-v_{3}}  \tag{3.16}\\
& {\left[v_{1}, v_{3}\right]=A v_{1}+B v_{2}} \\
& {\left[v_{2}, v_{3}\right]=C v_{1}-A v_{2}}
\end{align*}
$$

The following result follows: The Lie algebra of $G$ (see the Theorem) is of the type (3.14) or (3.15) or (3.16) resp.

Finally, let us prove the existence of the transitive $G$-structures of the types (3.14)- (3.16). A simple check shows that the vector fields

$$
\begin{align*}
& u_{1}=\frac{1}{2}\left(1+2 y-3 x^{2}\right) \frac{\partial}{\partial x}+\frac{1}{2}(2 x+z-3 x y) \frac{\partial}{\partial y}+\frac{3}{2}(y-x z) \frac{\partial}{\partial z}, \\
& u_{2}=\frac{1}{2}\left(1-2 y+3 x^{2}\right) \frac{\partial}{\partial x}+\frac{1}{2}(2 x-z+3 x y) \frac{\partial}{\partial y}+\frac{3}{2}(y+x z) \frac{\partial}{\partial z}, \\
& u_{3}=x \frac{\partial}{\partial x}+2 y \frac{\partial}{\partial y}+3 z \frac{\partial}{\partial z} \tag{3.17}
\end{align*}
$$

on $\mathbf{R}^{\mathbf{3}}$ satisfy

$$
\begin{equation*}
\left[u_{1}, u_{2}\right]=u_{3}, \quad\left[u_{1}, u_{3}\right]=u_{2}, \quad\left[u_{2}, u_{3}\right]=u_{1} \tag{3.18}
\end{equation*}
$$

In a suitable neighbourhood of the point $\left(\frac{1}{4} \pi, 0,0\right) \in \mathbf{R}^{3}$, consider the vector fields

$$
\begin{align*}
& w_{1}=\sin (y+z) \frac{\partial}{\partial x}+\frac{\cos x}{\sin x} \cos (y+z) \frac{\partial}{\partial y}-\frac{\sin x}{\cos x} \cos (y+z) \frac{\partial}{\partial z} \\
& w_{2}=\cos (y+z) \frac{\partial}{\partial x}-\frac{\cos x}{\sin x} \sin (y+z) \frac{\partial}{\partial y}+\frac{\sin x}{\cos x} \sin (y+z) \frac{\partial}{\partial z} \\
& w_{3}=\frac{\partial}{\partial y}+\frac{\partial}{\partial z} \tag{3.19}
\end{align*}
$$

the direct check proves

$$
\begin{equation*}
\left[w_{1}, w_{2}\right]=2 w_{3}, \quad\left[w_{1}, w_{3}\right]=-2 w_{2}, \quad\left[w_{2}, w_{3}\right]=2 w_{1} \tag{3.20}
\end{equation*}
$$

Now, consider the $G$-structure (3.14). Obviously, $\left[C v_{1}-B v_{2}, v_{3}\right]=0$. On a neighbourhood of a point $m_{0} \in M^{3}$, consider local coordinates $(x, y, z)$ such that

$$
C v_{1}-B v_{2}=\frac{\partial}{\partial y}, \quad v_{3}=\frac{\partial}{\partial x}
$$

this being always possible. Let

$$
v_{2}=\alpha \frac{\partial}{\partial x}+\beta \frac{\partial}{\partial y}+\gamma \frac{\partial}{\partial z}, \quad \text { i.e., } \quad C v_{1}=B \alpha \frac{\partial}{\partial x}+(B \beta+1) \frac{\partial}{\partial y}+B \gamma \frac{\partial}{\partial z} .
$$

From (3.14 1,2 ), we get

$$
\frac{\partial \alpha}{\partial y}=C ; \quad \frac{\partial \beta}{\partial y}=0, \quad \frac{\partial \gamma}{\partial y}=0, \quad \frac{\partial \alpha}{\partial x}=-C, \quad \frac{\partial \beta}{\partial x}=-A C, \quad \frac{\partial \gamma}{\partial x}=0
$$

Consider the particular solution $\alpha=C(y-x), \beta=-A C x, \gamma=1$. Then

$$
\begin{align*}
& v_{1}=B(y-x) \frac{\partial}{\partial x}+\left(C^{-1}-A B x\right) \frac{\partial}{\partial y}+B C^{-1} \frac{\partial}{\partial z}  \tag{3.21}\\
& v_{2}=C(y-x) \frac{\partial}{\partial x}-A C x \frac{\partial}{\partial y}+\frac{\partial}{\partial z} \\
& v_{3}=\frac{\partial}{\partial x}
\end{align*}
$$

this vectors being linearly independent and satisfying (3.14), they generate a $G$-structure of the type (3.14) on $\mathbf{R}^{3}$. Similarly, the vector fields

$$
\begin{equation*}
v_{1}=-(B x+y) \frac{\partial}{\partial x}-A x \frac{\partial}{\partial y}+\frac{\partial}{\partial z}, \quad v_{2}=\frac{\partial}{\partial y}, \quad v_{3}=\frac{\partial}{\partial x} \tag{3.22}
\end{equation*}
$$

generate a $G$-structure of the type (3.15) on $\mathbf{R}^{3}$. The type (3.16) is a little more complicated. First of all, suppose $A=B=0$; the $G$-structure of this type on $\mathbf{R}^{3}$ is generated by the vector fields

$$
\begin{equation*}
v_{1}=\frac{\partial}{\partial x}, \quad v_{2}=-C y \frac{\partial}{\partial x}+x \frac{\partial}{\partial y}+\frac{\partial}{\partial z}, \quad v_{3}=\frac{\partial}{\partial y} \tag{3.23}
\end{equation*}
$$

Similarly, the $G$-structure of the type (3.16) with $A=C=0$ is generated by the vector fields

$$
\begin{equation*}
v_{1}=-B y \frac{\partial}{\partial x}-x \frac{\partial}{\partial y}+\frac{\partial}{\partial z}, \quad v_{2}=\frac{\partial}{\partial x}, \quad v_{3}=\frac{\partial}{\partial y} . \tag{3.24}
\end{equation*}
$$

Now, consider the case $A^{2}+B C=0, A B \neq 0$, i.e.,

$$
\left[v_{1}, v_{2}\right]=v_{3}, \quad\left[v_{1}, v_{3}\right]=A v_{1}+B v_{2}, \quad\left[v_{2}, v_{3}\right]=-\frac{A^{2}}{B} v_{1}-A v_{2}
$$

We see that $\left[A v_{1}+B v_{2}, v_{3}\right]=0$, and the vector fields

$$
\begin{align*}
& v_{1}=-B y \frac{\partial}{\partial x}-x \frac{\partial}{\partial y}+\frac{\partial}{\partial z}  \tag{3.25}\\
& v_{2}=A B y \frac{\partial}{\partial x}+(1+A x) \frac{\partial}{\partial y}-A \frac{\partial}{\partial z} \\
& v_{3}=\frac{\partial}{\partial x}
\end{align*}
$$

generate the $G$-structure of this type on $\mathbf{R}^{3}$. If $A^{2}+B C \neq 0$ then the Lie algebra (3.16) $L$ satisfies $[L, L]=L$ and it contains a basis ( $u_{1}, u_{2}, u_{3}$ ) satisfying (3.18) or a basis ( $w_{1}, w_{2}, w_{3}$ ) satisfying (3.20).
4. Consider the space $\mathbf{C}^{2}$ and the pseudogroup $\Gamma$. The relation between the 1-parametric local subgroups of $\Gamma$ and the holomorphic vector fields on $\mathbf{C}^{2}$ is well known. Let

$$
\begin{equation*}
v=a(x, y) \frac{\partial}{\partial x}+b(x, y) \frac{\partial}{\partial y} \tag{4.1}
\end{equation*}
$$

be a (locally defined) holomorphic vector field; the corresponding local group $G_{v}$ consists of the maps
given by

$$
\begin{equation*}
\varphi_{t}: \quad \tilde{x}=f(x, y, t), \quad \tilde{y}=g(x, y, t), \quad t \in(-\varepsilon, \varepsilon) \tag{4.2}
\end{equation*}
$$

$$
\begin{gather*}
\frac{\partial f(x, y, t)}{\partial x}=a(f(x, y, t), g(x, y, t)), \frac{\partial g(x, y, t)}{\partial t}=b(f(x, y, t), g(x, y, t)), \\
f(x, y, 0)=x, \quad g(x, y, 0)=y \tag{4.3}
\end{gather*}
$$

We have $G_{v} \subset \Gamma_{\delta}$ if and only if

$$
\begin{equation*}
\frac{\partial a(x, y)}{\partial x}+\frac{\partial b(x, y)}{\partial y}=0 . \tag{4.4}
\end{equation*}
$$

Indeed, let us write

$$
D(x, y, t)=\frac{\partial f}{\partial x} \frac{\partial g}{\partial y}-\frac{\partial f}{\partial y} \frac{\partial g}{\partial x}
$$

We have $D(x, y, 0)=1$. From (4.3), we get

$$
\frac{\partial D}{\partial t}=\left(\frac{\partial a}{\partial x}+\frac{\partial b}{\partial y}\right) D
$$

and the result follows easily. Denote by $L_{s}$ the Lie algebra of holomorphic vector fields (4.1) on $\mathbf{C}^{2}$ satisfying (4.4).

Let $w_{1}, w_{2} \in L_{s}, w_{1} \neq 0 \neq w_{2},\left[w_{1}, w_{2}\right]=0$; then there are (locally) $\Gamma_{8}$-coordinates $(u, v)$ such that

$$
\begin{equation*}
w_{1}=\frac{\partial}{\partial u}, w_{2}=\alpha \frac{\partial}{\partial v}(0 \neq \alpha \in \mathbf{C}) \text { or } w_{2}=a(v) \frac{\partial}{\partial u} r e s p . \tag{4.5}
\end{equation*}
$$

Here, the $\Gamma_{s}$-coordinates ( $u, v$ ) are defined (locally) as holomorphic coordinates $u=u(x, y), v=v(x, y)$ satisfying $\partial(u, v) / \partial(x, y)=1$. Indeed, we may choose (at least locally) $\Gamma_{s}$-coordinates $r=r(x, y), s=s(x, y)$ such that $w_{1}=\partial / \partial r$. Let

$$
w_{2}=b(r, s) \frac{\partial}{\partial r}+c(r, s) \frac{\partial}{\partial s}, \frac{\partial b}{\partial r}+\frac{\partial c}{\partial s}=0 .
$$

From $\left[w_{1}, w_{2}\right]=0$, we get

$$
\frac{\partial b}{\partial r}=0, \quad \frac{\partial c}{\partial r}=0
$$

Thus $b=b(s), c=\alpha \in \mathbf{C}$. Now, consider the $\Gamma_{s}$-coordinates $u=u(r, s), v=v(r, s)$. Then

$$
\begin{aligned}
& w_{1}=\frac{\partial u}{\partial r} \frac{\partial}{\partial u}+\frac{\partial v}{\partial r} \frac{\partial}{\partial v}, \\
& w_{2}=b(s)\left(\frac{\partial u}{\partial r} \frac{\partial}{\partial u}+\frac{\partial v}{\partial r} \frac{\partial}{\partial v}\right)+\alpha\left(\frac{\partial u}{\partial s} \frac{\partial}{\partial u}+\frac{\partial v}{\partial s} \frac{\partial}{\partial v}\right) .
\end{aligned}
$$

We have

$$
\frac{\partial u}{\partial r}=1, \quad \frac{\partial v}{\partial r}=0 \quad \text { and } \quad \frac{\partial v}{\partial s}=1
$$

i.e., $u=r+g(s), \quad v=s+\varrho, \quad \varrho \in \mathbf{C}, \quad$ and

$$
w_{2}=\left(b+\alpha \frac{d g}{d s}\right) \frac{\partial}{\partial u}+\alpha \frac{\partial}{\partial v} .
$$

If $\alpha \neq 0$, let us choose $g(s)$ such that

$$
\frac{d g(s)}{d s}=-\frac{b(s)}{\alpha} .
$$

5. Let $L$ be a Lie algebra of the type (3.14), suppose $L \subset L_{8}$. Then

$$
\left[v_{2}-\frac{C}{B} v_{1}, v_{3}\right]=0
$$

and we may choose (locally) $\Gamma_{s}$-coordinates $(u, v)$ such that

$$
\begin{equation*}
v_{3}=\frac{\partial}{\partial u}, \quad v_{2}-\frac{C}{B} v_{1}=\alpha \frac{\partial}{\partial v} ; \quad 0 \neq \alpha \in \mathbf{C} \tag{5.1}
\end{equation*}
$$

or

$$
\begin{equation*}
v_{3}=\frac{\partial}{\partial u}, \quad v_{2}-\frac{C}{B} v_{1}=a(v) \frac{\partial}{\partial u} \tag{5.2}
\end{equation*}
$$

resp. Suppose (5.1) and

$$
v_{1}=b(u, v) \frac{\partial}{\partial u}+c(u, v) \frac{\partial}{\partial v}, \quad \frac{\partial b}{\partial u}+\frac{\partial c}{\partial v}=0 .
$$

From (3.142), we get

$$
-\frac{\partial b}{\partial u}=B, \quad \frac{\partial c}{\partial u}=A B^{2} \alpha,
$$

i.e.,

$$
b=-B u+\beta(v), \quad c=A B^{2} \alpha u+B v+\gamma_{0} ; \quad \gamma_{0} \in \mathbf{C} .
$$

From (3.141),

$$
\left[v_{1}, \frac{C}{B} v_{1}+\alpha \frac{\partial}{\partial v}\right]=-\alpha \frac{\partial b}{\partial v} \frac{\partial}{\partial u}-\alpha \frac{\partial c}{\partial v} \frac{\partial}{\partial v}=\frac{\partial}{\partial u},
$$

i.e., $\partial c / \partial v=B=0$. Thus we should have (5.2) because of $B \neq 0$. Let further $L$ be of the type (3.15) and $L \subset L_{s}$. Then there are (locally) $\Gamma_{s}$-coordinates ( $u, v$ ) such that

$$
\begin{equation*}
v_{3}=\frac{\partial}{\partial u}, \quad v_{2}=\alpha \frac{\partial}{\partial v} ; \quad 0 \neq \alpha \in \mathbf{C} \tag{5.3}
\end{equation*}
$$

or

$$
\begin{equation*}
v_{3}=\frac{\partial}{\partial u}, \quad v_{2}=a(v) \frac{\partial}{\partial u} \tag{5.4}
\end{equation*}
$$

resp. Suppose (5.3) and let us write

$$
v_{1}=b(u, v) \frac{\partial}{\partial u}+c(u, v) \frac{\partial}{\partial v}, \frac{\partial b}{\partial u}+\frac{\partial c}{\partial v}=0 .
$$

From (3.151,2),

$$
\alpha \frac{\partial b}{\partial v}=-1, \quad \frac{\partial c}{\partial v}=0, \quad \frac{\partial b}{\partial u}=-B, \quad \frac{\partial c}{\partial u}=-A \alpha
$$

Because of $B \neq 0$, we have (5.4).
Now, let $M^{3} \subset \mathbf{C}^{2} \equiv \mathbf{R}^{4}$ be the orbit of the group $G \subset \Gamma_{s}$ such that its Lie algebra $g$ is of the type (3.14) or (3.15) resp. Then we have shown that $g$ contains (in suitable $\Gamma_{s}$-coordinates) the vector fields $\partial / \partial x, a(y) \partial / \partial x$, and the vector fields

$$
\frac{\partial}{\partial x^{1}}, a_{1}\left(x^{2}, y^{2}\right) \frac{\partial}{\partial x^{1}}+a_{2}\left(x^{2}, y^{2}\right) \frac{\partial}{\partial y^{1}} ; \quad a(y)=a_{1}\left(x^{2}, y^{2}\right)+i a_{2}\left(x^{2}, y^{2}\right) ;
$$

are tangent to $M^{3} \subset \mathbf{R}^{4}$. The plane $\tau_{m}$ is thus spanned by the vectors $\partial / \partial x^{1}, \partial / \partial y^{1}$, and the field $\tau_{m}$ is integrable. The groups $G \subset \Gamma_{s}$ satisfying the suppositions of the Theorem and possessing the Lie algebra of the type (3.14) or (3.15) do not exist.
6. Let us investigate the case $L \subset L_{s}, L$ being of the type (3.16). Suppose $\operatorname{dim}[L, L]=1$, i.e.,

$$
\begin{equation*}
\left[v_{1}, v_{2}\right]=v_{3}, \quad\left[v_{1}, v_{3}\right]=0, \quad\left[v_{2}, v_{3}\right]=0 \tag{6.1}
\end{equation*}
$$

We may suppose the existence of $\Gamma_{s}$-coordinates $(u, v)$ such that

$$
v_{2}=\alpha \frac{\partial}{\partial v}, \quad v_{3}=\frac{\partial}{\partial u} ; \quad 0 \neq \alpha \in C .
$$

Let

$$
v_{1}=b(u, v) \frac{\partial}{\partial u}+c(u, v) \frac{\partial}{\partial v}, \quad \frac{\partial b}{\partial u}+\frac{\partial c}{\partial v}=0 .
$$

From ( $6.1_{1,2}$ ), we get

$$
\frac{\partial b}{\partial u}=0, \quad \frac{\partial b}{\partial v}=\frac{1}{\alpha}, \quad \frac{\partial c}{\partial u}=\frac{\partial c}{\partial v}=0
$$

i.e.,

$$
b=-\frac{v}{\alpha}+\beta, \quad c=\gamma ; \quad \beta, \gamma \in \mathbf{C}
$$

we have $\gamma \neq 0$ because of the non-integrability of the field $\tau_{m}$. Consider the $\Gamma_{s}$-coordinates $x=u, y=v-\alpha \beta$. Then

$$
v_{2}=\alpha \frac{\partial}{\partial y}, \quad v_{3}=\frac{\partial}{\partial x}, \quad v_{1}=-\frac{v}{\alpha} \frac{\partial}{\partial x}+\gamma \frac{\partial}{\partial y},
$$

and the general element of $L$ is

$$
\begin{equation*}
v=R\left(-\frac{v}{\alpha} \frac{\partial}{\partial x}+\gamma \frac{\partial}{\partial y}\right)+S \alpha \frac{\partial}{\partial y}+T \frac{\partial}{\partial x} ; \quad R, S, T \in \mathbf{R} . \tag{6.2}
\end{equation*}
$$

The associated local group $G_{v}$ is given by (4.3), i.e.,

$$
\frac{\partial f}{\partial t}=-\frac{R}{\alpha} g+T, \frac{\partial g}{\partial t}=R \gamma+S \alpha
$$

It is easy to see that its finite equations are

$$
f=x-\frac{R t}{\alpha} y-\frac{1}{2} R S t^{2}-\frac{1}{2} \frac{\gamma}{\alpha} R^{2} t^{2}+T t, \quad g=y+\gamma R t+\alpha S t
$$

Write $R t=a, \quad S t=b, \quad T t=c ; \quad$ we get

$$
\begin{equation*}
f=x-\frac{a}{\alpha} y-\frac{1}{2} a b-\frac{1}{2} \frac{\gamma}{\alpha} a^{2}+c, \quad g=y+\gamma a+\alpha b \tag{6.3}
\end{equation*}
$$

Thus

$$
\bar{f}=\bar{x}-\frac{a}{\bar{\alpha}} \bar{y}-\frac{1}{2} a b-\frac{1}{2} \frac{\bar{\gamma}}{\bar{\alpha}} a^{2}+c, \quad \bar{g}=\bar{y}+\bar{\gamma} a+\bar{\alpha} b,
$$

i.e.,

$$
\begin{gathered}
f-\bar{f}=x-\bar{x}-a\left(\frac{y}{\alpha}-\frac{\bar{y}}{\bar{\alpha}}\right)-\frac{1}{2} a^{2}\left(\frac{\gamma}{\alpha}-\frac{\bar{\gamma}}{\bar{\alpha}}\right), \\
\frac{g}{\alpha}-\frac{\bar{g}}{\bar{\alpha}}=\frac{y}{\alpha}-\frac{\bar{y}}{\bar{\alpha}}+a\left(\frac{\gamma}{\alpha}-\frac{\bar{\gamma}}{\bar{\alpha}}\right)
\end{gathered}
$$

the elimination of $a$ yields

$$
\left(\frac{y}{\alpha}-\frac{\bar{y}}{\bar{\alpha}}\right)^{2}+2\left(\frac{\gamma}{\alpha}-\frac{\bar{\gamma}}{\bar{\alpha}}\right)(x-\bar{x})=\left(\frac{g}{\alpha}-\frac{\bar{g}}{\bar{\alpha}}\right)^{2}+2\left(\frac{\gamma}{\alpha}-\frac{\bar{\gamma}}{\bar{\alpha}}\right)(f-\bar{f}),
$$

and we get the type (I).
Let us investigate the case $L \subset L_{8}, L$ being of the type (3.16) with $\operatorname{dim}[L, L]=2$. Then $A^{2}+B C=0$. First of all, suppose $A=B=0$, the case $A=C=0$ being symmetric. The algebra $L$ is of the type

$$
\begin{equation*}
\left[v, v_{2}\right]=v_{3}, \quad\left[v_{1}, v_{3}\right]=B v_{2}, \quad\left[v_{2}, v_{3}\right]=0 ; \quad B \neq 0 . \tag{6.4}
\end{equation*}
$$

In $\mathbf{C}^{2}$, there are $\Gamma_{s}$-coordinates $(u, v)$ such that

$$
v_{2}=\alpha \frac{\partial}{\partial v}, \quad v_{3}=\frac{\partial}{\partial u} ; \quad \alpha \neq 0 .
$$

Let

$$
v_{1}=b(u, v) \frac{\partial}{\partial u}+c(u, v) \frac{\partial}{\partial v}, \quad \frac{\partial b}{\partial u}+\frac{\partial c}{\partial v}=0 .
$$

From ( $6.4_{1}, 2$ ), we get

$$
\frac{\partial b}{\partial v}=-\frac{1}{\alpha}, \quad \frac{\partial c}{\partial v}=0, \quad \frac{\partial b}{\partial u}=0, \quad \frac{\partial c}{\partial u}=-\alpha B,
$$

i.e.,

$$
v_{1}=\left(-\frac{v}{\alpha}+b_{0}\right) \frac{\partial}{\partial u}+\left(-\alpha B u+c_{0}\right) \frac{\partial}{\partial v} ; \quad b_{0}, c_{0} \in \mathbf{C} .
$$

Consider the $\Gamma_{s}$-coordinates $x=u-c_{0} \alpha^{-1} B^{-1}, y=v-\alpha b_{0}$. Then

$$
v_{1}=-\frac{y}{\alpha} \frac{\partial}{\partial x}-\alpha B x \frac{\partial}{\partial y}, \quad v_{2}=\alpha \frac{\partial}{\partial y}, \quad v_{3}=\frac{\partial}{\partial x} .
$$

The general element $v \in L$ is

$$
\begin{equation*}
v=R\left(-\frac{y}{\alpha} \frac{\partial}{\partial x}-\alpha B x \frac{\partial}{\partial y}\right)+S \alpha \frac{\partial}{\partial y}+T \frac{\partial}{\partial x} ; R, S, T \in \mathbf{R} ; \tag{6.5}
\end{equation*}
$$

and the local group $G_{v}$ is given by

$$
\frac{\partial f}{\partial t}=-\frac{R}{\alpha} g+T, \quad \frac{\partial \varrho}{\partial t}=-R \alpha B f+S \alpha .
$$

Consider the group

$$
\begin{gather*}
f=a x-\frac{1}{\alpha} b y+c, \quad g=-a B b x+a y+\alpha d ;  \tag{6.6}\\
a, b, c, d \in \mathbf{R}, \quad a^{2}-B b^{2}=1 .
\end{gather*}
$$

We get its identity for $a=1, b=c=0$. Let $a(t), b(t), c(t), d(t)$ be its one-parametric subgroup $G_{1}$, let $t=0$ correspond to its identity. Then

$$
a \frac{\mathrm{~d} a}{\mathrm{~d} t}-B b \frac{\mathrm{~d} b}{\mathrm{~d} t}=0, \quad \frac{\mathrm{~d} a(0)}{\mathrm{d} t}=0 .
$$

The vector field

$$
v=\left(-\frac{1}{\alpha} \frac{\mathrm{~d} b(0)}{\mathrm{d} t} y+\frac{\mathrm{d} c(0)}{\mathrm{d} t}\right) \frac{\partial}{\partial x}+\left(-\alpha B \frac{\mathrm{~d} b(0)}{\mathrm{d} t} x+\alpha \frac{\mathrm{d} d(0)}{\mathrm{d} t}\right) \frac{\partial}{\partial y}
$$

being associated to $G_{1}$, we see that (6.6) corresponds to (6.5). We have

$$
\begin{gathered}
\bar{f}=a \bar{x}-\frac{1}{\bar{\alpha}} b \bar{y}+c, \quad \bar{g}=-\bar{\alpha} B b \bar{x}+a \bar{y}+\bar{\alpha} d, \\
f-\bar{f}=a(x-\bar{x})-b\left(\frac{y}{\alpha}-\frac{\bar{y}}{\bar{\alpha}}\right), \quad \bar{\alpha} g-\alpha \bar{g}=-\alpha \bar{\alpha} B b(x-\bar{x})+a(\bar{\alpha} y-\alpha \bar{y})
\end{gathered}
$$

and

$$
B(f-\bar{f})^{2}-\left(\frac{g}{\alpha}-\frac{\bar{g}}{\bar{\alpha}}\right)^{2}=B(x-\bar{x})^{2}-\left(\frac{y}{\alpha}-\frac{\bar{y}}{\bar{\alpha}}\right)^{2}
$$

Thus we have obtained the type (II).
Now, let $L$ be of the type (3.16) with $A^{2}+B C=0, A B \neq 0$, i.e.,

$$
\begin{equation*}
\left[v_{1}, v_{2}\right]=v_{3}, \quad\left[v_{1}, v_{3}\right]=A v_{1}+B v_{2}, \quad\left[v_{2}, v_{3}\right]=-\frac{A^{2}}{B} v_{1}-A v_{2} \tag{6.7}
\end{equation*}
$$

Then $\left[A v_{1}+B v_{2}, v_{3}\right]=0$, and there are $\Gamma_{s}$-coordinates $(u, v)$ such that

$$
\begin{aligned}
& A v_{1}+B v_{2}=\alpha \frac{\partial}{\partial v} \quad(0 \neq \alpha \in \mathbf{C}), \quad v_{3}=\frac{\partial}{\partial u} \\
& v_{1}=b(u, v) \frac{\partial}{\partial u}+c(u, v) \frac{\partial}{\partial v}, \quad \frac{\partial b}{\partial u}+\frac{\partial c}{\partial v}=0
\end{aligned}
$$

We have

$$
v_{2}=-\frac{A}{B} b \frac{\partial}{\partial u}+\frac{1}{B}(\alpha-A c) \frac{\partial}{\partial v}
$$

from (6.7 $7_{1,2}$ )

$$
\frac{\partial b}{\partial v}=-\frac{B}{\alpha}, \quad \frac{\partial c}{\partial v}=0, \quad \frac{\partial b}{\partial u}=0, \quad \frac{\partial c}{\partial u}=-\alpha
$$

i.e.,

$$
v_{1}=\left(-\frac{B}{\alpha} v+b_{0}\right) \frac{\partial}{\partial u}+\left(-\alpha u+c_{0}\right) \frac{\partial}{\partial v} ; \quad b_{0}, c_{0} \in \mathbf{C} .
$$

In the $\Gamma_{s}$-coordinates

$$
x=u-\frac{c_{0}}{\alpha}, \quad y=v-\frac{b_{0}}{\alpha} B
$$

we get

$$
A v_{1}+B v_{2}=\alpha \frac{\partial}{\partial y}, \quad v_{3}=\frac{\partial}{\partial x}, \quad v_{1}=-\frac{B}{\alpha} y \frac{\partial}{\partial x}-\alpha x \frac{\partial}{\partial y}
$$

The general element $v \in L$ being

$$
\begin{equation*}
v=R\left(\frac{y}{\alpha} \frac{\partial}{\partial x}+\frac{\alpha}{B} x \frac{\partial}{\partial y}\right)+S \alpha \frac{\partial}{\partial y}+T \frac{\partial}{\partial x} ; R, S, T \in \mathbf{R} \tag{6.8}
\end{equation*}
$$

we do not obtain now groups-compare (6.8) with (6.5).
7. Above we have considered all possibilities for $L \subset L_{8}$ with $\operatorname{dim}[L, L]<3$. Now, there are exactly two Lie algebras (over $\mathbf{R}$ ) with $\operatorname{dim} L=\operatorname{dim}[L, L]=3$ :

$$
\begin{equation*}
\left[w_{1}, w_{2}\right]=w_{3}, \quad\left[w_{1}, w_{3}\right]=-w_{2}, \quad\left[w_{2}, w_{3}\right]=w_{1} \tag{7.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\left[w_{1}, w_{2}\right]=w_{3}, \quad\left[w_{1}, w_{3}\right]=w_{2}, \quad\left[w_{2}, w_{3}\right]=w_{1} . \tag{7.2}
\end{equation*}
$$

First of all, let us consider the Lie algebra $L$ (7.2). The change $v_{1}=v_{3}, v_{2}=$ $=w_{2}-w_{1}, v_{3}=w_{2}+w_{1}$ of its basis yields

$$
\begin{equation*}
\left[v_{1}, v_{2}\right]=v_{2}, \quad\left[v_{1}, v_{3}\right]=-v_{3}, \quad\left[v_{2}, v_{3}\right]=-2 v_{1} . \tag{7.3}
\end{equation*}
$$

In $\mathbf{C}^{2}$, there are $\Gamma_{s}$-coordinates $(r, s)$ such that

$$
v_{2}=\frac{\partial}{\partial r}, \quad v_{1}=a(r, s) \frac{\partial}{\partial r}+b(r, s) \frac{\partial}{\partial s}, \quad \frac{\partial a}{\partial r}+\frac{\partial b}{\partial s}=0 .
$$

From (7.31),

$$
\frac{\partial a}{\partial r}=-1, \quad \frac{\partial b}{\partial r}=0
$$

and there exist a function $\alpha(s)$ and $b_{0} \in \mathbf{C}$ such that

$$
v_{1}=(-r+\alpha(s)) \frac{\partial}{\partial r}+\left(s+b_{0}\right) \frac{\partial}{\partial s} .
$$

Let us choose the $\Gamma_{s}$-coordinates

$$
u=r-\left(s+b_{0}\right) \int \alpha(s) d s, \quad v=s+b_{0}
$$

Then

$$
v_{2}=\frac{\partial}{\partial u}, \quad v_{1}=-u \frac{\partial}{\partial u}+v \frac{\partial}{\partial v} .
$$

Let

$$
v_{3}=e(u, v) \frac{\partial}{\partial u}+f(u, v) \frac{\partial}{\partial v}, \quad \frac{\partial e}{\partial u}+\frac{\partial f}{\partial v}=0 .
$$

From (7.33), we obtain

$$
\frac{\partial e}{\partial u}=2 u, \quad \frac{\partial f}{\partial u}=-2 v,
$$

and there exists a function $\varphi(v)$ and $f_{0} \in \mathbf{C}$ such that

From (7.32),

$$
v_{3}=\left(u^{2}+\varphi(v)\right) \frac{\partial}{\partial u}+\left(-2 u v+f_{0}\right) \frac{\partial}{\partial v} .
$$

$$
v \frac{\mathrm{~d} \varphi(v)}{\mathrm{d} v}+2 \varphi(v)=0
$$

and we obtain the existence of $\varphi_{0} \in \mathbf{C}$ such that

$$
v_{3}=\left(u^{2}-\frac{\varphi_{0}}{v^{2}}\right) \frac{\partial}{\partial u}+\left(-2 u v+f_{0}\right) \frac{\partial}{\partial v} .
$$

Finally, introduce the $\Gamma_{s}$-coordinates

$$
x=u+\frac{f_{0}}{2 v}, \quad y=v
$$

we have

$$
\begin{equation*}
v_{1}=-x \frac{\partial}{\partial x}+y \frac{\partial}{\partial y}, \quad v_{2}=\frac{\partial}{\partial x}, \quad v_{3}=\left(x^{2}-\frac{\alpha^{2}}{y^{2}}\right) \frac{\partial}{\partial x}-2 x y \frac{\partial}{\partial y} . \tag{7.4}
\end{equation*}
$$

Now, it is easy to check that (7.4) are the infinitesimal transformations of (III).
Let the vector fields $w_{1}, w_{2}, w_{3}$ on $\mathbf{C}^{2}$ generate the algebra (7.1). Then the vector fields $i w_{1}, i v_{2}, i w_{3}$ generate the algebra (7.2), and the vector fields $v_{1}=i w_{3}, v_{2}=$ $=w_{2}-i w_{1}, v_{3}=w_{2}+i w_{1}$ satisfy (7.3). Thus we obtain the existence of $\Gamma_{s}$-coordinates $(x, y)$ such that

$$
i w_{3}=-x \frac{\partial}{\partial x}+y \frac{\partial}{\partial y}, \quad w_{2}-i w_{1}=\frac{\partial}{\partial x}, w_{2}+i w_{1}=\left(x^{2}-\frac{\alpha^{2}}{y^{2}}\right) \frac{\partial}{\partial x}-2 x y \frac{\partial}{\partial y} .
$$

Our result is as follows: Let the vector fields $w_{1}, w_{2}, w_{3}$ satisfy (7.1), then there are (local) $\Gamma_{s}$-coordinates $(x, y)$ such that

$$
\begin{align*}
& w_{1}=\frac{1}{2} i\left(1-x^{2}+\frac{\alpha^{2}}{y^{2}}\right) \frac{\partial}{\partial x}+i x y \frac{\partial}{\partial y}, \\
& w_{2}=\frac{1}{2}\left(1+x^{2}-\frac{\alpha^{2}}{y^{2}}\right) \frac{\partial}{\partial x}-x y \frac{\partial}{\partial y},  \tag{7.5}\\
& w_{3}=i x \frac{\partial}{\partial x}-i y \frac{\partial}{\partial y} .
\end{align*}
$$

8. Consider the space $\mathbf{R}^{4}$ and its decomposition $\mathbf{R}^{4}=\mathbf{R}_{1}^{2} \oplus \mathbf{R}_{2}^{2}$. Denote by $H$ the group $\left\{\gamma \in G L\left(\mathbf{R}^{4}\right) ; \gamma\left(\mathbf{R}_{1}^{2}\right)=\mathbf{R}_{1}^{2}, \gamma\left(\mathbf{R}_{2}^{2}\right)=\mathbf{R}_{2}^{2}\right\}$, and let $\Gamma$ be the pseudogroup of local diffeomorphisms $\varphi: U \subset \mathbf{R}^{4} \rightarrow \mathbf{R}^{4}$ satisfying ( $\left.d \varphi\right)_{x} \in H$ for each $x \in U$. We wish to study hypersurfaces $M^{3} \subset \mathbf{R}^{4}$ with respect to $\Gamma$. Let $m \in M^{3}, T_{m}\left(M^{3}\right)$ the tangent space of $M^{3}$ at $m$; denote by $S_{i}^{2}(m) ; i=1,2$; the plane for which $m \in S_{i}^{2}(m)$ and $S_{i}^{2}(m) \cap R_{i}^{2}=\Phi$. In what follows, let us restrict ourselves to the study of hypersurfaces $M^{3} \subset \mathbf{R}^{4}$ satisfying the following conditions: (i) $M^{3}$ is analytic; (ii) $t_{i}(m)=T_{m}\left(M^{3}\right) \cap S_{i}^{2}(m)$ is one-dimensional for each $m \in M^{3}$ and
$=1,2$; (iii) $\tau_{m} \subset T_{m}\left(M^{3}\right)$ being the plane spanned by $t_{1}(m)$ and $t_{2}(m)$, the field $\tau_{m}$ is non-integrable. By means of the theory of systems of partial differential equations in involution (see, p.ex., K. Kuranishi, Lectures on involutive systems of partial differential equations; Publ. da Soc. Mat. de Sao Paulo, 1967), it is not difficult to prove

Theorem. Let $M^{3} \subset \mathbf{R}^{4}$ be a hypersurface and $\Phi: M^{3} \rightarrow \mathbf{R}^{4}$ an analytic mapping such that both $M^{3}$ and $\widetilde{M}^{3}=\Phi\left(M^{3}\right)$ are hypersurfaces satisfying the conditions mentioned above. Let $(d \Phi)_{m}\left(t_{i}(m)\right)=\tilde{t}_{i}(m)$ for each $m \in M^{3}$ and $i=1,2$; let $m_{0} \in M^{3}$ be a fixed point. Then there is a neighbourhood $U \subset M^{3}$ of $m_{0}$ and a diffeomorphism $\varphi \in \Gamma$ such that $\varphi$ is defined on $U$ and $\varphi / U=\Phi$.

To each hypersurface $M^{3} \subset \mathbf{R}^{4}$, we associate a $G$-structure $B_{G}\left(M^{3}\right)$ as follows. Let $\left(v_{1}, v_{2}, v_{3}\right)$ be a frame in $T_{m}\left(M^{3}\right)$. Then $\left(v_{1}, v_{2}, v_{3}\right) \in B_{G}\left(M^{3}\right)$ if and only if $v_{i}$ spans $t_{i}(m)$ for $i=1,2 . \quad\left(w_{1}, w_{2}, w_{3}\right) \in B_{G}\left(M^{3}\right)$ being another frame at $m \in M^{3}$, we have

$$
\begin{equation*}
w_{1}=\alpha v_{1}, \quad w_{2}=\beta v_{2}, \quad w_{3}=\gamma v_{1}+\delta v_{2}+\varphi v_{3} ; \quad \alpha \beta \gamma \neq 0 . \tag{8.1}
\end{equation*}
$$

In a neighbourhood $U$ of $m \in M^{3}$, let us choose an analytic section ( $v_{1}, v_{2}, v_{3}$ ) of $B_{G}\left(M^{3}\right)$; $\left(w_{1}, w_{2}, w_{3}\right)$ being another section of $B_{G}\left(M^{3}\right)$, we have (8.1) with $\alpha, \ldots, \varphi$ real-valued functions on $U$. The vector fields $v_{1}, v_{2},\left[v_{1}, v_{2}\right]$ being $\mathbf{R}$-linearly independent, we may write

$$
\begin{align*}
& {\left[v_{1},\left[v_{1}, v_{2}\right]\right]=a_{1} v_{1}+a_{2} v_{2}+a_{3}\left[v_{1}, v_{2}\right],}  \tag{8.2}\\
& {\left[v_{2},\left[v_{1}, v_{2}\right]\right]=b_{1} v_{1}+b_{2} v_{2}+b_{3}\left[v_{1}, v_{2}\right]}
\end{align*}
$$

and

$$
\begin{align*}
& {\left[w_{1},\left[w_{1}, w_{2}\right]\right]=A_{1} w_{1}+A_{2} w_{2}+A_{3}\left[w_{1}, w_{2}\right],}  \tag{8.3}\\
& {\left[w_{2},\left[w_{1}, w_{2}\right]\right]=B_{1} w_{1}+B_{2} w_{2}+B_{3}\left[w_{1}, w_{2}\right] .}
\end{align*}
$$

From the Jacobi identity

$$
\left[v_{1},\left[v_{2},\left[v_{1}, v_{2}\right]\right]\right]+\left[v_{2},\left[\left[v_{1}, v_{2}\right], v_{1}\right]\right]=0
$$

we get

$$
\begin{align*}
& v_{1} b_{1}-v_{2} a_{1}+a_{1} b_{3}-a_{3} b_{1}=0,  \tag{8.4}\\
& v_{1} b_{2}-v_{2} a_{2}+a_{2} b_{3}-a_{3} b_{2}=0, \\
& v_{1} b_{3}-v_{2} a_{3}+b_{2}+a_{1}=0
\end{align*}
$$

and analoguous equations for $A_{1}, \ldots, B_{3}$. Introduce the functions

$$
\begin{equation*}
p=\left(\alpha \beta^{2}\right)^{1 / 3}, \quad q=\left(\alpha^{2} \beta\right)^{1 / 3} \tag{8.5}
\end{equation*}
$$

over $U$ so that the equations $\left(8.1_{1,2}\right)$ become

$$
\begin{equation*}
w_{1}=p^{-1} q^{2} v_{1}, \quad w_{2}=p^{2} q^{-1} v_{2} . \tag{8.6}
\end{equation*}
$$

Then

$$
\begin{aligned}
& {\left[w_{1}, w_{2}\right]=\left[p^{-1} q^{2} v_{1}, p^{2} q^{-1} v_{2}\right]=} \\
& =\left(q \cdot v_{2} p-2 p . v_{2} q\right) v_{1}+\left(2 q \cdot v_{1} p-p . v_{1} q\right) v_{2}+p q\left[v_{1}, v_{2}\right], \\
& {\left[w_{1},\left[w_{1}, w_{2}\right]\right]=(.) v_{1}+(.) v_{2}+\left(q^{3} a_{3}+3 p^{-1} q^{3} \cdot v_{1} p\right)\left[v_{1}, v_{2}\right]=} \\
& =(.) v_{1}+(.) v_{2}+p q A_{3}\left[v_{1}, v_{2}\right], \\
& {\left[w_{2},\left[w_{1}, w_{2}\right]\right]=(.) v_{1}+(.) v_{2}+\left(p^{3} b_{3}+3 p^{3} q^{-1} . v_{2} q\right)\left[v_{1}, v_{2}\right]=} \\
& =(.) v_{1}+(.) v_{2}+p q B_{3}\left[v_{1}, v_{2}\right],
\end{aligned}
$$

and we have

$$
\begin{equation*}
p^{-1} q^{2}\left(a_{3}+3 p^{-1} \cdot v_{1} p\right)=A_{3}, \quad p^{2} q^{-1}\left(b_{3}+3 q^{-1} \cdot v_{2} q\right)=B_{3} \tag{8.8}
\end{equation*}
$$

The section ( $v_{1}, v_{2}, v_{3}$ ) of $B_{G}\left(M^{3}\right)$ being given, there exists (possibly in a small neighbourhood $U_{1} \subset U$ of $m \in M^{3}$ ) a section ( $w_{1}, w_{2}, w_{3}$ ) of $B_{G}\left(M^{3}\right)$ satisfying (8.3) with $A_{3}=B_{3}=0$; indeed, it is sufficient to take the section (8.6) where $p, q$ are any solutions of the system

$$
\begin{equation*}
v_{1} p=-\frac{1}{3} p a_{3}, \quad v_{2} q=-\frac{1}{3} q b_{3} . \tag{8.9}
\end{equation*}
$$

In what follows, let us restrict ourselves to the sections $\left(v_{1}, v_{2}, v_{3}\right),\left(w_{1}, w_{2}, v_{3}\right)$ of $B_{G}\left(M^{3}\right)$ satisfying

$$
\begin{equation*}
a_{3}=b_{3}=0 \quad \text { or } \quad A_{3}=B_{3}=0 \quad \text { resp. } ; \tag{8.10}
\end{equation*}
$$

we have $(8.6)+\left(8.1_{3}\right)$ with

$$
\begin{equation*}
v_{1} p=0, \quad v_{2} q=0 \tag{8.11}
\end{equation*}
$$

Now,
$\left[w_{1},\left[w_{1}, w_{2}\right]\right]=\left(2 p^{-1} q^{3} \cdot v_{1} v_{2} p+2 q^{2} \cdot v_{2} v_{1} q-2 p^{-1} q^{2} \cdot v_{2} p \cdot v_{1} q+\right.$

$$
\begin{equation*}
\left.+q^{3} a_{1}\right) v_{1}+\left(-q^{2} \cdot v_{1} v_{1} q+q^{3} a_{2}\right) v_{2}=p^{-1} q^{2} A_{1} v_{1}+p^{2} q^{-1} A_{2} v_{2} \tag{8.13}
\end{equation*}
$$

$\left[w_{2},\left[w_{1}, w_{2}\right]\right]=\left(p^{2} . v_{2} v_{2} p+p^{3} b_{1}\right) v_{1}+\left(-2 p^{2} . v_{1} v_{2} p-2 p^{3} q^{-1} \cdot v_{2} v_{1} q+\right.$

$$
\left.+2 p^{2} q^{-1} \cdot v_{2} p \cdot v_{1} q+p^{3} b_{2}\right) v_{2}=p^{-1} q^{2} B_{1} v_{1}+p^{2} q^{-1} B_{2} v_{2}
$$

i.e.,

$$
\begin{gather*}
-q^{3} \cdot v_{1} v_{1} q+q^{4} a_{2}=p^{2} A_{2}  \tag{8.14}\\
p^{3} \cdot v_{2} v_{2} p+p^{4} b_{1}=q^{2} B_{1} \\
2 q \cdot v_{1} v_{2} p+2 p \cdot v_{2} v_{1} q-2 v_{2} p \cdot v_{1} q+p q a_{1}=A_{1} \\
-2 q \cdot v_{1} v_{2} p-2 p \cdot v_{2} v_{1} q+2 v_{2} p \cdot v_{1} q+p q b_{2}=B_{2} . \tag{8.15}
\end{gather*}
$$

The equations (8.4) reduce to

$$
\begin{equation*}
v_{1} b_{1}-v_{2} a_{1}=0, \quad v_{1} b_{2}-v_{2} a_{2}=0, \quad b_{2}+a_{1}=0 \tag{8.16}
\end{equation*}
$$

and analoguous equations for $A_{1}, \ldots, B_{2}$; thus, (8.15) is a consequence of (8.143) and (8.163).

Let us consider the system (8.11) $+(8.14)$. From (8.11) and (8.141.2), we get

$$
\begin{array}{ll}
v_{1} v_{1} p=0, & v_{1} v_{2} q=0  \tag{8.17}\\
v_{2} v_{1} p=0, & v_{2} v_{2} q=0 \\
v_{2} v_{2} p=p^{-3} q^{2} B_{1}-p b_{1}, & v_{1} v_{1} q=q a_{2}-p^{2} q^{-3} A_{2}
\end{array}
$$

and

$$
\begin{align*}
& v_{1} v_{1} v_{1} p=v_{2} v_{1} v_{1} p=v_{1} v_{2} v_{1} p=v_{2} v_{2} v_{1} p=0,  \tag{8.18}\\
& v_{1} v_{2} v_{2} p=2 p^{-3} q B_{1} \cdot v_{1} q+p^{-3} q^{2} \cdot v_{1} B_{1}-p \cdot v_{1} b_{1}, \\
& v_{2} v_{2} v_{2} p=-3 p^{-4} q^{2} B_{1} \cdot v_{2} p-b_{1} \cdot v_{2} p+p^{-3} q^{2} \cdot v_{2} B_{1}-p . v_{2} b_{1}, \\
& v_{1} v_{1} v_{2} q=v_{2} v_{1} v_{2} q=v_{1} v_{2} v_{2} q=v_{2} v_{2} v_{2} q=0, \\
& v_{1} v_{1} v_{1} q=a_{2} \cdot v_{1} q+3 p^{2} q^{-4} A_{2} \cdot v_{1} q+q \cdot v_{1} a_{2}-p^{2} q^{-3} \cdot v_{1} A_{2}, \\
& v_{2} v_{1} v_{1} q=-2 p q^{-3} A_{2} \cdot v_{2} p+q \cdot v_{2} a_{2}-p^{2} q^{-3} \cdot v_{2} A_{2} .
\end{align*}
$$

The equations (8.2) may be rewritten as

$$
\begin{align*}
& v_{1} v_{1} v_{2}-2 v_{1} v_{2} v_{1}+v_{2} v_{1} v_{1}-a_{1} v_{1}-a_{2} v_{2}=0  \tag{8.19}\\
& 2 v_{2} v_{1} v_{2}-v_{2} v_{2} v_{1}-v_{1} v_{2} v_{2}-b_{1} v_{1}+a_{1} v_{2}=0
\end{align*}
$$

Applying them to the functions $p, q$, we get

$$
\begin{align*}
& v_{1} v_{1} v_{2} p=a_{2} \cdot v_{2} p  \tag{8.20}\\
& v_{2} v_{1} v_{2} p p^{-3} q B_{1} \cdot v_{1} q-\frac{1}{2} a_{1} \cdot v_{2} p+\frac{1}{2} p^{-3} q^{2} \cdot v_{1} B_{1}-\frac{1}{2} p \cdot v_{1} b_{1} \\
& v_{2} v_{2} v_{1} q=-b_{1} \cdot v_{1} q, \\
& v_{1} v_{2} v_{1} q=-p q^{-3} A_{2} \cdot v_{2} p-\frac{1}{2} a_{1} \cdot v_{1} q-\frac{1}{2} p^{2} q^{-3} \cdot v_{2} A_{2}+\frac{1}{2} q \cdot v_{2} a_{2}
\end{align*}
$$

Applying $v_{1}$ and $v_{2}$ to (8.143), we get

$$
2 q \cdot v_{1} v_{1} v_{2} p+2 p \cdot v_{1} v_{2} v_{1} q-2 v_{2} p \cdot v_{1} v_{1} q+p a_{1} \cdot v_{1} q+p q \cdot v_{1} a_{1}=v_{1} A_{1}
$$

$$
2 q \cdot v_{2} v_{1} v_{2} p+2 p \cdot v_{2} v_{2} v_{1} q-2 v_{1} q \cdot v_{2} v_{2} p+q a_{1} \cdot v_{2} p+p q \cdot v_{2} a_{1}=v_{2} A_{1}
$$ i.e.,

$$
\begin{aligned}
& q^{3}\left(v_{1} a_{1}+v_{2} a_{2}\right)=p^{-1} q^{2} \cdot v_{1} A_{1}+p^{2} q^{-1} \cdot v_{2} A_{2}=w_{1} A_{1}+w_{2} A_{2}, \\
& p^{3}\left(v_{2} a_{1}-v_{1} b_{1}\right)=p^{2} q^{-1} \cdot v_{2} A_{1}-p^{-1} q^{2} \cdot v_{1} B_{1}=w_{2} A_{1}-w_{1} B_{1}
\end{aligned}
$$

by means of (8.20). These equations being satisfied because of (8.16), we see that all the differential consequences of (8.143) are consequences of the system (8.11) + $+\left(8.14_{1,2}\right)$.

From (8.2) $+(8.10)$, we get

$$
\begin{align*}
& {\left[v_{1},\left[v_{1},\left[v_{1}, v_{2}\right]\right]\right]=v_{1} a_{1} \cdot v_{1}+v_{1} a_{2} \cdot v_{2}+a_{2}\left[v_{1}, v_{2}\right],}  \tag{8.21}\\
& {\left[v_{2},\left[v_{2},\left[v_{1}, v_{2}\right]\right]\right]=v_{2} b_{1} \cdot v_{1}-v_{2} a_{1} \cdot v_{2}-b_{1}\left[v_{1}, v_{2}\right],}
\end{align*}
$$

i.e.,

$$
\begin{align*}
L_{1} \equiv & v_{1} v_{1} v_{1} v_{2}-3 v_{1} v_{1} v_{2} v_{1}+3 v_{1} v_{2} v_{1} v_{1}-v_{2} v_{1} v_{1} v_{1}-v_{1} a_{1} \cdot v_{1}-v_{1} a_{2} \cdot v_{2}- \\
& -a_{2} \cdot v_{1} v_{2}+a_{2} \cdot v_{2} v_{1}=0 \\
L_{2} \equiv & v_{1} v_{2} v_{2} v_{2}-3 v_{2} v_{1} v_{2} v_{2}+3 v_{2} v_{2} v_{1} v_{2}-v_{2} v_{2} v_{2} v_{1}-v_{2} b_{1} \cdot v_{2} a_{1} \cdot v_{2}+ \\
& +b_{1} \cdot v_{1} v_{2}-b_{1} \cdot v_{2} v_{1}=0 . \tag{8.22}
\end{align*}
$$

Now,

From $L_{1} q=0, L_{2} p=0$, we obtain

$$
\begin{gather*}
3 p q^{-4} A_{2} \cdot v_{2} p \cdot v_{1} q-3 p q^{-3} A_{2} \cdot v_{1} v_{2} p-3 p^{2} q^{-4} A_{2} \cdot v_{2} v_{1} q+  \tag{8.23}\\
+\frac{3}{2} p^{2} q^{-4} \cdot v_{2} A_{2} \cdot v_{1} q-p q^{-3} \cdot v_{1} A_{2} \cdot v_{2} p+\frac{3}{2} q \cdot v_{1} v_{2} a_{2}+\frac{3}{2} q a_{1} a_{2}- \\
-\frac{3}{2} p^{2} q^{-3} a_{1} A_{2}-\frac{3}{2} p^{2} q^{-3} \cdot v_{1} v_{2} A_{2}-q \cdot v_{2} v_{1} a_{2}+p^{2} q^{-3} \cdot v_{2} v_{1} A_{2}=0, \\
3 p^{-4} q B_{1} \cdot v_{2} p \cdot v_{1} q-3 p^{-4} q^{2} B_{1} \cdot v_{1} v_{2} p-3 p^{-3} q B_{1} \cdot v_{2} v_{1} q+ \\
\quad+\frac{3}{2} p^{-4} q^{2} \cdot v_{1} B_{1} \cdot v_{2} p-p^{-3} q \cdot v_{2} B_{1} \cdot v_{1} q+p^{-3} q^{2} \cdot v_{1} v_{2} B_{1}- \\
-\frac{3}{2} p^{-3} q^{2} \cdot v_{2} v_{1} B_{1}-p \cdot v_{1} v_{2} b_{1}+\frac{3}{2} p \cdot v_{2} v_{1} b_{1}-\frac{3}{2} p^{-3} q^{2} a_{1} B_{1}+\frac{3}{2} p a_{1} b_{1}=0 .
\end{gather*}
$$

$$
\begin{aligned}
& v_{1} v_{1} v_{2} v_{1} q=3 p q^{-4} A_{2} \cdot v_{2} p \cdot v_{1} q-p q^{-3} \cdot v_{1} A_{1} \cdot v_{2} p-p q^{-3} A_{2} \cdot v_{1} v_{2} p-\frac{1}{2} v_{1} a_{1} \cdot v_{1} q- \\
& -\frac{1}{2} q a_{1} a_{2}+\frac{1}{2} p^{2} q^{-3} a_{1} A_{2}+\frac{3}{2} p^{2} q^{-4} \cdot v_{2} A_{2} \cdot v_{1} q-\frac{1}{2} p^{2} q^{-3} \cdot v_{1} v_{2} A_{2}+ \\
& +\frac{1}{2} v_{2} a_{2} \cdot v_{1} q+\frac{1}{2} q . v_{1} v_{2} a_{2}, \\
& v_{1} v_{2} v_{1} v_{1} q=6 p q^{-4} A_{2} \cdot v_{2} p . v_{1} q-2 p q^{-3} \cdot v_{1} A_{2} . v_{2} p-2 p q^{-3} A_{2} . v_{1} v_{2} p+ \\
& +v_{2} a_{2} \cdot v_{1} q+q . v_{1} v_{2} a_{2}+3 p^{2} q^{-4} . v_{2} A_{2} . v_{1} q-p^{2} q^{-3} . v_{1} v_{2} A_{2} \text {, } \\
& v_{2} v_{1} v_{1} v_{1} q=v_{2} a_{2} . v_{1} q+a_{2} . v_{2} v_{1} q+6 p q^{-4} A_{2} . v_{2} p . v_{1} q+3 p^{2} q^{-4} \cdot v_{2} A_{2} . v_{1} q+ \\
& +3 p^{2} q^{-4} A_{2} . v_{2} v_{1} q+q . v_{2} v_{1} a_{2}-2 p q^{-3} . v_{1} A_{2} . v_{2} p-p^{2} q^{-3} . v_{2} v_{1} A_{2} \text {, } \\
& v_{1} v_{2} v_{2} v_{2} p=-6 p^{-4} q B_{1} . v_{2} p . v_{1} q-3 p^{-4} q^{2} . v_{1} B_{1} . v_{2} p-3 p^{-4} q^{2} B_{1} . v_{1} v_{2} p- \\
& -v_{1} b_{1} . v_{2} p-b_{1} \cdot v_{1} v_{2} p+2 p^{-3} q . v_{2} B_{1} . v_{1} q+p^{-3} q^{2} \cdot v_{1} v_{2} B_{1}- \\
& -p . v_{1} v_{2} b_{1}, \\
& v_{2} v_{1} v_{2} v_{2} p=-6 p^{-4} q B_{1} \cdot v_{2} p \cdot v_{1} q+2 p^{-3} q \cdot v_{2} B_{1} \cdot v_{1} q+2 p^{-3} q B_{1} \cdot v_{2} v_{1} q- \\
& -3 p^{-4} q^{2} . v_{1} B_{1} . v_{2} p+p^{-3} q^{2} . v_{2} v_{1} B_{1}-v_{1} b_{1} \cdot v_{2} p-p . v_{2} v_{1} b_{1}, \\
& v_{2} v_{2} v_{1} v_{2} p=-3 p^{-4} q B_{1} \cdot v_{2} p . v_{1} q+p^{-3} q \cdot v_{2} B_{1} . v_{1} q+ \\
& +p^{-3} q B_{1} . v_{2} v_{1} q-\frac{1}{2} v_{2} a_{1} . v_{2} p-\frac{1}{2} p^{-3} q^{2} a_{1} B_{1}+ \\
& +\frac{1}{2} p a_{1} b_{1}-\frac{3}{2} p^{-4} q^{2} \cdot v_{1} B_{1} \cdot v_{2} p+\frac{1}{2} p^{-3} q^{2} . v_{2} v_{1} B_{1}- \\
& -\frac{1}{2} v_{1} b_{1} . v_{2} p-\frac{1}{2} p . v_{2} v_{1} b_{1} .
\end{aligned}
$$

Multiplying (8.143) by $\frac{3}{2} p q^{-4} A_{2}$ or $\frac{3}{2} p^{-4} q B_{1}$ resp. and adding it to (8.23 $)$ or (8.232) resp., we get

$$
\begin{gather*}
\frac{3}{2} p^{2} q^{-4} \cdot v_{2} A_{2} \cdot v_{1} q-p q^{-3} \cdot v_{1} A_{2} \cdot v_{2} p+\frac{3}{2} q \cdot v_{1} v_{2} a_{2}+\frac{3}{2} q a_{1} a_{2}-  \tag{8.24}\\
\quad-\frac{3}{2} p^{2} q^{-3} \cdot v_{1} v_{2} A_{2}-q \cdot v_{2} v_{1} a_{2}+p^{2} q^{-3} \cdot v_{2} v_{1} A_{2}-\frac{3}{2} p q^{-4} A_{1} A_{2}=0, \\
\frac{3}{2} p^{-4} q^{2} \cdot v_{1} B_{1} \cdot v_{2} p-p^{-3} q \cdot v_{2} B_{1} \cdot v_{1} q+p^{-3} q^{2} \cdot v_{1} v_{2} B_{1}- \\
-\frac{3}{2} p^{-3} q^{2} \cdot v_{2} v_{1} B_{1}-p \cdot v_{1} v_{2} b_{1}+\frac{3}{2} p \cdot v_{2} v_{1} b_{1}+\frac{3}{2} p a_{1} b_{1}-\frac{3}{2} p^{-4} q A_{1} B_{1}=0 .
\end{gather*}
$$

From (8.6),
$v_{2} v_{1}=q^{-2} \cdot v_{2} p . w_{1}+p^{-1} q^{-1} \cdot w_{2} w_{1}, v_{1} v_{2}=p^{-2} \cdot v_{1} q \cdot w_{2}+p^{-1} q^{-1} \cdot w_{1} w_{2}$.
Finally, we get

$$
\begin{array}{r}
p^{-1} q^{5}\left(3 v_{1} v_{2} a_{2}-2 v_{2} v_{1} a_{2}+3 a_{1} a_{2}\right)=3 w_{1} w_{2} A_{2}-2 w_{2} w_{1} A_{2}+3 A_{1} A_{2},  \tag{8.26}\\
p^{5} q^{-1}\left(3 v_{2} v_{1} b_{1}-2 v_{1} v_{2} b_{1}+3 a_{1} b_{1}\right)=3 w_{2} w_{1} B_{1}-2 w_{1} w_{2} B_{1}+3 A_{1} B_{1}
\end{array}
$$

from (8.6), (8.25) and (8.24).
Let us write

$$
\begin{array}{ll}
j_{1}=3 v_{1} v_{2} a_{2}-2 v_{2} v_{1} a_{2}+3 a_{1} a_{2}, & j_{2}=3 v_{2} v_{1} b_{1}-2 v_{1} v_{2} b_{1}+3 a_{1} b_{1},  \tag{8.27}\\
\mathscr{f}_{1}=3 w_{1} w_{2} A_{2}-2 w_{2} w_{1} A_{2}+3 A_{1} A_{2}, \mathcal{f}_{2}=3 w_{2} w_{1} B_{1}-2 w_{1} w_{2} B_{1}+3 A_{1} B_{1} .
\end{array}
$$

Then

$$
\begin{equation*}
\alpha^{3} \beta j_{1}=\mathfrak{f}_{1}, \quad \alpha \beta^{3} j_{2}=\mathscr{f}_{2} \tag{8.28}
\end{equation*}
$$

Suppose

$$
\begin{equation*}
j_{1} j_{2} \neq 0 \tag{8.29}
\end{equation*}
$$

and write

$$
\begin{equation*}
k_{1}=\left|j_{1}-3 j_{2}\right|^{1 / 8}, \quad k_{2}=\left|j_{1} j_{2}-3\right|^{1 / 8} \tag{8.30}
\end{equation*}
$$

Then

$$
\begin{equation*}
k_{1}=|\alpha| \cdot K_{1}, \quad k_{2}=|\beta| \cdot K_{2} \tag{8.31}
\end{equation*}
$$

and

$$
\begin{equation*}
K_{1} w_{1}=\operatorname{sgn} \alpha \cdot k_{1} v_{1}, \quad K_{2} w_{2}=\operatorname{sgn} \beta \cdot k_{2} v_{2} . \tag{8.32}
\end{equation*}
$$

Theorem. On $M^{3}$, be given a $G$-structure $B_{G}\left(M^{3}\right)$ of the considered type. In a neighbourhood of $m_{0} \in M^{3}$, let us choose its section ( $v_{1}, v_{2}, v_{3}$ ) in such a way that (8.2) and (8.10) are satisfied. Suppose that we have (8.29) for the functions $j_{1}, j_{2}$ defined by (8.27). Consider the vector fields

$$
\begin{equation*}
V=k_{1} v_{1} . \quad V_{2}=k_{2} v_{2} \tag{8.33}
\end{equation*}
$$

$k_{1}$ and $k_{2}$ being defined by (8.30). These vector fields are invariant up to the sign, i.e., choosing another section ( $w_{1}, w_{2}$, $w_{3}$ ) satisfying (8.3) and (8.10), we have $W_{1} \equiv K_{1} w_{1}=$ $= \pm V_{1}, W_{2} \equiv K_{2} w_{2}= \pm V_{2}$.
9. Consider the space $\mathbf{C}^{2}$, i.e., the space $\mathbf{R}^{4}$ endowed with a fixed automorphism $I: \mathbf{R}^{4} \rightarrow \mathbf{R}^{4}$ satisfying $I^{2}=-i d$. Let $H^{\prime} \subset G L\left(\mathbf{R}^{4}\right)$ be the subgroup of elements $\gamma \in G L\left(\mathbf{R}^{4}\right)$ satisfying $\gamma I=I \gamma$. The local diffeomorphism $\varphi: U \subset \mathbf{R}^{4} \rightarrow \mathbf{R}^{4}$ is called holomorphic if $(d \varphi)_{x} \in H^{\prime}$ for each $x \in U$. Our task is to study hypersurfaces $M^{3} \subset \mathbf{R}^{4}$ with respect to the pseudogroup $\Gamma^{\prime}$ of all local holomorphic diffeomorphisms.

Let $m \in M^{3}$. Write $\tau_{m}=T_{m}\left(M^{3}\right) \cap I T_{m}\left(M^{3}\right) ; \tau_{m}$ is always a plane. Let us restrict ourselves to hypersurfaces for which the field of planes $\tau_{m}$ is non-integrable. To $M^{3}$, we associate a $G^{\prime}$-structure $B_{G^{\prime}}^{\prime}\left(M^{3}\right)$ as follows. The frame ( $u_{1}, u_{2}, u_{3}$ ) of $T_{m}\left(M^{3}\right)$ belongs to $B_{G^{\prime}}^{\prime}\left(M^{3}\right)$ if and only if $u_{1} \in \tau_{m}, u_{2}=I u_{1} . \quad\left(\tilde{u}_{1}, \tilde{u}_{2}, \tilde{u}_{3}\right)$ being another frame of $B_{G^{\prime}}^{\prime}\left(M^{3}\right)$ over $m$, we have

$$
\begin{align*}
& \tilde{u}_{1}=\varrho u_{1}-\sigma u_{2},  \tag{9.1}\\
& \tilde{u}_{2}=\sigma u_{1}+\varrho u_{2}, \\
& \tilde{u}_{3}=\varkappa_{1} u_{1}+\varkappa_{2} u_{2}+\varkappa u_{3} ; \quad\left(\varrho^{2}+\sigma^{2}\right) \varkappa \neq 0 .
\end{align*}
$$

In a neighbourhood of $m \in M^{3}$, let us choose a section $\left(u_{1}, u_{2}, u_{3}\right)$ of $B_{G^{\prime}}^{\prime}\left(M^{3}\right)$. We may write

$$
\begin{align*}
& {\left[u_{1},\left[u_{1}, u_{2}\right]\right]=c_{1}, u_{1}+c_{2} u_{2}+c_{3}\left[u_{1}, u_{2}\right],}  \tag{9.2}\\
& {\left[u_{2},\left[u_{1}, u_{2}\right]\right]=d_{1} u_{1}+d_{2} u_{2}+d_{3}\left[u_{1}, u_{2}\right] .}
\end{align*}
$$

Consider the complexification $T^{C}\left(M^{3}\right)=T\left(M^{3}\right) \oplus i T\left(M^{3}\right)$ of the tangent bundle $T\left(M^{3}\right)$ and its vector fields

$$
\begin{equation*}
v_{1}=u_{1}+i u_{2}, v_{2}=u_{1}-i u_{2} \text { or } w_{1}=\tilde{u}_{1}+i \tilde{u}_{2}, w_{2}=\tilde{u}_{1}-i \tilde{u}_{2} \text { resp. } \tag{9.3}
\end{equation*}
$$

Then

$$
\begin{equation*}
w_{1}=\alpha v_{1}, w_{2}=\beta v_{2}, \quad \text { where } \quad \alpha=\varrho+i \sigma, \beta=\varrho-i \sigma \tag{9.4}
\end{equation*}
$$

Further,

$$
\begin{align*}
{\left[v_{1},\left[v_{1}, v_{2}\right]\right]=} & \left\{d_{1}-c_{2}-i\left(d_{2}+c_{1}\right)\right\} v_{1}+\left\{d_{1}+c_{2}+i\left(d_{2}-c_{1}\right)\right\} v_{2}+(9  \tag{9.5}\\
& +\left(c_{3}+i d_{3}\right)\left[v_{1}, v_{2}\right], \\
{\left[v_{2},\left[v_{1}, v_{2}\right]\right]=} & \left\{-d_{1}-c_{2}+i\left(d_{2}-c_{1}\right)\right\} v_{1}+\left\{-d_{1}+c_{2}-i\left(d_{2}+c_{1}\right)\right\} v_{2}+ \\
& +\left(c_{3}-i d_{3}\right)\left[v_{1}, v_{2}\right] .
\end{align*}
$$

To obtain invariants of $M^{3} \subset \mathbf{C}^{2}$, we proceed formally in the same way as we have done in the preceding section.

