## Mathematic Bohemica

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Mathematica Bohemica, Vol. 145 (2020), No. 3, 265-280

Persistent URL: http://dml.cz/dmlcz/148349

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# ON THE ORDER OF MAGNITUDE OF <br> WALSH-FOURIER TRANSFORM 

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Received June 27, 2018. Published online July 17, 2019.
Communicated by Jiří Spurný

Abstract. For a Lebesgue integrable complex-valued function $f$ defined on $\mathbb{R}^{+}:=[0, \infty)$ let $\hat{f}$ be its Walsh-Fourier transform. The Riemann-Lebesgue lemma says that $\hat{f}(y) \rightarrow 0$ as $y \rightarrow \infty$. But in general, there is no definite rate at which the Walsh-Fourier transform tends to zero. In fact, the Walsh-Fourier transform of an integrable function can tend to zero as slowly as we wish. Therefore, it is interesting to know for functions of which subclasses of $L^{1}\left(\mathbb{R}^{+}\right)$there is a definite rate at which the Walsh-Fourier transform tends to zero. We determine this rate for functions of bounded variation on $\mathbb{R}^{+}$. We also determine such rate of Walsh-Fourier transform for functions of bounded variation in the sense of Vitali defined on $\left(\mathbb{R}^{+}\right)^{N}, N \in \mathbb{N}$.

Keywords: function of bounded variation over $\mathbb{R}^{+}$; function of bounded variation over $\left(\mathbb{R}^{+}\right)^{2}$; function of bounded variation over $\left(\mathbb{R}^{+}\right)^{N}$; order of magnitude; Riemann-Lebesgue lemma; Walsh-Fourier transform

MSC 2010: 42C20, 26A12, 26A45, 26B30, 26D15

## 1. Introduction

We consider the Walsh orthonormal system $\left\{w_{m}(x): m \in \mathbb{N}_{0}\right\}, \mathbb{N}_{0}=\mathbb{N} \cup\{0\}$, defined on the unit interval $\mathbb{\square}:=[0,1$ ) in the Paley enumeration (see [19]). To go into some details, let

$$
r_{0}(x):=\left\{\begin{aligned}
1 & \text { if } x \in\left[0, \frac{1}{2}\right) \\
-1 & \text { if } x \in\left[\frac{1}{2}, 1\right)
\end{aligned}\right.
$$

[^0]and extend $r_{0}(x)$ for the half-real axis $\mathbb{R}^{+}:=[0, \infty)$ with period 1 . The Rademacher orthonormal system $\left\{r_{k}(x): k \in \mathbb{N}\right\}$ is defined by
$$
r_{k}(x):=r_{0}\left(2^{k} x\right), \quad k=1,2, \ldots, x \in \mathbb{0} .
$$

Now, the $m$ th Walsh function $w_{m}(x)$ in the Paley enumeration is defined as follows: If

$$
m=\sum_{k=0}^{\infty} m_{k} 2^{k}, \quad \text { where each } m_{k}=0 \text { or } 1
$$

is the binary decomposition of $m \in \mathbb{N}$, then let

$$
\begin{equation*}
w_{m}(x):=\prod_{k=0}^{\infty} r_{k}^{m_{k}}(x), \quad x \in 0 \tag{1.1}
\end{equation*}
$$

Clearly, $m_{k}=0$ except for a finite number of $k$ 's. Thus, the right-hand side of (1.1) is a finite product for each $m \in \mathbb{N}$. In particular, we have

$$
w_{0}(x) \equiv 1 \quad \text { and } \quad w_{2^{m}}=r_{m}(x), \quad m \in \mathbb{N}_{0}
$$

It is well known that $\left\{w_{m}(x): m \in \mathbb{N}_{0}\right\}$ is a complete orthonormal system on $\mathbb{\square}$.
Any $x \in \mathbb{\square}$ can be written in the form

$$
x=\sum_{k=0}^{\infty} x_{k} 2^{-k-1}, \quad \text { where each } x_{k}=0 \text { or } 1
$$

For each $x \in \mathbb{\backslash} \backslash$ there is only one expression of this form, where $Q$ is the collection of dyadic rationals in $\mathbb{\square}$. When $x \in Q$, there are two expressions of this form, one which terminates in 0's and other which terminates in 1's. Now the dyadic sum of $x, y \in \mathbb{Q}$ is defined by

$$
x \dot{+} y:=\sum_{k=0}^{\infty}\left|x_{k}-y_{k}\right| 2^{-k-1} .
$$

A remarkable property of the Walsh functions is that for each $m \in \mathbb{N}_{0}$ we have

$$
w_{m}(x \dot{+} y)=w_{m}(x) w_{m}(y), \quad x, y \in \mathbb{\mathbb { 1 }}, x \dot{+} y \notin Q
$$

Next, we consider the generalized Walsh functions $\psi_{x}, x \in \mathbb{R}^{+}$(see [20], Chapter 9 ), and recall the following properties:
(i) $\psi_{k}(x)=w_{k}(x)$ for $k \in \mathbb{N}_{0}, x \in \mathbb{\square}$;
(ii) $\psi_{y}(x \dot{+} t)=\psi_{y}(x) \psi_{y}(t)$ for $x, t \in \mathbb{R}^{+}$and $x \dot{+} t$ dyadic irrational;
(iii) $\psi_{y}(x)=\psi_{x}(y), \psi_{y}(x)=\psi_{[y]}(x) \psi_{[x]}(y)$ for $x, y \in \mathbb{R}^{+}$, where for $u \in \mathbb{R}^{+}$, [u] represents the greatest integer in $u$;
(iv) the functions $\psi_{j}, j \in \mathbb{N}_{0}$ form a complete orthonormal system in each of the intervals of the form $[k, k+1), k \in \mathbb{N}_{0}$;
(v) $\psi_{j}$ is a periodic extension of $w_{j}$ from $\mathbb{\square}$ to $\mathbb{R}^{+}$.

Now we recall (see e.g. [20], page 421) that the Walsh-Fourier transform of an $f \in L^{1}\left(\mathbb{R}^{+}\right)$is defined by

$$
\begin{equation*}
\hat{f}(y):=\int_{0}^{\infty} f(x) \psi_{y}(x) \mathrm{d} x, \quad y \in \mathbb{R}^{+} \tag{1.2}
\end{equation*}
$$

We also recall that the Riemann-Lebesgue lemma holds for Walsh-Fourier transform (see [20], page 422), that is, $\hat{f}(y) \rightarrow 0$ as $y \rightarrow \infty$. But in general, there is no definite rate at which the Walsh-Fourier transform tends to zero. In fact, the Walsh-Fourier transform of an integrable function can tend to zero as slowly as we wish (see e.g. [13], 32.47 (b)). Therefore, it is interesting to know for functions of which subclasses of $L^{1}\left(\mathbb{R}^{+}\right)$there is a definite rate at which the Walsh-Fourier transform tends to zero.

Looking to the periodic version, for the case of trigonometric Fourier series, that is, for functions on one-dimensional torus $\mathbb{T}:=[0,2 \pi)$, the study of order of magnitude of Fourier coefficients is done extensively (see e.g. [15], [21], see also [3], Section 2.3, page 30 and [22], Section 4, page 45). This study in periodic version for trigonometric Fourier series is done even for more general cases, that is, for the case of functions on two-dimensional torus, or more generally, on the $N$-dimensional torus $\mathbb{T}^{N}:=[0,2 \pi)^{N}$, $N \in \mathbb{N}$ (see e.g. [16], [5], [6], [8]).

Also looking to the periodic version, for the case of Walsh-Fourier series, that is, for functions defined on $\mathbb{\square}$, the study of order of magnitude of Walsh-Fourier coefficients is done (see e.g. [4], [11]). This study in periodic version for Walsh-Fourier series is done even for more general cases, that is, for the case of functions on two-dimensional torus $\rrbracket^{2}$, or more generally, on the $N$-dimensional torus $\rrbracket^{N}, N \in \mathbb{N}$ (see e.g. [7]).

Recently, in 2015 (see [9]), we have studied the order of magnitude of trigonometric Fourier transform for functions of bounded variation on $\mathbb{R}$ and for functions of bounded variation in the sense of Vitali on $\mathbb{R}^{N}$ and obtained results analogous to the periodic case. But it appears that such a study for the Walsh-Fourier transform has not yet been done. In this paper we carry out this study and determine the rate of decay of Walsh-Fourier transform for functions of bounded variation on $\mathbb{R}^{+}$. We also determine such rate of Walsh-Fourier transform for functions of bounded variation in the sense of Vitali defined on $\left(\mathbb{R}^{+}\right)^{N}, N \in \mathbb{N}$.

## 2. One-dimensional case

We recall that a function $f: \mathbb{R}^{+} \rightarrow \mathbb{C}$ is said to be of bounded variation over $\mathbb{R}^{+}$, in symbol $f \in \operatorname{BV}\left(\mathbb{R}^{+}\right)$, if

$$
\begin{equation*}
\sup _{\mathcal{S}} \sum_{k=1}^{n}\left|f\left(x_{k}\right)-f\left(x_{k-1}\right)\right|<\infty \tag{2.1}
\end{equation*}
$$

where the supremum is extended over all finite sequences

$$
\mathcal{S}: 0 \leqslant x_{0}<x_{1}<x_{2}<\ldots<x_{n}<\infty \quad \text { and } \quad n=1,2, \ldots
$$

The supremum in (2.1), denoted by $V(f)$, is called the total variation of $f$ over $\mathbb{R}^{+}$.
It is clear that the above definition of bounded variation over $\mathbb{R}^{+}$can be reformulated equivalently as follows. A function $f$ is of bounded variation over $\mathbb{R}^{+}$if and only if $f$ is of bounded variation over any closed and bounded interval $[a, b] \subset \mathbb{R}^{+}$ in the ordinary sense and the set of the total variations $V(f,[a, b])$ of $f$ over all such closed and bounded intervals $[a, b]$ is bounded. Furthermore, if this is the case, then the supremum of the total variations over all such closed and bounded intervals is equal to $V(f)$ defined above (see e.g. [18], page 238).

In a similar way, one can define the notion of bounded variation over the intervals of the form $[a, \infty)$, where $a \in \mathbb{R}$ is arbitrary.

Given $f \in \operatorname{BV}\left(\mathbb{R}^{+}\right)$, let $V(f, x):=V(f,[0, x])$ denote the total variation of $f$ over the interval $[0, x]$. Then it is evident that

$$
\begin{equation*}
\lim _{x \rightarrow \infty} V(f, x)=V(f) \tag{2.2}
\end{equation*}
$$

We note that the variation of $f$ over $[x, \infty)$ is given by $V(f,[x, \infty))=V(f)-V(f, x)$ (see e.g. $[9],(9)$ in Lemma 1) and hence from (2.2) it follows that

$$
\lim _{x \rightarrow \infty} V(f,[x, \infty))=0
$$

In this section we prove a theorem concerning definite rate of decay of WalshFourier transform for functions of bounded variation on $\mathbb{R}^{+}$. Our main theorem of this section is as follows.

Theorem 2.1. If $f \in L^{1}\left(\mathbb{R}^{+}\right) \cap \mathrm{BV}\left(\mathbb{R}^{+}\right)$, then $\hat{f}(y)=O(1 / y), y \rightarrow \infty$.

We need the following lemma whose proof is similar to that of Lemma 1 in [9].

Lemma 2.1. If $f \in \operatorname{BV}\left(\mathbb{R}^{+}\right)$and $\left\{a_{n}: n \in \mathbb{N}_{0}\right\}$ is an increasing sequence of non-negative real numbers with $\lim _{n \rightarrow \infty} a_{n}=\infty$, then the series $\sum_{n=1}^{\infty} V\left(f,\left[a_{n-1}, a_{n}\right]\right)$ converges and

$$
\begin{equation*}
V\left(f,\left[a_{0}, \infty\right)\right)=\sum_{n=1}^{\infty} V\left(f,\left[a_{n-1}, a_{n}\right]\right) \tag{2.3}
\end{equation*}
$$

Proof of Theorem 2.1. We present a proof using Taibleson-like technique (see [21]), which we have developed for $\mathbb{\square}$ in [11] and here for $\mathbb{R}^{+}$.

Fix $y \in \mathbb{R}^{+}, y \geqslant 1$. Put $n=[y]$. Then $n \in \mathbb{N}$, so there exists a unique $m \in \mathbb{N}_{0}$ such that $2^{m} \leqslant n<2^{m+1}$. Now, put $a_{i}=i / 2^{m}$ for $i=0,1,2,3, \ldots, 2^{m}$. Then by the definition of Walsh functions, $w_{n}$ takes the value 1 on one half of each of the intervals $\left(a_{i-1}, a_{i}\right)$ and the value -1 on the other half. Therefore we have

$$
\int_{a_{i-1}}^{a_{i}} w_{n}(x) \mathrm{d} x=0 \quad \text { for } i=1,2,3, \ldots, 2^{m}
$$

Since $\psi_{n}$ is a periodic extension of $w_{n}$ from $\mathbb{\square}$ to $\mathbb{R}^{+}$, for each $k \in \mathbb{N}_{0}$ and $i=$ $1,2,3, \ldots, 2^{m}$, we have

$$
\begin{equation*}
\int_{k+a_{i-1}}^{k+a_{i}} \psi_{n}(x) \mathrm{d} x=\int_{a_{i-1}}^{a_{i}} \psi_{n}(x) \mathrm{d} x=\int_{a_{i-1}}^{a_{i}} w_{n}(x) \mathrm{d} x=0 \tag{2.4}
\end{equation*}
$$

Define a step function $g$ on $\mathbb{R}^{+}$by $g(x):=f\left(k+a_{i-1}\right)$ on $\left[k+a_{i-1}, k+a_{i}\right)$, $i=1,2,3, \ldots, 2^{m}, k \in \mathbb{N}_{0}$. Since $f \in L^{1}\left(\mathbb{R}^{+}\right)$, it follows that $g \in L^{1}\left(\mathbb{R}^{+}\right)$. Then in view of (2.4) for each $k \in \mathbb{N}_{0}$ and $i=1,2,3, \ldots, 2^{m}$, we have

$$
\begin{equation*}
\int_{k+a_{i-1}}^{k+a_{i}} g(x) \psi_{n}(x) \mathrm{d} x=f\left(k+a_{i-1}\right) \int_{k+a_{i-1}}^{k+a_{i}} \psi_{n}(x) \mathrm{d} x=0 . \tag{2.5}
\end{equation*}
$$

Therefore by Property (iii) of generalized Walsh functions, as stated above and from (2.5) we have

$$
\begin{align*}
\int_{k+a_{i-1}}^{k+a_{i}} g(x) \psi_{y}(x) \mathrm{d} x & =\int_{k+a_{i-1}}^{k+a_{i}} g(x) \psi_{n}(x) \psi_{[x]}(y) \mathrm{d} x  \tag{2.6}\\
& =\psi_{k}(y) \int_{k+a_{i-1}}^{k+a_{i}} g(x) \psi_{n}(x) \mathrm{d} x=0
\end{align*}
$$

for each $k \in \mathbb{N}_{0}$ and $i=1,2,3, \ldots, 2^{m}$. Now, by definition (1.2) of $\hat{f}(y)$ and (2.6) we have

$$
\begin{aligned}
|\hat{f}(y)| & =\left|\int_{0}^{\infty} f(x) \psi_{y}(x) \mathrm{d} x\right|=\left|\sum_{k=0}^{\infty} \sum_{i=1}^{2^{m}} \int_{k+a_{i-1}}^{k+a_{i}} f(x) \psi_{y}(x) \mathrm{d} x\right| \\
& =\left|\sum_{k=0}^{\infty} \sum_{i=1}^{2^{m}} \int_{k+a_{i-1}}^{k+a_{i}}(f(x)-g(x)) \psi_{y}(x) \mathrm{d} x\right| \leqslant \sum_{k=0}^{\infty} \sum_{i=1}^{2^{m}} \int_{k+a_{i-1}}^{k+a_{i}}|f(x)-g(x)| \mathrm{d} x \\
& =\sum_{k=0}^{\infty} \sum_{i=1}^{2^{m}} \int_{k+a_{i-1}}^{k+a_{i}}\left|f(x)-f\left(k+a_{i-1}\right)\right| \mathrm{d} x \\
& \leqslant \sum_{k=0}^{\infty} \sum_{i=1}^{2^{m}} V\left(f,\left[k+a_{i-1}, k+a_{i}\right]\right)\left(k+a_{i}-\left(k+a_{i-1}\right)\right) \\
& =\sum_{k=0}^{\infty} V(f,[k, k+1]) \frac{1}{2^{m}}=\frac{1}{2^{m}} V(f) \leqslant \frac{2}{n} V(f) \leqslant \frac{4}{y} V(f) .
\end{aligned}
$$

Note that we have used (2.3) of Lemma 2.1 in the last step. Therefore we have

$$
\begin{equation*}
|\hat{f}(y)| \leqslant \frac{4 V(f)}{y} \tag{2.7}
\end{equation*}
$$

This completes the proof of Theorem 2.1.
Problem 2.1. What can be said about the exactness of the constant in (2.7)?

## 3. Two-dimensional case

There is a number of definitions extending the concept of bounded variation for functions of two variables defined on closed and bounded rectangle (see e.g. [1], [2], [12]). We recall one of them.

Let $R:=\left[a_{1}, b_{1}\right] \times\left[a_{2}, b_{2}\right]$ be a closed and bounded rectangle on the real plane $\mathbb{R}^{2}$. We recall that (see e.g. [10], page 21) a collection of points $\left(x_{0}, y_{0}\right),\left(x_{0}, y_{1}\right), \ldots$, $\left(x_{m}, y_{n}\right)$ in $R$, where $m, n \in \mathbb{N}$, satisfying

$$
a_{1}=x_{0} \leqslant x_{1} \leqslant x_{2} \leqslant \ldots \leqslant x_{m}=b_{1} \quad \text { and } \quad a_{2}=y_{0} \leqslant y_{1} \leqslant y_{2} \leqslant \ldots \leqslant y_{n}=b_{2},
$$

is called a collection of grid points of $R$. If $P$ is any such collection of grid points of $R$ and $f: R \rightarrow \mathbb{C}$ is any function, we put

$$
\begin{equation*}
S(P, f)=\sum_{j=1}^{m} \sum_{k=1}^{n}\left|f\left(x_{j}, y_{k}\right)-f\left(x_{j-1}, y_{k}\right)-f\left(x_{j}, y_{k-1}\right)+f\left(x_{j-1}, y_{k-1}\right)\right| \tag{3.1}
\end{equation*}
$$

Now, such a function $f: R \rightarrow \mathbb{C}$ is said to be of bounded variation over the rectangle $R$ in the sense of Vitali (-Lebesgue, -Fréchet, -de la Vallée Poussin, as indicated in [2]), in symbol $f \in \mathrm{BV}_{V}(R)$, if

$$
\begin{equation*}
V(f)=V(f, R):=\sup S(P, f)<\infty \tag{3.2}
\end{equation*}
$$

where the supremum is extended over all collections $P$ of grid points of $R$, while $V(f)$ defined in (3.2) is called the total variation of $f$ over $R$.

Next, we recall the concept of bounded variation for functions on $\left(\mathbb{R}^{+}\right)^{2}$ which is defined as follows (see e.g. [17], Section 2). To do this, analogously to the grid points of a rectangle as defined above, we say that a collection of points $\left(x_{0}, y_{0}\right),\left(x_{0}, y_{1}\right), \ldots,\left(x_{m}, y_{n}\right)$ in $\left(\mathbb{R}^{+}\right)^{2}$, where $m, n \in \mathbb{N}$, satisfying

$$
0 \leqslant x_{0} \leqslant x_{1} \leqslant x_{2} \leqslant \ldots \leqslant x_{m}<\infty
$$

and

$$
0 \leqslant y_{0} \leqslant y_{1} \leqslant y_{2} \leqslant \ldots \leqslant y_{n}<\infty
$$

is called a collection of grid points of $\left(\mathbb{R}^{+}\right)^{2}$. If $P$ is any such collection of grid points of $\left(\mathbb{R}^{+}\right)^{2}$ and $f:\left(\mathbb{R}^{+}\right)^{2} \rightarrow \mathbb{C}$ is any function, we define $S(P, f)$ as in (3.1).

Now, such a function $f:\left(\mathbb{R}^{+}\right)^{2} \rightarrow \mathbb{C}$ is said to be of bounded variation over the set $\left(\mathbb{R}^{+}\right)^{2}$ in the sense of Vitali, in symbol $f \in \mathrm{BV}_{V}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$, if

$$
\begin{equation*}
V(f)=V\left(f,\left(\mathbb{R}^{+}\right)^{2}\right):=\sup S(P, f)<\infty, \tag{3.3}
\end{equation*}
$$

where the supremum is extended over all collections $P$ of grid points of $\left(\mathbb{R}^{+}\right)^{2}$, while $V(f)$ defined in (3.3) is called the total variation of $f$ over $\left(\mathbb{R}^{+}\right)^{2}$.

Similarly to the case of functions $f \in \operatorname{BV}\left(\mathbb{R}^{+}\right)$, the above definition can also be equivalently reformulated as follows. A function $f:\left(\mathbb{R}^{+}\right)^{2} \rightarrow \mathbb{C}$ is of bounded variation over $\left(\mathbb{R}^{+}\right)^{2}$ if and only if $f$ is of bounded variation over all closed and bounded rectangles

$$
\left[a_{1}, b_{1}\right] \times\left[a_{2}, b_{2}\right], \quad 0 \leqslant a_{1}<b_{1}<\infty \text { and } 0 \leqslant a_{2}<b_{2}<\infty
$$

in the sense of Vitali, and in addition, the set of the total variations of $f$ over all such closed and bounded rectangles $\left[a_{1}, b_{1}\right] \times\left[a_{2}, b_{2}\right]$ is bounded. Furthermore, if this is the case, then the supremum of the set of these total variations over all such closed and bounded rectangles $\left[a_{1}, b_{1}\right] \times\left[a_{2}, b_{2}\right]$ is equal to $V(f)$ defined in (3.3).

Next, for a complex-valued Lebesgue integrable function $f$ on $\left(\mathbb{R}^{+}\right)^{2}$, in symbol $f \in L^{1}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$, we consider its Walsh-Fourier transform defined as

$$
\begin{equation*}
\hat{f}(\xi, \eta):=\int_{0}^{\infty} \int_{0}^{\infty} f(x, y) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y, \quad(\xi, \eta) \in\left(\mathbb{R}^{+}\right)^{2} \tag{3.4}
\end{equation*}
$$

We observe that a version of Riemann-Lebesgue lemma holds for the Walsh-Fourier transform defined above. In fact, we have the following.

Lemma 3.1 (Riemann-Lebesgue). If $f \in L^{1}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$, then

$$
\begin{equation*}
\lim _{\xi, \eta \rightarrow \infty} \hat{f}(\xi, \eta)=0 \tag{3.5}
\end{equation*}
$$

Proof. We give a proof of this lemma, which is similar to its one-dimensional version (see [20], page 422). Let $\varepsilon>0$ be given. Since $f \in L^{1}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$, we have

$$
\lim _{m, n \rightarrow \infty} \int_{0}^{m} \int_{0}^{n}|f(x, y)| \mathrm{d} x \mathrm{~d} y=\int_{0}^{\infty} \int_{0}^{\infty}|f(x, y)| \mathrm{d} x \mathrm{~d} y
$$

Therefore, we can choose $m, n$ so large that

$$
\int_{0}^{\infty} \int_{0}^{\infty}|f(x, y)| \mathrm{d} x \mathrm{~d} y-\int_{0}^{m} \int_{0}^{n}|f(x, y)| \mathrm{d} x \mathrm{~d} y<\varepsilon
$$

that is,

$$
\begin{equation*}
\int_{0}^{m} \int_{n}^{\infty}|f(x, y)| \mathrm{d} x \mathrm{~d} y+\int_{m}^{\infty} \int_{0}^{n}|f(x, y)| \mathrm{d} x \mathrm{~d} y+\int_{m}^{\infty} \int_{n}^{\infty}|f(x, y)| \mathrm{d} x \mathrm{~d} y<\varepsilon . \tag{3.6}
\end{equation*}
$$

Now, we notice that

$$
\begin{array}{rl}
\int_{0}^{m} \int_{0}^{n} & f(x, y) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y  \tag{3.7}\\
& =\sum_{k=0}^{m-1} \sum_{l=0}^{n-1} \int_{k}^{k+1} \int_{l}^{l+1} f(x, y) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y \\
& =\sum_{k=0}^{m-1} \sum_{l=0}^{n-1} \int_{k}^{k+1} \int_{l}^{l+1} f(x, y) \psi_{[\xi]}(x) \psi_{[x]}(\xi) \psi_{[\eta]}(y) \psi_{[y]}(\eta) \mathrm{d} x \mathrm{~d} y \\
& =\sum_{k=0}^{m-1} \sum_{l=0}^{n-1} \psi_{k}(\xi) \psi_{l}(\eta) \int_{k}^{k+1} \int_{l}^{l+1} f(x, y) \psi_{[\xi]}(x) \psi_{[\eta]}(y) \mathrm{d} x \mathrm{~d} y \\
& =\sum_{k=0}^{m-1} \sum_{l=0}^{n-1} \psi_{k}(\xi) \psi_{l}(\eta) \int_{0}^{1} \int_{0}^{1} f(x, y) w_{[\xi]}(x) w_{[\eta]}(y) \mathrm{d} x \mathrm{~d} y
\end{array}
$$

In view of (3.6) and (3.7), we see that $\hat{f}(\xi, \eta)$ is dominated by $\varepsilon$ plus a fixed sum of double Walsh-Fourier coefficients of order $([\xi],[\eta])$. Since each of these double Walsh-Fourier coefficients tend to zero as $\xi, \eta \rightarrow \infty$, we conclude that (3.5) holds. This completes the proof of Lemma 3.1.

By the Riemann-Lebesgue lemma, as above, it is certain that the Walsh-Fourier transform $\hat{f}(\xi, \eta) \rightarrow 0$ as $\xi, \eta \rightarrow \infty$. But in general, there is no definite rate at which the Walsh-Fourier transform tends to zero. In fact, the Walsh-Fourier transform of an integrable function on $\left(\mathbb{R}^{+}\right)^{2}$ can tend to zero as slowly as we wish (see e.g. [13], 32.47 (b)). Therefore, as in the one-dimensional case, it is interesting to know for functions of which subclasses of $L^{1}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$ there is a definite rate at which the WalshFourier transform tends to zero. In this section, we carry out this study for functions of bounded variation on $\left(\mathbb{R}^{+}\right)^{2}$ in the sense of Vitali. Our main theorem of this section is as follows.

Theorem 3.1. If $f \in L^{1}\left(\left(\mathbb{R}^{+}\right)^{2}\right) \cap \mathrm{BV}_{V}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$ and $(\xi, \eta) \in\left(\mathbb{R}^{+}\right)^{2}$ is such that $\xi \eta \neq 0$, then

$$
\hat{f}(\xi, \eta)=O\left(\frac{1}{\xi \eta}\right), \quad \xi, \eta \rightarrow \infty .
$$

We need the following lemma, whose proof is similar to that of Lemma 2 in [9].
Lemma 3.2. If $f \in \operatorname{BV}_{V}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$ and if $\left\{a_{n}: n \in \mathbb{N}_{0}\right\}$ and $\left\{b_{n}: n \in \mathbb{N}_{0}\right\}$ are two increasing sequences of non-negative real numbers with

$$
\lim _{n \rightarrow \infty} a_{n}=\lim _{n \rightarrow \infty} b_{n}=\infty
$$

then

$$
V(f)=\sum_{m=1}^{\infty} \sum_{n=1}^{\infty} V\left(f,\left[a_{m-1}, a_{m}\right] \times\left[b_{n-1}, b_{n}\right]\right)
$$

the series on the right-hand side being convergent in the Pringsheim's sense.
Proof of Theorem 3.1. As in the proof of Theorem 2.1, here we present a proof using Taibleson-like technique (see [21]) developed in [7], and developed here for $\left(\mathbb{R}^{+}\right)^{2}$.

Fix $(\xi, \eta) \in\left(\mathbb{R}^{+}\right)^{2}$ with $\xi \geqslant 1, \eta \geqslant 1$. Put $m=[\xi]$ and $n=[\eta]$. Then $m, n \in \mathbb{N}$, so there exist unique $s, t \in \mathbb{N}_{0}$ such that $2^{s} \leqslant m<2^{s+1}$ and $2^{t} \leqslant n<2^{t+1}$.

Now, put $a_{i}=i / 2^{s}$ for $i=0,1,2,3, \ldots, 2^{s}$. Then by the definition of Walsh functions, $w_{m}$ takes the value 1 on one half of each of the intervals ( $a_{i-1}, a_{i}$ ) and the value -1 on the other half. Similarly, if we put $b_{j}=j / 2^{t}$ for $j=0,1,2,3, \ldots, 2^{t}$, then again by the definition of Walsh functions, $w_{n}$ takes the value 1 on one half of each of the intervals $\left(b_{j-1}, b_{j}\right)$ and the value -1 on the other half. Therefore we have

$$
\begin{equation*}
\int_{a_{i-1}}^{a_{i}} w_{m}(x) \mathrm{d} x=0 \quad \text { for } i=1,2,3, \ldots, 2^{s} \tag{3.8}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{b_{j-1}}^{b_{j}} w_{n}(y) \mathrm{d} y=0 \quad \text { for } j=1,2,3, \ldots, 2^{t} \tag{3.9}
\end{equation*}
$$

Since $\psi_{m}$ and $\psi_{n}$ are periodic extensions of $w_{m}$ and $w_{n}$, respectively, from $\mathbb{\square}$ to $\mathbb{R}^{+}$, for each $k, l \in \mathbb{N}_{0}, i=1,2,3, \ldots, 2^{s}$, and $j=1,2,3, \ldots, 2^{t}$, in view of (3.8)-(3.9), we have

$$
\begin{equation*}
\int_{k+a_{i-1}}^{k+a_{i}} \psi_{m}(x) \mathrm{d} x=\int_{a_{i-1}}^{a_{i}} \psi_{m}(x) \mathrm{d} x=\int_{a_{i-1}}^{a_{i}} w_{m}(x) \mathrm{d} x=0 \tag{3.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{l+b_{j-1}}^{l+b_{j}} \psi_{n}(y) \mathrm{d} y=\int_{b_{j-1}}^{b_{j}} \psi_{n}(y) \mathrm{d} y=\int_{b_{j-1}}^{b_{j}} w_{n}(y) \mathrm{d} y=0 . \tag{3.11}
\end{equation*}
$$

Define three functions $f_{1}, f_{2}, f_{3}$ on $\left(\mathbb{R}^{+}\right)^{2}$ by setting

$$
f_{1}(x, y):=f\left(x, l+b_{j-1}\right), \quad x \in \mathbb{R}^{+}, l+b_{j-1} \leqslant y<l+b_{j}
$$

for $j=1,2,3, \ldots, 2^{t}, l \in \mathbb{N}_{0}$;

$$
f_{2}(x, y):=f\left(k+a_{i-1}, y\right), \quad k+a_{i-1} \leqslant x<k+a_{i}, y \in \mathbb{R}^{+}
$$

for $i=1,2,3, \ldots, 2^{s}, k \in \mathbb{N}_{0}$; and

$$
f_{3}(x, y):=f\left(k+a_{i-1}, l+b_{j-1}\right), \quad k+a_{i-1} \leqslant x<k+a_{i}, l+b_{j-1} \leqslant y<l+b_{j}
$$

for $i=1,2,3, \ldots, 2^{s}, j=1,2,3, \ldots, 2^{t}$, and $k, l \in \mathbb{N}_{0}$. Since $f \in L^{1}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$, it follows that $f_{1}, f_{2}, f_{3} \in L^{1}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$. Now, in view of Fubini's theorem and relations (3.10)-(3.11), for each $i=1,2,3, \ldots, 2^{s}, j=1,2,3, \ldots, 2^{t}$ and $k, l \in \mathbb{N}_{0}$, we have

$$
\begin{align*}
\int_{k+a_{i-1}}^{k+a_{i}} & \int_{l+b_{j-1}}^{l+b_{j}} f_{1}(x, y) \psi_{\xi}(x) \psi_{n}(y) \mathrm{d} x \mathrm{~d} y  \tag{3.12}\\
& =\int_{k+a_{i-1}}^{k+a_{i}}\left(f\left(x, l+b_{j-1}\right) \int_{l+b_{j-1}}^{l+b_{j}} \psi_{n}(y) \mathrm{d} y\right) \psi_{\xi}(x) \mathrm{d} x=0
\end{align*}
$$

$$
\begin{align*}
\int_{k+a_{i-1}}^{k+a_{i}} & \int_{l+b_{j-1}}^{l+b_{j}} f_{2}(x, y) \psi_{m}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y  \tag{3.13}\\
& =\int_{l+b_{j-1}}^{l+b_{j}}\left(f\left(k+a_{i-1}, y\right) \int_{k+a_{i-1}}^{k+a_{i}} \psi_{m}(x) \mathrm{d} x\right) \psi_{\eta}(y) \mathrm{d} y=0
\end{align*}
$$

and

$$
\begin{align*}
& \int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} f_{3}(x, y) \psi_{m}(x) \psi_{n}(y) \mathrm{d} x \mathrm{~d} y  \tag{3.14}\\
& \quad=f\left(k+a_{i-1}, l+b_{j-1}\right)\left(\int_{k+a_{i-1}}^{k+a_{i}} \psi_{m}(x) \mathrm{d} x\right)\left(\int_{l+b_{j-1}}^{l+b_{j}} \psi_{n}(y) \mathrm{d} y\right)=0
\end{align*}
$$

Therefore by Property (iii) of generalized Walsh functions and (3.12)-(3.14) for each $i=1,2,3, \ldots, 2^{s}, j=1,2,3, \ldots, 2^{t}$, and $k, l \in \mathbb{N}_{0}$, we have

$$
\begin{align*}
\int_{k+a_{i-1}}^{k+a_{i}} & \int_{l+b_{j-1}}^{l+b_{j}} f_{1}(x, y) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y  \tag{3.15}\\
& =\int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} f_{1}(x, y) \psi_{\xi}(x) \psi_{n}(y) \psi_{[y]}(\eta) \mathrm{d} x \mathrm{~d} y \\
& =\psi_{l}(\eta) \int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} f_{1}(x, y) \psi_{\xi}(x) \psi_{n}(y) \mathrm{d} x \mathrm{~d} y=0
\end{align*}
$$

$$
\begin{align*}
\int_{k+a_{i-1}}^{k+a_{i}} & \int_{l+b_{j-1}}^{l+b_{j}} f_{2}(x, y) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y  \tag{3.16}\\
& =\int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} f_{2}(x, y) \psi_{m}(x) \psi_{[x]}(\xi) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y \\
& =\psi_{k}(\xi) \int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} f_{2}(x, y) \psi_{m}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y=0
\end{align*}
$$

and

$$
\begin{align*}
\int_{k+a_{i-1}}^{k+a_{i}} & \int_{l+b_{j-1}}^{l+b_{j}} f_{3}(x, y) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y  \tag{3.17}\\
& =\int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} f_{3}(x, y) \psi_{m}(x) \psi_{[x]}(\xi) \psi_{n}(y) \psi_{[y]}(\eta) \mathrm{d} x \mathrm{~d} y \\
& =\psi_{k}(\xi) \psi_{l}(\eta) \int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} f_{3}(x, y) \psi_{m}(x) \psi_{n}(y) \mathrm{d} x \mathrm{~d} y=0 .
\end{align*}
$$

Using (3.15)-(3.17) in definition (3.4) of $\hat{f}(\xi, \eta)$ we get

$$
\begin{aligned}
|\hat{f}(\xi, \eta)| & =\left|\int_{0}^{\infty} \int_{0}^{\infty} f(x, y) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y\right| \\
& =\left|\sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \sum_{i=1}^{2^{s}} \sum_{j=1}^{2^{t}} \int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} f(x, y) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y\right|
\end{aligned}
$$

$$
\begin{aligned}
& =\mid \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \sum_{i=1}^{2^{s}} \sum_{j=1}^{2^{t}} \int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}}\left(f(x, y)-f_{1}(x, y)\right. \\
& \left.-f_{2}(x, y)+f_{3}(x, y)\right) \psi_{\xi}(x) \psi_{\eta}(y) \mathrm{d} x \mathrm{~d} y \mid \\
& \leqslant \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \sum_{i=1}^{2^{s}} \sum_{j=1}^{2^{t}} \int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} \mid f(x, y)-f_{1}(x, y) \\
& -f_{2}(x, y)+f_{3}(x, y) \mid \mathrm{d} x \mathrm{~d} y \\
& =\sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \sum_{i=1}^{2^{s}} \sum_{j=1}^{2^{t}} \int_{k+a_{i-1}}^{k+a_{i}} \int_{l+b_{j-1}}^{l+b_{j}} \mid f(x, y)-f\left(x, l+b_{j-1}\right) \\
& -f\left(k+a_{i-1}, y\right)+f\left(k+a_{i-1}, l+b_{j-1}\right) \mid \mathrm{d} x \mathrm{~d} y \\
& \leqslant \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \sum_{i=1}^{2^{s}} \sum_{j=1}^{2^{t}} V\left(f,\left[k+a_{i-1}, k+a_{i}\right] \times\left[l+b_{j-1}, l+b_{j}\right]\right) \\
& \times\left(a_{i}-a_{i-1}\right)\left(b_{j}-b_{j-1}\right) \\
& \leqslant \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} V(f,[k, k+1] \times[l, l+1]) \frac{1}{2^{s}} \frac{1}{2^{t}} \\
& =\frac{1}{2^{s} 2^{t}} V(f) \leqslant \frac{4}{m n} V(f) \leqslant \frac{16 V(f)}{\xi \eta},
\end{aligned}
$$

in view of Lemma 3.2. Thus, we get

$$
\begin{equation*}
|\hat{f}(\xi, \eta)| \leqslant \frac{16 V(f)}{\xi \eta} \tag{3.18}
\end{equation*}
$$

The proof of Theorem 3.1 is complete.
Problem 3.1. How to estimate $\hat{f}(\xi, 0), \xi \neq 0$ (or $\hat{f}(0, \eta), \eta \neq 0)$ in terms of $\xi$ (or $\eta$ ), even assuming that $f$ is of bounded variation over $\left(\mathbb{R}^{+}\right)^{2}$ in the sense of Hardy (see [17] for definition)?

Problem 3.2. What can be said about the exactness of the constant in (3.18)?

## 4. Extension of the Result $\operatorname{to}\left(\mathbb{R}^{+}\right)^{N}, N \in \mathbb{N}$

We start by defining the concept of bounded variation for functions on $\left(\mathbb{R}^{+}\right)^{N}$, $N \in \mathbb{N}$ in the sense of Vitali.

For a function $f:\left(\mathbb{R}^{+}\right)^{N} \rightarrow \mathbb{C}$ and for any rectangle $R=\left[\alpha_{1}, \beta_{1}\right] \times \ldots \times\left[\alpha_{N}, \beta_{N}\right]$ in $\left(\mathbb{R}^{+}\right)^{N}$ with $0 \leqslant \alpha_{i}<\beta_{i}<\infty$ for all $i=1,2, \ldots, N$, we define $\Delta f(R)$ as follows:

When $N=2$, we put

$$
\Delta f(R):=\Delta f\left(\left[\alpha_{1}, \beta_{1}\right] \times\left[\alpha_{2}, \beta_{2}\right]\right)=f\left(\beta_{1}, \beta_{2}\right)-f\left(\beta_{1}, \alpha_{2}\right)-f\left(\alpha_{1}, \beta_{2}\right)+f\left(\alpha_{1}, \alpha_{2}\right) ;
$$

for $N=3$

$$
\begin{aligned}
\Delta f(R):= & \Delta f\left(\left[\alpha_{1}, \beta_{1}\right] \times\left[\alpha_{2}, \beta_{2}\right] \times\left[\alpha_{3}, \beta_{3}\right]\right) \\
= & \left(f\left(\beta_{1}, \beta_{2}, \beta_{3}\right)-f\left(\beta_{1}, \alpha_{2}, \beta_{3}\right)-f\left(\alpha_{1}, \beta_{2}, \beta_{3}\right)+f\left(\alpha_{1}, \alpha_{2}, \beta_{3}\right)\right) \\
& -\left(f\left(\beta_{1}, \beta_{2}, \alpha_{3}\right)-f\left(\beta_{1}, \alpha_{2}, \alpha_{3}\right)-f\left(\alpha_{1}, \beta_{2}, \alpha_{3}\right)+f\left(\alpha_{1}, \alpha_{2}, \alpha_{3}\right)\right) \\
= & \Delta_{\left[\alpha_{3}, \beta_{3}\right]} \Delta f\left(\left[\alpha_{1}, \beta_{1}\right] \times\left[\alpha_{2}, \beta_{2}\right]\right)
\end{aligned}
$$

and successively for any $N \geqslant 3$

$$
\begin{aligned}
\Delta f(R) & :=\Delta f\left(\left[\alpha_{1}, \beta_{1}\right] \times \ldots \times\left[\alpha_{N}, \beta_{N}\right]\right) \\
& =\Delta_{\left[\alpha_{N}, \beta_{N}\right]} \Delta f\left(\left[\alpha_{1}, \beta_{1}\right] \times \ldots \times\left[\alpha_{N-1}, \beta_{N-1}\right]\right)
\end{aligned}
$$

A collection of points $\left(x_{1}^{0}, \ldots, x_{N}^{0}\right),\left(x_{1}^{0}, \ldots, x_{N-1}^{0}, x_{N}^{1}\right), \ldots,\left(x_{1}^{s_{1}}, \ldots, x_{N}^{s_{N}}\right)$ of $\left(\mathbb{R}^{+}\right)^{N}$ satisfying

$$
0 \leqslant x_{j}^{0} \leqslant x_{j}^{1} \leqslant \ldots \leqslant x_{j}^{s_{j}}<\infty, \quad s_{j} \in \mathbb{N}, j=1,2, \ldots, N
$$

is called a collection of grid points of $\left(\mathbb{R}^{+}\right)^{N}$. If $P$ is any such collection of grid points of $\left(\mathbb{R}^{+}\right)^{N}$ and $f:\left(\mathbb{R}^{+}\right)^{N} \rightarrow \mathbb{C}$ is any function, we put

$$
S(P, f)=\sum_{i_{1}=1}^{s_{1}} \ldots \sum_{i_{N}=1}^{s_{N}}\left|\Delta f\left(\left[x_{1}^{i_{1}-1}, x_{1}^{i_{1}}\right] \times \ldots \times\left[x_{N}^{i_{N}-1}, x_{N}^{i_{N}}\right]\right)\right|
$$

Now, a function $f:\left(\mathbb{R}^{+}\right)^{N} \rightarrow \mathbb{C}$ is said to be of bounded variation over the set $\left(\mathbb{R}^{+}\right)^{N}$ in the sense of Vitali, in symbol $f \in \operatorname{BV}_{V}\left(\left(\mathbb{R}^{+}\right)^{N}\right)$, if

$$
\begin{equation*}
V(f)=V\left(f,\left(\mathbb{R}^{+}\right)^{N}\right):=\sup S(P, f)<\infty \tag{4.1}
\end{equation*}
$$

where the supremum is extended over all collections $P$ of grid points of $\left(\mathbb{R}^{+}\right)^{N}$, while $V(f)$ defined in (4.1) is called the total variation of $f$ over $\left(\mathbb{R}^{+}\right)^{N}$.

Similarly to the case of functions $f \in \mathrm{BV}\left(\mathbb{R}^{+}\right)$and $f \in \mathrm{BV}_{V}\left(\left(\mathbb{R}^{+}\right)^{2}\right)$, the above definition can also be equivalently reformulated as follows. A function $f:\left(\mathbb{R}^{+}\right)^{N} \rightarrow \mathbb{C}$ is of bounded variation over $\left(\mathbb{R}^{+}\right)^{N}$ in the sense of Vitali if and only if $f$ is of bounded variation over all closed and bounded $N$-rectangles

$$
\left[a_{1}, b_{1}\right] \times \ldots \times\left[a_{N}, b_{N}\right], \quad 0 \leqslant a_{i}<b_{i}<\infty, i=1,2, \ldots, N
$$

in the sense of Vitali (see e.g. [14] or [5] for definition), and in addition, the set of total variations of $f$ over all such closed and bounded $N$-rectangles $\left[a_{1}, b_{1}\right] \times \ldots \times\left[a_{N}, b_{N}\right]$ is bounded. Furthermore, if this is the case, then the supremum of the set of these total variations over all such closed and bounded $N$-rectangles $\left[a_{1}, b_{1}\right] \times \ldots \times\left[a_{N}, b_{N}\right]$ is equal to $V(f)$ defined in (4.1).

Next, for a complex-valued Lebesgue integrable function $f$ on $\left(\mathbb{R}^{+}\right)^{N}$, in symbol $f \in L^{1}\left(\left(\mathbb{R}^{+}\right)^{N}\right)$, we consider its Walsh-Fourier transform defined as

$$
\hat{f}\left(\xi_{1}, \ldots, \xi_{N}\right):=\int_{\left(\mathbb{R}^{+}\right)^{N}} f\left(x_{1}, \ldots, x_{N}\right) \psi_{\xi_{1}}\left(x_{1}\right) \ldots \psi_{\xi_{N}}\left(x_{N}\right) \mathrm{d} x_{1} \ldots \mathrm{~d} x_{N}
$$

where $\left(\xi_{1}, \ldots, \xi_{N}\right) \in\left(\mathbb{R}^{+}\right)^{N}$.
In this case also, similar to Lemma 3.1, a version of Riemann-Lebesgue lemma holds, that is, $\hat{f}\left(\xi_{1}, \ldots, \xi_{N}\right) \rightarrow 0$ as $\xi_{1}, \ldots, \xi_{N} \rightarrow \infty$. But in general, there is no definite rate at which the Walsh-Fourier transform tends to zero. In fact, the WalshFourier transform of an integrable function on $\left(\mathbb{R}^{+}\right)^{N}$ can tend to zero as slowly as we wish (see e.g. [13], 32.47 (b)). Therefore, as in one and two dimensional cases, it is interesting to know for functions of which subclasses of $L^{1}\left(\left(\mathbb{R}^{+}\right)^{N}\right)$ there is a definite rate at which the Walsh-Fourier transform tends to zero. In this section, we state the result for functions of bounded variation on $\left(\mathbb{R}^{+}\right)^{N}$ in the sense of Vitali, which is an extension of our theorems in Sections 2 and 3. The proof of this theorem is similar to that of Theorem 3.1.

Theorem 4.1. If $f \in L^{1}\left(\left(\mathbb{R}^{+}\right)^{N}\right) \cap \operatorname{BV}_{V}\left(\left(\mathbb{R}^{+}\right)^{N}\right)$ and $\left(\xi_{1}, \ldots, \xi_{N}\right) \in\left(\mathbb{R}^{+}\right)^{N}$ is such that $\prod_{i=1}^{N} \xi_{i} \neq 0$, then

$$
\hat{f}\left(\xi_{1}, \ldots, \xi_{N}\right)=O\left(1 / \prod_{i=1}^{N} \xi_{i}\right), \quad \xi_{1}, \ldots, \xi_{N} \rightarrow \infty
$$

More precisely,

$$
\begin{equation*}
\left|\hat{f}\left(\xi_{1}, \ldots, \xi_{N}\right)\right| \leqslant 4^{N} V(f) / \prod_{i=1}^{N} \xi_{i}, \quad \xi_{1}, \ldots, \xi_{N} \geqslant 1 \tag{4.2}
\end{equation*}
$$

Problem 4.1. How to estimate $\hat{f}\left(\xi_{1}, \ldots, \xi_{N}\right)$ if $\left(\xi_{1}, \ldots, \xi_{N}\right) \neq(0, \ldots, 0)$, but $\xi_{j}=0$ for some $j \in\{1, \ldots, N\}$, even assuming that $f$ is of bounded variation over $\left(\mathbb{R}^{+}\right)^{N}$ in the sense of Hardy (which can be defined similarly as in the case of $\left.\left(\mathbb{R}^{+}\right)^{2}\right)$ ?

Problem 4.2. What can be said about the exactness of the constant in (4.2)?

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[^0]:    This research was completed while the first author was visiting Bolyai Institute, University of Szeged, Szeged, Hungary, under the Hungarian State Scholarship Grant Award during the academic year 2016-2017 between May 16, 2017 and June 15, 2017.

