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# Two Classes of Numerical Methods for Stiff Problems 

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#### Abstract

Two classes of numerical methods for stiff problems are shown. Formulae contained in the second clase require to solve a system of only linear algebr. eqs. to obtain the solution of a nonlinear system of differential eqs. at each step. One such formula is tested on a very stiff problem and the comparison with other often used methods is given.


Let us consider the differential equation of the form:

$$
\begin{equation*}
y^{\prime}=f(x, y) \quad y\left(x_{0}\right)=y_{0} \quad x \in\left\langle x_{0}, x_{d}\right\rangle . \tag{1}
\end{equation*}
$$

Under obvious conditions on $f$ relations (1) are equivalent to:

$$
\begin{equation*}
y^{\prime \prime}=f^{\prime}(x, y) \quad y\left(x_{0}\right)=y_{0} \quad y^{\prime}\left(x_{k}\right)=f\left(x_{k}, y\left(x_{k}\right)\right) \quad x, x_{k} \in\left\langle x_{0}, x_{d}\right\rangle \tag{2}
\end{equation*}
$$

Defining the mesh $x_{i}=x_{0}+i h$ on the interval $\left\langle x_{0}, x_{k}\right\rangle \subset\left\langle x_{0}, x_{d}\right\rangle$ and approximating the relations (2) by finite differences at mesh-points, we can derive different formulae having the following general form:

$$
\begin{align*}
& I\left[\begin{array}{l}
y_{n+1} \\
\vdots \\
y_{n+k}
\end{array}\right]=\left[\begin{array}{l}
y_{n} \\
\vdots \\
y_{n}
\end{array}\right]+h\left[\begin{array}{l}
d_{1} \\
: \\
d_{k}
\end{array}\right] f\left(x_{n}, y_{n}\right)+h^{2}\left[\begin{array}{l}
e_{1} \\
\vdots \\
e_{k}
\end{array}\right] f^{\prime}\left(x_{n}, y_{n}\right)+ \\
& \quad+h B\left[\begin{array}{c}
f\left(x_{n+1}, y_{n+1}\right) \\
\vdots \\
f\left(x_{n+k}, y_{n+k}\right)
\end{array}\right]+h^{2} C\left[\begin{array}{c}
f^{\prime}\left(x_{n+1}, y_{n+1}\right) \\
\vdots \\
f^{\prime}\left(x_{n+k}, y_{n+k}\right)
\end{array}\right] \tag{3}
\end{align*}
$$

where: $I$ - unit $k \times k$ matrix, $B, C-k \times k$ matrices, $d_{i}, e_{i}$ - real numbers, $h$ - mesh size.

We assume $y_{n}$ to be a known starting value. The unknown values $y_{n+1} \ldots, y_{n+k}$ are to be calculated from the formula (3) and $y_{n+k}$ is used as a new starting value only.

We have shown for a sufficiently smooth right-hand side $f$ the necessary and sufficient conditions for method (3) to be of order $p$. Further, we have proved that every formula (3) having order $p \geqslant 1$ is convergent and the rate of convergence is $O\left(h^{p}\right)$.

The class (3) contains as a subset the class of selfstarting overimplicit methods (SOM). It has been shown in [1] that there exist $A$-stable methods of arbitrarily high order in the class SOM. We have derived the $A$-stable formulae up to the order 6 which contain the second derivatives of the solution and therefore do not belong to the class SOM. We believe it is possible to show in a way similar to that referred in [1] that the formulae of class (3) (containing the second derivatives) can also yield $A$-stable methods of any arbitrary order. But we have not proved it yet.

By applying the formula (3) to a nonlinear system we have to solve a set of nonlinear eqs. at each step. We suggest to use a certain iterative procedure (resembling the Newton method) requiring an evaluation of the Jacobian (not Hessian) matrix of the right-hand side of original differential system only. This procedure is described in [2].

The main goal of this communication is to devise a way of avoiding any kind of iteration.

We shall illustrate our procedure on the simplest type of the class (3). The class (3) takes for $k=1$ the form:
$y_{n+1}=y_{n}+h d f\left(x_{n}, y_{n}\right)+h^{2} e f^{\prime}\left(x_{n}, y_{n}\right)+h b f\left(x_{n+1}, y_{n+1}\right)+h^{2} c f^{\prime}\left(x_{n+1}, y_{n+1}\right)$.
Let us now consider a system of diff. eqs. Then $f$ and $f^{\prime}$ are vector functions and it holds: $f_{i}^{\prime}=\frac{\partial f_{i}}{\partial x}+\mathscr{F}_{i} f_{i}$, where $\mathcal{F}$ is the Jacobian matrix of $f$ and the subscript $i$ denotes that all values are taken at the point $x_{i}, y_{i}$. We replace in the formula (4) $f\left(x_{n+1}, y_{n+1}\right)$ by $f_{n}+h \frac{\partial f_{n}}{\partial x}+\mathscr{F}_{n} \Delta_{n}$ and $f^{\prime}\left(x_{n+1}, y_{n+1}\right)$ by $\frac{\partial f_{n}}{\partial x}+\mathscr{f}_{n} f_{n}+$ $+\mathscr{f}_{n} \frac{\partial f_{n}}{\partial x}+\mathscr{F}_{n}^{2} \Delta_{n}$ where $\Delta_{n}=y_{n+1}-y_{n}$. Requiring the formula (4) to be of the order $p \geqslant 2$ we finally obtain:
$\left[I-h b \mathcal{F}_{n}-h^{2} c \mathcal{F}_{n}^{2}\right] \Delta_{n}=h f_{n}+h^{2}\left[(0.5-b) \mathscr{f}_{n} f_{n}+0.5 \frac{\partial f_{n}}{\partial x}+h c \mathcal{F}_{n} \frac{\partial f_{n}}{\partial x}\right]$
This formula beeing applied to $y^{\prime}=a y$ yields the same expression as the original formula (4). Therefore (5) is $A$-stable if and only if (4) is $A$-stable. We have proved that for a sufficiently smooth right-hand side $f$ the formula (5) is convergent if (4) is convergent. If the order of (4) is $p \geqslant 2$, then the rate of convergence of (5) is $O\left(h^{2}\right)$. Referring to (5) a system of only linear algebraic Eqs. is to be solved at each step. For $b=1$ and $c=-0.5$ the formula (5) is A-stable and of the second order.

We have tested this formula on a very stiff system arising in the reactor kinetics (taken from [3]).

$$
\begin{array}{ll}
y_{1}^{\prime}=-0.04 y_{1}+10^{4} y_{2} y_{3} & y_{2}^{\prime}=0.04 y_{1}-10^{4} y_{2} y_{3}-3 \cdot 10^{7} y_{2}^{2} \\
y_{3}^{\prime}=3.10^{7} y_{2}^{2} & y_{1}(0)=1 y_{2}(0)=y_{3}(0)=0
\end{array}
$$

Results obtained for $x=4$ are in following table:

| $h$ | $y_{1}$ | $10^{4} y_{2}$ | $10 y_{3}$ |
| :---: | :---: | :---: | :---: |
| 0.4 | 0.98477 | 0.38157 | 0.35192 |
| 0.2 | 0.92398 | 0.24645 | 0.75995 |
| 0.05 | 0.90683 | 0.22557 | 0.93147 |
| 0.02 | 0.90561 | 0.22416 | 0.94361 |
| 0.01 | 0.90553 | 0.22406 | 0.94449 |
| $\$$ | 0.90552 | 0.22404 | 0.94458 |

$\S-$ reference solution

- Runge-Kutta method
of 4-th order, $h=0.001$

The comparison with several other methods has shown that our technique can very succesfully compete with all methods considered.

| Method | $\left\|y_{1}-y_{1}^{*}\right\|$ | $\left\|y_{2}-y_{2}^{*}\right\|$ | $\left\|y_{3}-y_{3}^{*}\right\|$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| BVT | $2.2 \mathrm{E}-4$ | $3.8 \mathrm{E}-8$ | $2.2 \mathrm{E}-4$ |  |  |
| Calahan | $1.4 \mathrm{E}+0$ | $4.0 \mathrm{E}-5$ | $1.4 \mathrm{E}+0$ |  |  |
| Allen | $9.8 \mathrm{E}-1$ | $4.7 \mathrm{E}-4$ | $2.4 \mathrm{E}+1$ |  |  |
| ISI3 ( -100 ) | $2.2 \mathrm{E}-{ }^{-3}$ | $4.0 \mathrm{E}-7$ | $2.2 \mathrm{E}-3$ |  |  |
| ISI3 ( $-\infty$ ) | $2.2 \mathrm{E}-3$ | $3.9 \mathrm{E}-7$ | $2.2 \mathrm{E}-3$ | $h=0.02$ | $x=0.4$ |
| LW1 | $1.6 \mathrm{E}-4$ | $2.4 \mathrm{E}-4$ | $3.2 \mathrm{E}-3$ | BVT | - method (5) |
| LW2 | $5.9 \mathrm{E}-4$ | $2.9 \mathrm{E}-3$ | $4.0 \mathrm{E}-2$ | LW1, LW2 | - derived in [2] <br> - derived in [4] |

The one-step nature of our method allows to implement an automatic step-size control. Results calculated by a step-size control procedure are compared with those obtained using the constant step-size in the following table (for $x=10$ ):

|  | $h$ | $\left\|y_{1}-y_{1}^{*}\right\|$ | $\left\|y_{2}-y_{2}^{*}\right\| .10^{4}$ | $\left\|y_{3}-y_{3}^{*}\right\| .10$ | evaluation <br> in $R H S^{\star}$ |
| :--- | :--- | :--- | :--- | :--- | :---: |
| constant | 0.4 | 0.027 | 0.023 | 0.27 | 25 |
| step | 0.05 | 0.001 | 0.001 | 0.01 | 200 |
|  | 0.02 | 0.000 | 0.000 | 0.00 | 500 |
| step-size <br> control <br> procedure |  | 0.000 | 0.000 | 0.00 | 38 |

[^0]
## References

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[^0]:    * RHS -- right-hand side

